

A duality principle in spin glasses

Antonio Auffinger* Wei-Kuo Chen†

Abstract

We prove a duality principle that connects the thermodynamic limits of the free energies of the Hamiltonians and their squared interactions. Under the main assumption that the limiting free energy is concave in the squared temperature parameter, we show that this relation is valid in a large class of disordered systems. In particular, when applied to mean field spin glasses, this duality provides an interpretation of the Parisi formula as an inverted variational principle, establishing a prediction of Guerra [13].

Keywords: large deviation; Legendre duality; Parisi formula; spin glass.

AMS MSC 2010: 60F10; 82D30.

Submitted to EJP on February 5, 2017, final version accepted on May 21, 2017.

1 Introduction

A fundamental goal in spin glass theory is to study the thermodynamic limit of random Hamiltonian systems that simultaneously exhibit ferromagnetic and anti-ferromagnetic properties. Roughly speaking, one seeks to obtain information on macroscopic phenomena such as phase transitions through the analysis of stochastic microscopic interactions. Applications of spin glass theory include many problems in the fields of physics, biology, neurology and computer science, see for instance [14]. More importantly, it also has led to a vast and challenging collection of beautiful mathematical questions [24].

For each $N \geq 1$, let Σ_N be a configuration space and ν_N be a random measure on Σ_N . Let H_N be a random Hamiltonian indexed by Σ_N . In this paper, we study the thermodynamic limit of the free energy associated to the (inverse) temperature $\beta \in \mathbb{R}$,

$$F(\beta) := \lim_{N \rightarrow \infty} \frac{1}{N} \log Z_N(\beta),$$

where

$$Z_N(\beta) := \int_{\Sigma_N} \exp(\beta H_N(\sigma)) \nu_N(d\sigma)$$

is called the partition function. There are two classical viewpoints that are often employed to guess or compute an expression for this limit. The first one is physicists'

*Northwestern University, United States of America. E-mail: auffing@math.northwestern.edu

†University of Minnesota, United States of America. E-mail: wkchen@umn.edu

entropic principle in statistical mechanics. It suggests that $F(\beta)$ has a variational representation that involves a maximization between the entropy of a generic thermodynamical state and its internal energy. Another approach is via large deviation theory, from which Laplace-Varadhan's lemma writes $F(\beta)$ as a maximization problem between an energy functional and a rate function. The reader interested in these classical approaches can see [10, 11] and the references therein.

In contrast to the classical methods, in the ground breaking treatment of the famous Sherrington-Kirkpatrick (SK) model, Giorgio Parisi [22] predicted that the thermodynamic limit of the free energy can be computed through a sharply different variational problem. Introducing the notion of replica symmetry breaking, he expressed $F(\beta)$ as a minimization problem of a functional \mathcal{P}_β over the space of all probability measures on the interval $[0, 1]$. The minimizer of this variational problem is known as the functional order parameter and plays a fundamental role in describing the system. Following Guerra's beautiful discovery of the replica symmetry breaking bound, Parisi's formula was firstly verified in the celebrated work of Talagrand [25] and then extended to all mixed p -spin models by Panchenko [18]. Further generalizations have also been pushed forward in the generalized random energy model [4, 5], the spherical mixed p -spin model [6], the Ghatak-Sherrington model [15], the multi-species SK model [19], the mixed p -spin model with vector spins [20] and the Potts spin glass model [21].

Despite all remarkable progress during the past 35 years, some aspects of the Parisi solution remain to be understood. In particular, it is unclear how to explain the Parisi solution directly from the aforementioned classical methods. The aim of this paper is to provide a rigorous framework that connects the Parisi formula and the classical approaches. It is motivated by an observation of Francesco Guerra [13], who suggested that the thermodynamic limit of the free energy in the mixed p -spin model is concave in the squared temperature. From such concavity, he conjectured a Legendre duality between the Parisi formula and the Legendre transform Γ of the scaled Parisi functional $\mathcal{P}_{\sqrt{\beta}}$, where the temperature and the functional order parameter are conjugate variables. This duality was rigorously established for the mixed p -spin models in the recent work of Auffinger and Chen [3], but no interpretation of the functional Γ was given.

In this paper, we show that Guerra's insight not only holds for the mixed p -spin glass model, but also extends to many disordered models. We establish that if the limiting free energy in the squared temperature is concave, then the system will exhibit a general Legendre duality principle between the limiting free energies corresponding to the original and squared Hamiltonians. Analogously to the classical entropic principle where temperature and energy parameters are Legendre conjugates, this duality also connects these two parameters. As an immediate consequence, we recover the Legendre structure obtained in [3]. Foremost, we give a direct interpretation of the functional Γ , as the limiting free energy of the squared Hamiltonian. While the Parisi formula relates temperature parameter with the Parisi measure, our new representation conjugates energy and temperature parameters, as in the classical approaches. From these, we also derive a new maximum variational representation for the limiting free energy, which shares some features similar to those obtained through large deviation principles in many classical examples, see [11, 23].

In addition, we apply our Legendre duality to the random energy model, from which we obtain the corresponding Parisi formula previously derived by Guerra [12] and Bolthausen-Kistler [4]. Our result extends to the spherical mixed p -spin model as well, where it explains the nature of the Crisanti-Sommers representation of the limiting free energy.

Acknowledgements. Both authors thank Francesco Guerra for bringing their attention to the papers [12, 13] and for several useful suggestions. We also thank Dmitry Panchenko

for valuable comments on the presentation of the paper. The research of A. A. is partly supported by NSF grant DMS-1597864 and NSF CAREER Grant DMS-1653552. The research of W.-K. C. is partly supported by NSF grant DMS-1642207 and Hong Kong Research Grants Council GRF-14302515.

2 Legendre duality

2.1 Main results

For each $N \geq 1$, let Σ_N be a configuration space and ν_N be a random measure on Σ_N with $\nu_N(\Sigma_N) < \infty$. Suppose that $(H_{N,p})_{p \geq 1}$ is a sequence of independent random Hamiltonians indexed by Σ_N and is independent of ν_N . Denote the normed vector space $\ell^2(\mathbb{N})$ as

$$\mathcal{T} = \left\{ \beta = (\beta_p)_{p \geq 1} : \sum_{p \geq 1} \beta_p^2 < \infty \right\}.$$

The sequence β here is called the (inverse) temperature parameter. Throughout Section 2, we assume that the following conditions hold:

(A0) Suppose that for each $p \geq 1$, the Hamiltonian $H_{N,p}$ is symmetric, that is, the families of random variables $(H_{N,p}(\sigma))_{\sigma \in \Sigma_N}$ and $(-H_{N,p}(\sigma))_{\sigma \in \Sigma_N}$ share the same joint distribution.

(A1) There exist constants $C_1, C_2 > 0$ such that for any $N \geq 1$ and $p \geq 1$,

$$\mathbb{E} \sup_{\sigma \in \Sigma_N} H_{N,p}(\sigma) \leq \frac{C_1 N}{2^{p/2}},$$

and

$$\mathbb{P} \left(\left| \sup_{\sigma \in \Sigma_N} H_{N,p}(\sigma) - \mathbb{E} \sup_{\sigma \in \Sigma_N} H_{N,p}(\sigma) \right| \geq t \right) \leq C_2 \exp \left(-\frac{2^p t^2}{C_2 N} \right), \quad \forall t > 0.$$

(A2) For each $\beta \in \mathcal{T}$, the limit of the free energy

$$F_N(\beta) := \frac{1}{N} \log \int_{\Sigma_N} \exp \sum_{p \geq 1} \beta_p H_{N,p}(\sigma) \nu_N(d\sigma)$$

converges a.s. to some nonrandom $F(\beta)$.

Remark 2.1 (Universality). Throughout the paper, we neither impose any other assumptions on the Hamiltonians $(H_{N,p})$, nor on the measure space (Σ_N, ν_N) . In particular, our results do not require $H_{N,p}$ to be a Gaussian process.

In the main examples that we are about to discuss in Subsection 2.2 below, the configuration space Σ_N is either the discrete hypercube $\{\pm 1\}^N$ or a sphere in \mathbb{R}^N and the Hamiltonians are the mixed p -spin models. They are known to satisfy assumptions (A0), (A1) and (A2).

Our first main theorem establishes a general duality principle that conjugates the temperature and energy parameters. Let

$$\mathcal{B} = \left\{ \beta = (\beta_p)_{p \geq 1} \in \mathcal{T} : \beta_p > 0, \forall p \geq 1 \right\}$$

be the space of positive temperatures and

$$\mathcal{M} = \left\{ \mathbf{m} = (m_p)_{p \geq 1} : m_p > 0, \forall p \geq 1 \right\}$$

denote the space of scaled energies. Define the free energy for the squared Hamiltonian,

$$V_N(\mathbf{m}) = \frac{1}{N} \log \int_{\Sigma_N} \exp \sum_{p \geq 1} \frac{N}{2m_p} \left(\frac{H_{N,p}(\sigma)}{N} \right)^2 \nu_N(d\sigma), \quad \forall \mathbf{m} \in \mathcal{M}. \quad (2.1)$$

Denote by V the limit of $(V_N)_{N \geq 1}$ whenever it exists in $\bar{\mathbb{R}} := \mathbb{R} \cup \{\infty\}$. For any $\beta \in \mathcal{B}$, $\mathbf{m} \in \mathcal{M}$ and $f : [0, \infty)^2 \rightarrow \mathbb{R}$, we set $f(\beta, \mathbf{m}) = (f(\beta_p, m_p))_{p \geq 1}$. We say that F is Gâteaux differentiable at β if

$$\lim_{t \rightarrow 0} \frac{F(\beta + t\beta') - F(\beta)}{t}$$

exists for any $\beta' \in \mathcal{T}$. Our first result shows that V exists and it can be expressed as a variational formula in terms of the free energy F . Moreover, if $F(\beta^{1/2})$ is concave in $\beta \in \mathcal{B}$, then this variational representation is invertible.

Theorem 2.2 (Legendre duality). *If (A0), (A1) and (A2) hold, then we have*

(i) *For any $\mathbf{m} \in \mathcal{M}$,*

$$V(\mathbf{m}) = \sup_{\beta \in \mathcal{B}} \left(F(\beta) - \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m_p \right). \quad (2.2)$$

(ii) *If $F(\beta)$ is Gâteaux differentiable everywhere in \mathcal{B} and $F(\beta^{1/2})$ is concave on \mathcal{B} , then for any $\beta \in \mathcal{B}$,*

$$F(\beta) = \inf_{\mathbf{m} \in \mathcal{M}} \left(V(\mathbf{m}) + \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m_p \right). \quad (2.3)$$

Theorem 2.2 establishes the Legendre duality between F and V , where the temperature and the scaled energy are conjugated variables.

Remark 2.3. It is clear that the variational representation (2.3) is concave in $\beta^{1/2}$. A simple application of Hölder's inequality also implies the convexity of F . Such convexity can be obtained from the variational formula of (2.3). Indeed, if we set $\mathbf{t} = (t_p)_{p \geq 1}$ for $t_p = \beta_p^2 m_p$,

$$\inf_{\mathbf{m} \in \mathcal{M}} \left(V(\mathbf{m}) + \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m_p \right) = \inf_{\mathbf{t} \in \mathcal{M}} \left(U\left(\frac{\beta^2}{\mathbf{t}}\right) + \frac{1}{2} \sum_{p \geq 1} t_p \right),$$

where $U(\mathbf{m}) := V(\mathbf{m}^{-1})$ for any $\mathbf{m} \in \mathcal{M}$. For any $0 < c < 1$, $\beta, \beta' \in \mathcal{B}$ and $\mathbf{t}, \mathbf{t}' \in \mathcal{M}$, the Cauchy-Schwarz inequality says

$$\frac{((1-c)\beta_p + c\beta'_p)^2}{(1-c)t_p + ct'_p} \leq (1-c)\frac{\beta_p^2}{t_p} + c\frac{\beta'^2_p}{t'_p}.$$

Observe that from (2.2), U is nondecreasing in each coordinate and convex. This implies that

$$\begin{aligned} & U\left(\frac{((1-c)\beta + c\beta')^2}{(1-c)\mathbf{t} + c\mathbf{t}'}\right) + \frac{1}{2} \sum_{p \geq 1} ((1-c)t_p + ct'_p) \\ & \leq (1-c)\left(U\left(\frac{\beta^2}{\mathbf{t}}\right) + \frac{1}{2} \sum_{p \geq 1} t_p\right) + c\left(U\left(\frac{\beta'^2}{\mathbf{t}'}\right) + \frac{1}{2} \sum_{p \geq 1} t'_p\right), \end{aligned}$$

which clearly implies the convexity of F by taking infimum on both sides.

The next theorem describes the optimizers of the above two variational problems. Assume that F is Gâteaux differentiable on \mathcal{B} . Denote by $\partial_p F$ the partial derivative of F with respect to β_p . We define

$$\mathcal{B}_0 = \left\{ \beta \in \mathcal{B} : \beta_p^{-1} \partial_p F(\beta) > 0, \forall p \geq 1 \right\}$$

and

$$\mathcal{M}_0 = \left\{ m \in \mathcal{M} : m_p = \beta_p^{-1} \partial_p F(\beta) > 0, \forall p \geq 1 \text{ for some } \beta \in \mathcal{B}_0 \right\}.$$

Theorem 2.4 (Optimality). *Given the assumptions (A0), (A1) and (A2), if F is Gâteaux differentiable everywhere in \mathcal{B} and $F(\beta^{1/2})$ is concave on \mathcal{B} , then we have*

(i) For any $m \in \mathcal{M}_0$,

$$V(m) = \max_{\beta \in \mathcal{B}_0} \left(F(\beta) - \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m_p \right) \tag{2.4}$$

and any $\beta \in \mathcal{B}$ satisfying $m_p = \beta_p^{-1} \partial_p F(\beta)$ for all $p \geq 1$ is a maximizer.

(ii) For any $\beta \in \mathcal{B}_0$,

$$F(\beta) = \min_{m \in \mathcal{M}_0} \left(V(m) + \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m_p \right) \tag{2.5}$$

and $m \in \mathcal{M}$ defined by $m_p = \beta_p^{-1} \partial_p F(\beta)$ for all $p \geq 1$ is a minimizer.

A word of comment is needed here. The dual parameter m plays the role of a scaled energy. If both free energies F and V are comparable, that is, if the energies satisfy

$$\beta_p H_{N,p} \sim \frac{1}{m_p} H_{N,p}^2,$$

we can see that the parameter m_p has the same scale as $H_{N,p}/\beta_p$.

Remark that although here we present results for the system with an infinite sequence of the Hamiltonians $(H_{N,p})_{p \geq 1}$, the statements of Theorems 2.2 and 2.4 remain valid if the system involves only a finite number of Hamiltonians, $(H_{N,p})_{1 \leq p \leq s}$, for some fixed $s \in \mathbb{N}$ independent of N . The case $s = 1$ reads as follows.

Assume that H_N is a Hamiltonian indexed by the spin configuration space Σ_N and is independent of ν_N . Define the free energies by

$$\begin{aligned} F_N(\beta) &= \frac{1}{N} \log \int_{\Sigma_N} \exp \beta H_N(\sigma) \nu_N(d\sigma), \quad \forall \beta \in \mathbb{R}, \\ V_N(m) &= \frac{1}{N} \log \int_{\Sigma_N} \exp \frac{N}{2m} \left(\frac{H_N(\sigma)}{N} \right)^2 \nu_N(d\sigma), \quad \forall m > 0. \end{aligned} \tag{2.6}$$

Suppose that (A0), (A1) and (A2) are valid for H_N . Denote by F the limit of (F_N) . From Theorem 2.2, the limit of V_N exists, which is denoted by V . Theorem 2.4 becomes:

Theorem 2.5. *If $F(\beta)$ is differentiable and $F(\beta^{1/2})$ is concave in $\beta > 0$, then for any $\beta \in \mathcal{B}_0 := \{\beta' > 0 : F'(\beta') > 0\}$ and $m \in \mathcal{M}_0 := \{F'(\beta')/\beta' : \text{for some } \beta' > 0\}$, we have*

$$V(m) = \max_{\beta' \in \mathcal{B}_0} \left(F(\beta') - \frac{\beta'^2 m}{2} \right), \tag{2.7}$$

$$F(\beta) = \min_{m' \in \mathcal{M}_0} \left(V(m') + \frac{\beta^2 m'}{2} \right), \tag{2.8}$$

where any $\beta' > 0$ satisfying $m = F'(\beta')/\beta'$ is the maximizer of (2.7) and $m' = F'(\beta)/\beta$ is the minimizer of (2.8).

2.2 Examples

We now discuss the Legendre structure established in the previous section for some mean-field spin glass models. They include the random energy model (REM) and the mixed p -spin models with both Ising and spherical spin configuration spaces. With

the representations in Theorem 2.4, we will show that the limiting free energy can be naturally written as a minimization problem that is related to Parisi's formulation of the same quantity.

2.2.1 Random energy model

The REM model is defined on the configuration space $\Sigma_N := \{-1, 1\}^N$ and its Hamiltonian H_N is a family of Gaussian process with covariance structure $\mathbb{E}H_N(\sigma^1)H_N(\sigma^2) = N\delta_{\sigma^1, \sigma^2}$ for all $\sigma^1, \sigma^2 \in \Sigma_N$, where $\delta_{\sigma^1, \sigma^2}$ equals 1 if $\sigma^1 = \sigma^2$ and 0 otherwise. Let ν_N be the uniform probability measure on Σ_N . It is well-known (see for instance [24]) that the conditions (A1) and (A2) hold and the limiting free energy has an explicit expression, for $\beta_c := \sqrt{2 \log 2}$,

$$F(\beta) = \lim_{N \rightarrow \infty} \frac{1}{N} \log \sum_{\sigma} \frac{1}{2^N} \exp \beta H_N(\sigma) = \begin{cases} \frac{\beta^2}{2}, & \text{if } \beta \leq \beta_c, \\ \beta \beta_c - \log 2, & \text{if } \beta > \beta_c. \end{cases} \quad (2.9)$$

Since $F(\beta^{1/2})$ is concave in $\beta > 0$, using Theorem 2.5, one readily sees that F and V satisfy the Legendre duality with $\mathcal{B}_0 = (0, \infty)$ and $\mathcal{M}_0 = (0, 1]$,

$$V(m) = \max_{\beta > 0} \left(F(\beta) - \frac{\beta^2 m}{2} \right), \quad 0 < m \leq 1,$$

$$F(\beta) = \min_{0 < m \leq 1} \left(V(m) + \frac{\beta^2 m}{2} \right), \quad \beta > 0.$$

Note that a direct computation gives $V(m) = \log 2/m - \log 2$ for $0 < m \leq 1$. Plugging this into the variational formula for F , we can express the limiting free energy of the REM in Parisi's formulation, previously derived in [4, 12]:

Theorem 2.6 (Parisi formula for the REM). *We have that*

$$F(\beta) = \inf_{0 < m \leq 1} \left(\frac{\beta^2 m}{2} + \frac{\log 2}{m} - \log 2 \right).$$

Remarkably, the same argument also yields the Parisi formula for the REM model with random external field, see [1].

2.2.2 Ising mixed p -spin model

The Ising mixed p -spin model is defined on the same configuration space as the REM. The Hamiltonian H_N is a centered Gaussian process indexed by Σ_N and has the covariance structure,

$$\mathbb{E}H_N(\sigma^1)H_N(\sigma^2) = N\xi(R_{1,2})$$

for $\sigma^1, \sigma^2 \in \Sigma_N$, where $\xi(s) := \sum_{p \geq 2} c_p^2 s^p$ with $\sum_{p \geq 2} 2^p c_p^2 < \infty$ and $R_{1,2} := N^{-1} \sum_{1 \leq i \leq N} \sigma_i^1 \sigma_i^2$ is the overlap between σ^1 and σ^2 . Using the same definitions as (2.6), the conditions (A0), (A1) and (A2) are valid for H_N . In particular, it is famously known that the limiting free energy can be computed through Parisi's formula (see [17, 27]),

$$F(\beta) = \lim_{N \rightarrow \infty} \frac{1}{N} \log \sum_{\sigma} \frac{1}{2^N} \exp \beta H_N(\sigma) = \inf_{\alpha \in \mathcal{D}} \mathcal{P}_{\beta}(\alpha), \quad (2.10)$$

where letting \mathcal{D} be the space of all probability distribution functions on $[0, 1]$, \mathcal{P}_{β} is defined as

$$\mathcal{P}_{\beta}(\alpha) = \Phi_{\alpha, \beta}(0, 0) - \frac{\beta^2}{2} \int_0^1 \alpha(s) s \xi''(s) ds$$

for $\Phi_{\alpha, \beta}$ being the solution to the Parisi PDE with boundary condition $\Phi_{\alpha, \beta}(1, x) = \log \cosh(x)$,

$$\partial_s \Phi_{\alpha, \beta} = -\frac{\beta^2 \xi''(s)}{2} \left(\partial_{xx}^2 \Phi_{\alpha, \beta} + \alpha(s) (\partial_x \Phi_{\alpha, \beta})^2 \right), \quad (s, x) \in [0, 1] \times \mathbb{R}.$$

It is proven in [2] that the Parisi formula has a unique minimizer, denoted by $\alpha_{P, \beta}$. Set \mathcal{D}_0 the collection of all $\alpha_{P, \beta}$ associated to any temperatures $\beta > 0$. It was recently proved in [3] that $F(\beta^{1/2})$ is a concave function in $\beta > 0$. As one can compute the derivative of F from the Parisi formula (see [16]),

$$F'(\beta) = \beta \int_0^1 \alpha_{P, \beta}(s) \xi'(s) ds > 0, \quad \forall \beta > 0, \tag{2.11}$$

it follows that the duality in Theorem 2.5 holds, where $\mathcal{B}_0 = (0, \infty)$ and

$$\mathcal{M}_0 = \left\{ \int_0^1 \alpha(s) \xi'(s) ds : \alpha \in \mathcal{D}_0 \right\}.$$

Define the Legendre transform of \mathcal{P}_β by

$$\Gamma(\alpha) = \sup_{\beta > 0} \left(\mathcal{P}_\beta(\alpha) - \frac{\beta^2}{2} \int_0^1 \alpha(s) \xi'(s) ds \right), \quad \forall \alpha \in \mathcal{D}. \tag{2.12}$$

Our result below says that the limiting free energy of the squared Hamiltonian V can be identified as Γ , from which it allows to rewrite the duality obtained from Theorem 2.5 in terms of Γ . This alternative expression was firstly conjectured by Guerra [13] and established in [3].

Theorem 2.7 (Legendre structure). *Given $\alpha \in \mathcal{D}_0$, $m := \int_0^1 \alpha(s) \xi'(s) ds$, we have that*

$$V(m) = \Gamma(\alpha). \tag{2.13}$$

Moreover, for any $\alpha \in \mathcal{D}_0$ and $\beta \in \mathcal{B}_0$,

$$\Gamma(\alpha) = \max_{\beta' \in \mathcal{B}_0} \left(F(\beta') - \frac{\beta'^2}{2} \int_0^1 \alpha(s) \xi'(s) ds \right), \tag{2.14}$$

$$F(\beta) = \min_{m' \in \mathcal{M}_0} \left(\Gamma(\alpha') + \frac{\beta^2}{2} \int_0^1 \alpha'(s) \xi'(s) ds \right), \tag{2.15}$$

where the maximizer of the first variational principle is equal to β if $\alpha = \alpha_{P, \beta}$ and the minimizer of the second one is given by $\alpha_{P, \beta}$.

Proof. First we verify (2.13). Let $m = \int_0^1 \alpha(s) \xi'(s) ds$, where $\alpha = \alpha_{P, \beta}$ for some $\beta > 0$. Note that from (2.7), the characterization of the maximizer of (2.7) and the derivative of F in (2.11), we have

$$\begin{aligned} V(m) &= \max_{\beta' \in \mathcal{B}_0} \left(F(\beta') - \frac{\beta'^2}{2} \int_0^1 \alpha(s) \xi'(s) ds \right) \\ &= F(\beta) - \frac{\beta^2}{2} \int_0^1 \alpha(s) \xi'(s) ds. \end{aligned} \tag{2.16}$$

To see that this matches $\Gamma(\alpha)$, we shall need a few properties of the Parisi functional $\mathcal{P}_\beta(\alpha)$ previously established in [3, 7]. First, from [3, Equation (13)], for any $\beta' > 0$,

$$\frac{d}{d\beta'} \left(\mathcal{P}_{\beta'}(\alpha) - \frac{\beta'^2}{2} \int_0^1 \alpha(s) \xi'(s) ds \right) = -\beta' \int_0^1 \xi'(s) (\mathbb{E} u_{\alpha, \beta'}(s)^2 - s) d\alpha(s), \tag{2.17}$$

where $u_{\alpha, \beta'}$ is a uniformly bounded progressively measurable process with respect to the filtration generated by a standard Brownian motion. From [3, Equation (28)], it

is understood that for any $s \in [0, 1]$, $\mathbb{E}u_{\alpha, \beta'}(s)^2$ is nondecreasing in β' . In addition, from [7, Proposition 1], we also have that $\mathbb{E}u_{\alpha, \beta'}(s)^2 = s$ for any s in the support of α if $\beta' = \beta$. As a result, the derivative (2.17) is ≥ 0 for $\beta' < \beta$ and ≤ 0 if $\beta' > \beta$. From the definition (2.12) of $\Gamma(\alpha)$ and noting that α is the minimizer in Parisi's formula, this concludes

$$\Gamma(\alpha) = \mathcal{P}_\beta(\alpha) - \frac{\beta^2}{2} \int_0^1 \alpha(s) \xi'(s) ds,$$

which combined with (2.16) finishes the proof of (2.13).

Next, note that (2.14) follows straightforwardly from (2.13) and (2.16), while (2.15) can be obtained from (2.8) and (2.13),

$$\begin{aligned} F(\beta) &= \inf_{m \in \mathcal{M}_0} \left(V(m) + \frac{\beta^2 m}{2} \right) \\ &= \inf_{\beta' > 0} \left(\Gamma(\alpha_{P, \beta'}) + \frac{\beta^2}{2} \int_0^1 \alpha_{P, \beta'}(s) \xi'(s) ds \right) \\ &= \inf_{\alpha \in \mathcal{D}_0} \left(\Gamma(\alpha) + \frac{\beta^2}{2} \int_0^1 \alpha(s) \xi'(s) ds \right) \end{aligned}$$

where the minimizer is obtained by for $\alpha = \alpha_{P, \beta}$. This completes our proof. \square

2.2.3 Spherical mixed p -spin model

Let $\Sigma_N = \{\sigma \in \mathbb{R}^N : \sum_{i=1}^N \sigma_i^2 = N\}$ and ν_N be a uniform probability measure on Σ_N . For every $p \geq 1$, set the pure p -spin Hamiltonian as

$$H_{N,p}(\sigma) = \frac{1}{2^{p/2} N^{(p-1)/2}} \sum_{1 \leq i_1, \dots, i_p \leq N} g_{i_1, \dots, i_p} \sigma_{i_1} \cdots \sigma_{i_p},$$

where g_{i_1, \dots, i_p} 's are i.i.d. standard Gaussian for all $1 \leq i_1, \dots, i_p \leq N$ and $p \geq 1$. Throughout this subsection, we shall use the notations defined in Subsection 2.1 with this family of Hamiltonians $(H_{N,p})_{p \geq 1}$. In this case, the assumption (A1) holds by standard Gaussian concentration of measure for Lipschitz functions and Dudley's entropy integral (see for instance [8, Appendix]). Assumption (A2) holds and the limiting free energy $F(\beta)$ can be computed by a spherical version of the Parisi formula [6, 26]. Here, Parisi's formula has a simpler expression, first derived by Crisanti-Sommers [9]. It reads

$$F(\beta) = \inf_{\alpha \in \mathcal{C}} \mathcal{Q}_\beta(\alpha), \tag{2.18}$$

where \mathcal{C} is the space of all distribution functions α on $[0, 1]$ with $\alpha(\hat{q}) = 1$ for some $\hat{q} < 1$ and setting

$$\xi_\beta(q) = \sum_{p \geq 1} 2^{-p} \beta_p^2 q^p,$$

the functional \mathcal{Q}_β is defined by

$$\mathcal{Q}_\beta(\alpha) = \frac{1}{2} \left(\int_0^1 \alpha(q) \xi'_\beta(q) dq + \int_0^{\hat{q}} \frac{dq}{\int_q^1 \alpha(s) ds} + \log(1 - \hat{q}) \right).$$

Note that for any fixed α , this quantity is independent of the choice of $\hat{q} < 1$ with $\alpha(\hat{q}) = 1$. As this functional \mathcal{Q}_β is strictly convex, it ensures that the optimization problem (2.18) has a unique minimizer, denoted by $\alpha_{P, \beta}$. If now one replaces β by $\beta^{1/2}$, then $\mathcal{Q}_{\beta^{1/2}}(\alpha)$ is linear in $\beta \in \mathcal{B}$ and as a result, $F(\beta^{1/2})$ is concave on \mathcal{B} . Consequently, knowing the fact that

$$\partial_p F(\beta) = \frac{p \beta_p}{2^p} \int_0^1 \alpha_{P, \beta}(q) q^{p-1} dq > 0$$

as long as $\beta_p > 0$ (see [26]), it implies that $\mathcal{B}_0 = \mathcal{B}$ and

$$\mathcal{M}_0 = \left\{ \mathbf{m} : m_p = \frac{p}{2^p} \int_0^1 \alpha_{P,\beta}(q) q^{p-1} dq, \forall p \geq 1, \beta \in \mathcal{B} \right\}.$$

Using Theorem 2.4, we then obtain the Legendre duality between V and F . Let \mathcal{C}_0 be the collection of all $\alpha_{P,\beta}$ for some $\beta \in \mathcal{B}$. Consider the Legendre transform of \mathcal{Q}_β ,

$$\begin{aligned} \Lambda(\alpha) &:= \sup_{\beta \in \mathcal{B}} \left(\mathcal{Q}_\beta(\alpha) - \frac{1}{2} \int_0^1 \alpha(q) \xi'_\beta(q) dq \right) \\ &= \int_0^{\hat{q}} \frac{dq}{\int_q^1 \alpha(s) ds} + \log(1 - \hat{q}) \end{aligned}$$

for any $\alpha \in \mathcal{C}$. We show that the free energy associated to the squared Hamiltonians V is equal to Λ . This and the duality in Theorem 2.4 together imply that the Crisanti-Sommers formula can be regarded as the Legendre transform of V . More precisely, below is the statement of our main result in this subsection.

Theorem 2.8. *For any \mathbf{m} with*

$$m_p = \frac{p}{2^p} \int_0^1 \alpha(q) q^{p-1} dq, \forall p \geq 1,$$

for some $\alpha \in \mathcal{C}_0$, we have

$$V(\mathbf{m}) = \Lambda(\alpha). \tag{2.19}$$

Moreover, for any $\alpha \in \mathcal{C}_0$ and $\beta \in \mathcal{B}_0$,

$$\Lambda(\alpha) = \max_{\beta' \in \mathcal{B}_0} \left(F(\beta') - \frac{1}{2} \int_0^1 \alpha(s) \xi'_{\beta'}(s) ds \right), \tag{2.20}$$

$$F(\beta) = \min_{\alpha' \in \mathcal{C}_0} \left(\Lambda(\alpha') + \frac{1}{2} \int_0^1 \alpha'(s) \xi'_\beta(s) ds \right), \tag{2.21}$$

where the maximizer of the first variational principle is equal to β if $\alpha = \alpha_{P,\beta}$ and the minimizer of the second one is given by $\alpha_{P,\beta}$.

Proof. Let $\mathbf{m} \in \mathcal{M}_0$ satisfy $m_p = 2^{-p} p \int_0^1 \alpha(q) q^{p-1} dq$ for all $p \geq 1$ and some $\alpha \in \mathcal{C}_0$. Assume that $\alpha = \alpha_{P,\beta}$ for $\beta \in \mathcal{B}$. Note that $m_p = \beta_p^{-1} \partial_p F(\beta)$. From (2.4) and the characterization of the maximizer, we have that

$$\begin{aligned} V(\mathbf{m}) &= \max_{\beta' \in \mathcal{B}_0} \left(F(\beta') - \frac{1}{2} \sum_{p \geq 1} \beta_p'^2 m_p \right) \\ &= F(\beta) - \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m_p \\ &= F(\beta) - \frac{1}{2} \int_0^1 \alpha(q) \xi'_\beta(q) dq \\ &= \Lambda(\alpha), \end{aligned}$$

where the last equation used the Crisanti-Sommers formula (2.18) and the definition of $\alpha_{P,\beta}$. This gives (2.19) and evidently also (2.20), while (2.21) follows directly from (2.5) and (2.19). The characterizations of the optimizers of (2.20) and (2.21) are valid from Theorem 2.4. \square

2.3 The Parisi formula as a maximization problem

In this subsection, we return to the Ising mixed p -spin model, considered in Subsection 2.2.2 and show one application of Theorem 2.7. Recall $F(\beta)$ from (2.10) and $\Gamma(\alpha)$ from

(2.12). Consider the space \mathbf{B} of bounded functions on $[0, 1]$. For $f \in \mathbf{B}$, we define the following two real-valued functionals,

$$L_*(f) = \inf_{\alpha \in \mathcal{D}} \left(\int_0^1 \alpha(s)f(s)ds + \frac{\beta^2}{2} \int_0^1 \alpha(s)\xi'(s)ds \right)$$

and

$$\Gamma^*(f) = \sup_{\alpha \in \mathcal{D}} \left(\int_0^1 \alpha(s)f(s)ds - \Gamma(\alpha) \right).$$

The following corollary is a combination of Theorem 2.7 and convex duality.

Corollary 2.9. For any $\beta \geq 0$,

$$F(\beta) = \max_{f \in \mathbf{B}} (L_*(f) - \Gamma^*(f)). \tag{2.22}$$

The variational problem (2.22) mimicks a maximizer coming from an entropy-energy variational principle. First, only L_* depends on the temperature parameter. Second, the functional inside the variational principle defining L_* is linear and the function Γ^* is convex and lower semi-continuous. We comment and derive further properties of (2.22) at the end of Section 2.4.

Proof of Corollary 2.9. Recall that from Theorem 2.7 and the Parisi formula, we can alternatively rewrite (2.15) as

$$F(\beta) = \inf_{\alpha \in \mathcal{D}} \left(\Gamma(\alpha) + \frac{\beta^2}{2} \int_0^1 \alpha(s)\xi'(s)ds \right).$$

Since

$$\begin{aligned} \Gamma(\alpha) + \frac{\beta^2}{2} \int_0^1 \alpha(s)\xi'(s)ds &= \left(\int_0^1 \alpha(s)f(s)ds + \frac{\beta^2}{2} \int_0^1 \alpha(s)\xi'(s)ds \right) \\ &\quad + \left(\Gamma(\alpha) - \int_0^1 \alpha(s)f(s)ds \right), \end{aligned}$$

we have that for all $f \in \mathbf{B}$,

$$\Gamma(\alpha) + \frac{\beta^2}{2} \int_0^1 \alpha(s)\xi'(s)ds \geq L_*(f) - \Gamma^*(f)$$

and thus,

$$F(\beta) \geq \sup_f (L_*(f) - \Gamma^*(f)).$$

To see the reverse inequality, we choose $f_\beta(s) = \beta^2\xi'(s)/2$ for $s \in [0, 1]$. Then $L_*(f_\beta) = 0$ and obviously the Parisi formula says $\Gamma^*(f_\beta) = -F(\beta)$. From these,

$$\sup_f (L_*(f) - \Gamma^*(f)) \geq F(\beta)$$

and so (2.22) follows. □

2.4 Proofs of Theorems 2.2 and 2.4

The idea of establishing the variational formula for V in Theorem 2.2 is to approximate the squared Hamiltonians using linear approximations. One of the main difficulties here comes from the fact that the interaction involves infinitely many components. This is the reason why we need assumption (A1), which will be used to control all Hamiltonians at once. More precisely, the following lemma gives uniform bounds for $H_{N,p}$ on Σ_N for all $N, p \geq 1$ if the condition (A1) is given.

Lemma 2.10. *Given assumption (A1), there exists a sequence of positive real numbers $(b_p)_{p \geq 1}$ with $b_p < C_1 2^{-p/4}$ such that*

$$\mathbb{P}\left(\exists N_0 \geq 1 \text{ such that } \sup_{\sigma \in \Sigma_N} \frac{|H_{N,p}(\sigma)|}{N} < b_p, \forall p \geq 1, N \geq N_0\right) = 1. \quad (2.23)$$

Proof. From (A1), we have that

$$\mathbb{E} \sup_{\sigma \in \Sigma_N} \frac{|H_{N,p}(\sigma)|}{N} \leq \frac{C_1}{2^{p/2}}$$

and that with probability at most $C_2 e^{-2^p N t^2 / C_2}$, the following event holds,

$$\sup_{\sigma \in \Sigma_N} \frac{|H_{N,p}(\sigma)|}{N} \geq \mathbb{E} \sup_{\sigma \in \Sigma_N} \frac{|H_{N,p}(\sigma)|}{N} + t.$$

Combining these together, with probability at most $C_2 e^{-2^p N t^2 / C_2}$,

$$\sup_{\sigma \in \Sigma_N} \frac{|H_{N,p}(\sigma)|}{N} \geq t - \frac{C_1}{2^{p/2}}.$$

In particular, taking $t = 2^{-p/4} C_1$, we obtain that with probability at most $C_2 e^{-2^{p/2} N C_1^2 / C_2}$,

$$\sup_{\sigma \in \Sigma_N} \frac{|H_{N,p}(\sigma)|}{N} \geq b_p := \frac{C_1}{2^{p/4}} \left(1 - \frac{1}{2^{p/4}}\right).$$

Denote by Ω_N the event that there exists some $p \geq 1$ such that

$$\sup_{\sigma \in \Sigma_N} \frac{|H_{N,p}(\sigma)|}{N} \geq b_p.$$

Since

$$\begin{aligned} \sum_{N \geq 1} \mathbb{P}(\Omega_N) &\leq \sum_{N \geq 1} \sum_{p \geq 1} \mathbb{P}\left(\sup_{\sigma \in \Sigma_N} \frac{|H_{N,p}(\sigma)|}{N} \geq b_p\right) \\ &\leq C_2 \sum_{N \geq 1} \sum_{p \geq 1} e^{-2^{p/2} N C_1^2 / C_2} < \infty, \end{aligned}$$

the Borel-Cantelli lemma says that the event that Ω_N occurs infinitely often has zero probability. This gives the announced result. \square

For each $k \geq 1$, set

$$\begin{aligned} \mathcal{B}_k &= \{\beta_k = (\beta_p)_{1 \leq p \leq k} : \beta_p > 0, \forall 1 \leq p \leq k\}, \\ \mathcal{M}_k &= \{\mathbf{m}_k = (m_p)_{1 \leq p \leq k} : m_p > 0, \forall 1 \leq p \leq k\}. \end{aligned}$$

Consider the following free energies,

$$\begin{aligned} F_{N,k}(\beta_k) &= \frac{1}{N} \log \int_{\Sigma_N} \exp \sum_{p=1}^k \beta_p H_{N,p}(\sigma) \nu_N(d\sigma), \quad \forall \beta_k \in \mathcal{B}_k, \\ V_{N,k}(\mathbf{m}_k) &= \frac{1}{N} \log \int_{\Sigma_N} \frac{1}{2^N} \exp \sum_{p=1}^k \frac{N}{2m_p} \left(\frac{H_{N,p}(\sigma)}{N}\right)^2 \nu_N(d\sigma), \quad \forall \mathbf{m}_k \in \mathcal{M}_k. \end{aligned}$$

We use F_k and V_k to denote the limits of $(F_{N,k})_{N \geq 1}$ and $(V_{N,k})_{N \geq 1}$ whenever they exist in $\overline{\mathbb{R}}$.

Lemma 2.11. *Assume that (A0), (A1) and (A2) hold. Consider any $I \subseteq \mathbb{N}$ and $(\beta_p)_{p \notin I}$ with $\sum_{p \notin I} \beta_p^2 < \infty$. For any $\gamma = (\gamma_p)_{p \in I}$ with $\sum_{p \in I} \gamma_p^2 < \infty$, define the restriction of F to I by*

$$g(\gamma) = F(\phi(\gamma)),$$

where $\phi(\gamma) := (\delta_p)_{p \geq 1}$ satisfies $\delta_p = \gamma_p$ if $p \in I$ and $\delta_p = \beta_p$ if $p \notin I$. Then g is continuous at $\mathbf{0} := (0, 0, \dots)$.

Proof. From the convexity of F_N on the set of coordinates I ,

$$-\sum_{p \in I} \partial_p F_N(\phi(\gamma)) \gamma_p \leq F_N(\phi(\gamma)) - F_{N,k}(\phi(\mathbf{0})) \leq \sum_{p \in I} \partial_p F_N(\phi(\mathbf{0})) \gamma_p \tag{2.24}$$

for any $\gamma = (\gamma_p)_{p \in I}$ with $\sum_{p \in I} \gamma_p^2 < \infty$. To control the two sides, observe that

$$\inf_{\sigma \in \Sigma_N} \frac{H_{N,p}(\sigma)}{N} \leq \partial_p F_N(\phi(\gamma)) = \left\langle \frac{H_{N,p}(\sigma)}{N} \right\rangle_\gamma \leq \sup_{\sigma \in \Sigma_N} \frac{H_{N,p}(\sigma)}{N}, \quad \forall p \in I, \tag{2.25}$$

where $\langle \cdot \rangle_\gamma$ is the expectation with respect to the Gibbs measure,

$$G_\gamma(d\sigma) = \frac{\exp\left(\sum_{p \notin I} \beta_p H_{N,p}(\sigma) + \sum_{p \in I} \gamma_p H_{N,p}(\sigma)\right) \nu_N(d\sigma)}{\int_{\Sigma_N} \exp\left(\sum_{p \notin I} \beta_p H_{N,p}(\sigma') + \sum_{p \in I} \gamma_p H_{N,p}(\sigma')\right) \nu_N(d\sigma')}.$$

From the condition (A0), the random variables $-\inf_{\sigma \in \Sigma_N} H_{N,p}(\sigma)$ and $\sup_{\sigma \in \Sigma_N} H_{N,p}(\sigma)$ have the same distribution. Taking expectations in (2.24) and applying (2.25) lead to

$$|\mathbb{E}F_N(\phi(\gamma)) - \mathbb{E}F_N(\phi(\mathbf{0}))| \leq C_1 \sum_{p \in I} \frac{\gamma_p}{2^{p/2}},$$

where the inequality used (A1). Note that the condition (A2) implies $\lim_{N \rightarrow \infty} \mathbb{E}F_N(\phi(\gamma)) = F(\phi(\gamma))$ by the dominated convergence theorem. Hence, passing to the limit in the last inequality gives

$$|g(\gamma) - g(\mathbf{0})| \leq C_1 \sum_{p \in I} \frac{\gamma_p}{2^{p/2}}.$$

The continuity of F at $\mathbf{0}$ follows evidently. □

The next lemma is a finite dimensional version of (2.2).

Lemma 2.12. *Suppose that (A0), (A1) and (A2) hold. For any $\mathbf{m}_k \in \mathcal{M}_k$, we have*

$$V_k(\mathbf{m}_k) = \sup_{\beta_k \in \mathcal{B}_k} \left(F_k(\beta_k) - \frac{1}{2} \sum_{p=1}^k \beta_p^2 m_p \right).$$

Proof. From Lemma 2.10, there exists a sequence of positive real numbers $(b_p)_{p \geq 1}$ such that with probability one, there exists some $N_0 > 0$ such that

$$\frac{H_{N,p}(\sigma)}{N} \in [-b_p, b_p], \quad \forall \sigma \in \Sigma_N \text{ and } \forall p \geq 1, N \geq N_0,$$

Since we only care about the limiting behavior of $(V_{k,N})_{N \geq 1}$, we assume without loss of generality that $|H_{N,p}/N| \leq b_p$ for all $1 \leq p \leq k$ and $N \geq 1$. For any $\varepsilon > 0$, let $\ell_p \in \mathbb{N}$ with $\ell_p \geq b_p/\varepsilon$. Set $\beta_{p,i} = i\varepsilon/m_p$ for $i \in [-\ell_p, \ell_p] \cap \mathbb{Z}$. For any $i_p \in [-\ell_p, \ell_p] \cap \mathbb{Z}$ for $1 \leq p \leq k$, define

$$A_{p,i_1, \dots, i_k} = \left\{ \sigma \in \Sigma_N \mid \frac{H_{N,p}(\sigma)}{N} \in [(i_p - 1)\varepsilon, (i_p + 1)\varepsilon], \forall 1 \leq p \leq k \right\}.$$

Observe that for any $\sigma \in A_{p,i_1,\dots,i_k}$, since $\beta_{p,i_p} m_p = i_p \varepsilon$, we have

$$\begin{aligned} \sum_{p=1}^k \frac{N}{2m_p} \left(\frac{H_{N,p}(\sigma)}{N} \right)^2 &= \sum_{p=1}^k \frac{N}{2m_p} \left(\frac{H_{N,p}(\sigma)}{N} - \beta_{p,i_p} m_p + \beta_{p,i_p} m_p \right)^2 \\ &= \sum_{p=1}^k \frac{N}{2m_p} \left(\frac{H_{N,p}(\sigma)}{N} - \beta_{p,i_p} m_p \right)^2 \\ &\quad + \sum_{p=1}^k N \beta_{p,i_p} \left(\frac{H_{N,p}(\sigma)}{N} - \beta_{p,i_p} m_p \right) + \sum_{p=1}^k \frac{N \beta_{p,i_p}^2 m_p}{2} \\ &\leq N \varepsilon^2 \sum_{p=1}^k \frac{1}{2m_p} + \sum_{p=1}^k \beta_{p,i_p} H_{N,p}(\sigma) - N \sum_{p=1}^k \frac{\beta_{p,i_p}^2 m_p}{2}, \end{aligned}$$

from which

$$\begin{aligned} &\int_{\Sigma_N} \exp \sum_{p=1}^k \frac{N}{2m_p} \left(\frac{H_{N,p}(\sigma)}{N} \right)^2 \nu_N(d\sigma) \\ &\leq \sum_{|i_1| \leq \ell_1, \dots, |i_k| \leq \ell_k} \exp \left(-N \sum_{p=1}^k \frac{\beta_{p,i_p}^2 m_p}{2} + N \varepsilon^2 \sum_{p=1}^k \frac{1}{2m_p} \right) \\ &\quad \cdot \int_{\Sigma_N} 1_{A_{p,i_1,\dots,i_k}}(\sigma) \exp \left(\sum_{p=1}^k \beta_{p,i_p} H_{N,p}(\sigma) \right) \nu_N(d\sigma) \\ &\leq \sum_{|i_1| \leq \ell_1, \dots, |i_k| \leq \ell_k} \exp \left(-N \sum_{p=1}^k \frac{\beta_{p,i_p}^2 m_p}{2} + N \varepsilon^2 \sum_{p=1}^k \frac{1}{2m_p} \right) \\ &\quad \cdot \int_{\Sigma_N} \exp \left(\sum_{p=1}^k \beta_{p,i_p} H_{N,p}(\sigma) \right) \nu_N(d\sigma). \end{aligned}$$

Consequently, taking $N^{-1} \log$ on both sides leads to

$$\begin{aligned} V_{N,k}(\mathbf{m}_k) &\leq \frac{\log \prod_{p=1}^k (2\ell_p + 1)}{N} + \varepsilon^2 \sum_{p=1}^k \frac{1}{2m_p} \\ &\quad + \max_{-\ell_p \leq i_p \leq \ell_p, \forall 1 \leq p \leq k} \left(F_{N,k}(\beta_{k,i_1}, \dots, \beta_{k,i_p}) - \frac{1}{2} \sum_{p=1}^k \beta_{p,i}^2 m_p \right) \end{aligned}$$

and from (A2), passing to limit gives

$$\begin{aligned} \limsup_{N \rightarrow \infty} V_{N,k}(\mathbf{m}_k) &\leq \varepsilon^2 \sum_{p=1}^k \frac{1}{m_p} + \max_{-\ell_p \leq i_p \leq \ell_p, \forall 1 \leq p \leq k} \left(F_k(\beta_{k,i_1}, \dots, \beta_{k,i_p}) - \frac{1}{2} \sum_{p=1}^k \beta_{p,i}^2 m_p \right) \\ &\leq \varepsilon^2 \sum_{p=1}^k \frac{1}{m_p} + \sup_{\beta_k \in \mathcal{B}_k} \left(F_k(\beta_k) - \frac{1}{2} \sum_{p=1}^k \beta_p^2 m_p \right), \end{aligned}$$

where the last inequality used the fact that F_k is an even function in each coordinate and F_k is continuous at $\mathbf{0}_k = (0, \dots, 0)$ by applying Lemma 2.11. Since this holds for any ε , it follows

$$\limsup_{N \rightarrow \infty} V_{N,k}(\mathbf{m}_k) \leq \sup_{\beta_k \in \mathcal{B}_k} \left(F_k(\beta_k) - \frac{1}{2} \sum_{p=1}^k \beta_p^2 m_p \right).$$

To prove the reverse direction, note that for any $\beta_k \in \mathcal{B}_k$,

$$\beta_p H_{N,p}(\sigma) = \beta_p (Nm_p)^{1/2} \cdot \frac{H_{N,p}(\sigma)}{(Nm_p)^{1/2}} \leq \frac{\beta_p^2 Nm_p}{2} + \frac{H_{N,p}(\sigma)^2}{2Nm_p}, \quad (2.26)$$

which implies

$$V_{N,k}(\mathbf{m}_k) \geq F_{N,k}(\beta_k) - \frac{1}{2} \sum_{p=1}^k \beta_p^2 m_p, \quad \forall \beta_k \in \mathcal{B}_k$$

and thus,

$$\liminf_{N \rightarrow \infty} V_{N,k}(\mathbf{m}_k) \geq \sup_{\beta_k \in \mathcal{B}_k} \left(F_k(\beta_k) - \frac{1}{2} \sum_{p=1}^k \beta_p^2 m_p \right).$$

This finishes our proof. □

As we have seen from Lemma 2.12, the variational representation of V is based on the assumptions (A0), (A1) and (A2). In what follows, we continue to provide the proof of Theorem 2.2. As we will see below, the same conditions in Lemma 2.12 also yield (2.2). However, the validity of the formula (2.3) for F will require concavity of the limiting free energy.

Proof of Theorem 2.2. Let $\mathbf{m} \in \mathcal{M}$ and let $\bar{\mathbf{m}}_k$ be the projection of the first k -coordinates of \mathbf{m} . Note that from Lemmas 2.11 and 2.12,

$$V_k(\bar{\mathbf{m}}_k) = \sup_{\beta_k \in \mathcal{B}_k} \left(F_k(\beta_k) - \frac{1}{2} \sum_{p=1}^k \beta_p^2 m_p \right) \leq \sup_{\beta \in \mathcal{B}} \left(F(\beta) - \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m_p \right).$$

Since $V_{N,k}(\bar{\mathbf{m}}_k) \leq V_{N,k+1}(\bar{\mathbf{m}}_{k+1})$ for all $k \geq 1$, we may use the monotone convergence theorem to get

$$V_N(\mathbf{m}) \leq \sup_{\beta \in \mathcal{B}} \left(F(\beta) - \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m_p \right)$$

and passing to limit,

$$\limsup_{N \rightarrow \infty} V_N(\mathbf{m}) \leq \sup_{\beta \in \mathcal{B}} \left(F(\beta) - \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m_p \right).$$

On the other hand, from (2.26), we could also obtain that for any $\beta \in \mathcal{B}$,

$$V_N(\mathbf{m}) \geq F_N(\beta) - \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m_p,$$

from which letting $N \rightarrow \infty$ and then taking supremum,

$$\liminf_{N \rightarrow \infty} V_N(\mathbf{m}) \geq \sup_{\beta \in \mathcal{B}} \left(F(\beta) - \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m_p \right).$$

This finishes the proof of (i).

As for (ii), we assume that F is Gâteaux differentiable on \mathcal{B} and $F(\beta^{1/2})$ is concave on \mathcal{B} . Let $\beta \in \mathcal{B}$ and set $C(\beta) = \{p \geq 1 : \partial_p F(\beta) = 0\}$. Take any $\mathbf{m}' \in \mathcal{M}$ with $m'_p = \beta_p^{-1} \partial_p F(\beta)$ for all $p \notin C(\beta)$. From (i), we rewrite

$$V(\mathbf{m}') = \sup_{\gamma \in \mathcal{B}} L(\gamma^2) \quad (2.27)$$

for

$$L(\mathbf{t}) := F(\mathbf{t}^{1/2}) - \frac{1}{2} \sum_{p \geq 1} t_p m'_p.$$

Since $F(\beta^{1/2})$ is concave in $\beta \in \mathcal{B}$, the function L is concave as well. As now,

$$\partial_p L(\mathbf{t}) = \frac{1}{2} \left(\frac{\partial_p F(\mathbf{t}^{1/2})}{t_p^{1/2}} - m'_p \right),$$

if we let $\mathbf{t} = \beta^2$, from our choice of \mathbf{m}' , this partial derivative is equal to 0 for all $p \notin C(\beta)$ and moreover, it is equal to $-m'_p/2$ for all $p \in C(\beta)$. It then follows from the concavity of L that

$$\begin{aligned} L(\gamma^2) - L(\beta^2) &\leq \sum_{p \geq 1} \partial_p L(\beta^2) (\gamma_p^2 - \beta_p^2) \\ &= -\frac{1}{2} \sum_{p \in C(\beta)} m'_p (\gamma_p^2 - \beta_p^2) \\ &\leq \frac{1}{2} \sum_{p \in C(\beta)} m'_p \beta_p^2, \quad \forall \gamma \in \mathcal{B}, \end{aligned}$$

from which

$$V(\mathbf{m}') \leq F(\beta) - \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m'_p + \frac{1}{2} \sum_{p \in C(\beta)} m'_p \beta_p^2 \tag{2.28}$$

and thus,

$$\begin{aligned} \inf_{\mathbf{m} \in \mathcal{M}} \left(V(\mathbf{m}) + \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m_p \right) &\leq V(\mathbf{m}') + \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m'_p \\ &\leq F(\beta) + \frac{1}{2} \sum_{p \in C(\beta)} \beta_p^2 m'_p. \end{aligned}$$

Since this inequality is valid for any $\mathbf{m}' \in \mathcal{M}$ with $m'_p = \beta_p^{-1} \partial_p F(\beta)$ for all $p \notin C(\beta)$. Letting $m'_p \downarrow 0$ for all $p \in C(\beta)$, we obtain

$$\inf_{\mathbf{m} \in \mathcal{M}} \left(V(\mathbf{m}) + \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m_p \right) \leq F(\beta).$$

Finally, from (2.26), it can be easily seen that

$$F(\beta) \leq \inf_{\mathbf{m} \in \mathcal{M}} \left(V(\mathbf{m}) + \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m_p \right).$$

Our proof is completed by this and last inequalities. □

Proof of Theorem 2.4. In view of the second-half of the proof of Theorem 2.2 and the notation $C(\beta)$ therein, if $\beta \in \mathcal{B}_0$ and $\mathbf{m} \in \mathcal{M}_0$ satisfy $m_p = \beta_p^{-1} \partial_p F(\beta)$ for all $p \geq 1$, then $C(\beta) = \emptyset$ and from (2.2) and (2.28),

$$V(\mathbf{m}) = F(\beta) - \frac{1}{2} \sum_{p \geq 1} \beta_p^2 m_p.$$

This clearly gives (i) and (ii) by using (2.27) and (2.28). □

References

- [1] Arguin, L.-P., Kistler, N: Microcanonical analysis of the random energy model in a random magnetic field. *J. Stat. Phys.*, **157**, no. 1, 1–16 (2014) [MR-3249901](#)
- [2] Auffinger, A., Chen, W.-K.: The Parisi formula has a unique minimizer. *Comm. Math. Phys.*, **335**, no. 3, 1429–1444 (2015) [MR-3320318](#)
- [3] Auffinger, A., Chen, W.-K.: The Legendre structure of the Parisi formula. To appear in *Comm. Math. Phys.* (2015) [MR-3555353](#)
- [4] Bolthausen, E., Kistler, N.: Universal structures in some mean field spin glasses and an application. *J. Math. Phys.*, **49**, 25–205 (2008) [MR-2484336](#)
- [5] Bolthausen, E., Kistler, N.: A quenched large deviation principle and a Parisi formula for a Perceptron version of the GREM. *Probability in Complex Physical Systems*, Springer Proceedings in Mathematics, **11**, 425–442 (2012) [MR-3372858](#)
- [6] Chen, W.-K.: The Aizenman-Sims-Starr scheme and Parisi formula for mixed p-spin spherical models. *Elec. Journal Probab.*, **18**, no. 94, 1–14 (2013) [MR-3126577](#)
- [7] Chen, W.-K.: Variational representation for the Parisi functional and the two-dimensional Guerra-Talagrand bound. [arXiv:1501.06635](#) (2015)
- [8] Chen, W.-K., Sen, A.: Parisi formula, disorder chaos and fluctuation for the ground state energy in the spherical mixed p-spin models. [arXiv:1512.08492](#) (2015) [MR-3606472](#)
- [9] Crisanti, A., Sommers, H.-J.: The spherical p-spin interaction spin glass model: the statics. *Zeitschrift für Physik B Condensed Matter*, **87**, 341–354 (1992)
- [10] Dembo, A., Zeitouni, O.: *Large Deviations Techniques and Applications*. 2nd ed., Springer-Verlag, New York (2009) [MR-1619036](#)
- [11] Ellis, R.: *Entropy, Large Deviations, and Statistical Mechanics*. 1st ed., Springer-Verlag, Berlin (2006) [MR-2189669](#)
- [12] Guerra, F.: The phenomenon of spontaneous replica symmetry breaking in complex statistical mechanics systems. *J. Phys.: Conf. Ser.*, **442**, 012013 (2013)
- [13] Guerra, F.: Legendre structures in statistical mechanics for ordered and disordered systems, Cambridge University Press, in *Advances in disordered systems, random processes and some applications*, Contucci, P. et al, eds, in print (2016)
- [14] Mézard, M., Parisi, G., Virasoro, M.: *Spin glass theory and beyond*. World Scientific, **9**, Singapore (2004) [MR-1026102](#)
- [15] Panchenko, D.: Free energy in the generalized Sherrington-Kirkpatrick mean field model. *Rev. Math. Phys.*, **17**, no. 7, 793–857 (2005) [MR-2159369](#)
- [16] Panchenko, D.: On differentiability of the Parisi formula. *Electron. Commun. Probab.*, **13**, 241–247 (2008) [MR-2399285](#)
- [17] Panchenko, D.: *The Sherrington-Kirkpatrick model*. Springer Monographs in Mathematics. Springer, New York (2013) [MR-3052333](#)
- [18] Panchenko, D.: The Parisi formula for mixed p-spin models. *Ann. Probab.*, **42**, no. 3, 946–958 (2014) [MR-3189062](#)
- [19] Panchenko, D.: The free energy in a multi-species Sherrington-Kirkpatrick model. *Ann. of Prob.*, **43**, no. 6, 3494–3513 (2015) [MR-3433586](#)
- [20] Panchenko, D.: Free energy in the Potts spin glass. [arXiv:1512.00370](#) (2015)
- [21] Panchenko, D.: Free energy in the mixed p-spin models with vector spins. [arXiv:1512.04441](#) (2015)
- [22] Parisi, G.: Infinite number of order parameters for spin-glasses. *Phys. Rev. Lett.*, **43**, 1754–1756 (1979) [MR-0702601](#)
- [23] Rassoul-ahga, F., Seppäläinen, T.: *A course on large deviations with an introduction to Gibbs measures*. Graduate Studies in Mathematics, American Mathematical Society (2015) [MR-3309619](#)
- [24] Talagrand, M.: *Spin Glasses: A Challenge for Mathematicians: Cavity and Mean Field Models*. *Ergebnisse der Mathematik und ihrer Grenzgebiete. 3. Folge. A Series of Modern Surveys in Mathematics*, **46**, Springer-Verlag, Berlin (2003)

A duality principle in spin glasses

- [25] Talagrand, M.: The Parisi formula. *Ann. of Math. (2)*, **163**, no. 1, 221–263 (2006) [MR-2195134](#)
- [26] Talagrand, M.: The free energy of the spherical mean-field model. *Probab. Theory Relat. Fields*, **134**, 339–382 (2006) [MR-2226885](#)
- [27] Talagrand, M.: Mean field models for spin glasses. *Ergebnisse der Mathematik und ihrer Grenzgebiete. 3. Folge. A Series of Modern Surveys in Mathematics*, **55**, Springer-Verlag, Berlin (2011) [MR-1993891](#)

Electronic Journal of Probability

Electronic Communications in Probability

Advantages of publishing in EJP-ECP

- Very high standards
- Free for authors, free for readers
- Quick publication (no backlog)
- Secure publication (LOCKSS¹)
- Easy interface (EJMS²)

Economical model of EJP-ECP

- Non profit, sponsored by IMS³, BS⁴, ProjectEuclid⁵
- Purely electronic

Help keep the journal free and vigorous

- Donate to the IMS open access fund⁶ (click here to donate!)
- Submit your best articles to EJP-ECP
- Choose EJP-ECP over for-profit journals

¹LOCKSS: Lots of Copies Keep Stuff Safe <http://www.lockss.org/>

²EJMS: Electronic Journal Management System <http://www.vtex.lt/en/ejms.html>

³IMS: Institute of Mathematical Statistics <http://www.imstat.org/>

⁴BS: Bernoulli Society <http://www.bernoulli-society.org/>

⁵Project Euclid: <https://projecteuclid.org/>

⁶IMS Open Access Fund: <http://www.imstat.org/publications/open.htm>