# ON THE LOCATION OF THE PEAKS OF LEAST-ENERGY SOLUTIONS TO SEMILINEAR DIRICHLET PROBLEMS WITH CRITICAL GROWTH 

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We study the location of the peaks of solution for the critical growth problem $-\varepsilon^{2} \Delta u+u=f(u)+u^{2^{*}-1}, u>0$ in $\Omega, u=0$ on $\partial \Omega$, where $\Omega$ is a bounded domain; $2^{*}=2 N /(N-2), N \geq 3$, is the critical Sobolev exponent and $f$ has a behavior like $u^{p}, 1<p<2^{*}-1$.

## 1. Introduction

In this paper, we will study the location of the peaks of least-energy solution for the problem

$$
\begin{align*}
-\varepsilon^{2} \Delta u+u & =f(u)+u^{2^{2}-1} \quad \text { in } \Omega, \\
u & >0  \tag{1.1}\\
u & \text { in } \Omega, \\
u & \text { on } \partial \Omega,
\end{align*}
$$

where $\Omega$ is a bounded domain in $\mathbb{R}^{N}, \varepsilon>0$, and $f$ is a function satisfying some subcritical conditions. Here $2^{*}=2 N /(N-2), N \geq 3$, is the critical Sobolev exponent.

By least-energy solution for problem (1.1) we mean a critical point at the Mountain-Pass level of the associated energy functional

$$
\begin{equation*}
J_{\varepsilon}(u)=\frac{1}{2} \int_{\Omega}\left(\varepsilon^{2}|\nabla u|^{2}+u^{2}\right) d z-\int_{\Omega}\left[F(u)+\frac{1}{2^{*}}\left(u^{+}\right)^{2^{*}}\right] d z, \tag{1.2}
\end{equation*}
$$

(where $u^{+}=\max \{u, 0\}$ ), defined on the Hilbert space $H_{o}^{1}(\Omega)$ endowed with the norm

$$
\begin{equation*}
\|u\|_{\varepsilon}^{2}=\int_{\Omega}\left(\varepsilon^{2}|\nabla u|^{2}+u^{2}\right) d z . \tag{1.3}
\end{equation*}
$$

The Mountain-Pass level of $J_{\varepsilon}$ is defined by

$$
\begin{equation*}
c_{\varepsilon}=\inf _{g \in \Gamma} \max _{0 \leq t \leq 1} J_{\varepsilon}(g(t)) \tag{1.4}
\end{equation*}
$$

where $\Gamma$ is the set of all continuous paths joining the origin and a fixed nonzero element $e$ in $H_{o}^{1}(\Omega)$, such that $e \neq 0$ and $J_{\varepsilon}(e) \leq 0$. Under suitable hypothesis (e.g., $\left(f_{1}\right),\left(f_{4}\right),\left(f_{5}\right)$ below), it is not hard to check that $c_{\varepsilon}>0$ does not depend on the element $0 \neq v \in H_{o}^{1}(\Omega)$ and $u$ is a least-energy solution if and only if $J_{\varepsilon}(u)=c$ and $J_{\varepsilon}^{\prime}(u)=0$, and $J_{\varepsilon}(u) \leq J_{\varepsilon}(v)$ for all $v \neq 0$ such that $J_{\varepsilon}^{\prime}(v)=0$.

The existence of least-energy solution of problem (1.1) was given in Brézis and Nirenberg in [3, Theorem 2.1] (see Lemma 2.4 in this paper).

In this paper, we will study some properties of the least-energy solution $u_{\varepsilon}$ of problem (1.1) when $\varepsilon$ is small. In order to describe these properties, we introduce the hypotheses on the function $f$.

Suppose that $f: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}$is a $C^{1, \alpha}$ function such that
$\left(\mathrm{f}_{1}\right) f(0)=f^{\prime}(0)=0$;
( $\mathrm{f}_{2}$ ) there is $q_{1} \in(1,(N+2) /(N-2))$ such that

$$
\begin{equation*}
\lim _{s \rightarrow \infty} \frac{f(s)}{s^{q_{1}}}=0 \tag{1.5}
\end{equation*}
$$

$\left(\mathrm{f}_{3}\right)$ there are $q_{2} \in(1,(N+2) /(N-2))$ and $\lambda>0$ such that

$$
\begin{equation*}
f(s) \geq \lambda s^{q_{2}}, \quad \forall s>0 \tag{1.6}
\end{equation*}
$$

(when $N=3$, we need $q_{2}>2$, otherwise we require a sufficiently large $\lambda$ );
$\left(\mathrm{f}_{4}\right)$ if $F(s)=\int_{o}^{s} f(t) d t$, for some $\theta \in\left(2, q_{1}+1\right)$ we have

$$
\begin{equation*}
0<\theta F(s) \leq f(s) s, \quad \forall s>0 \tag{1.7}
\end{equation*}
$$

$\left(\mathrm{f}_{5}\right)$ the function $f(s) / s$ is increasing for $s>0$.
Since our interest is on positive solutions we define $f(s)=0$, in $s \leq 0$.
Now we will state our main result.
Theorem 1.1. Suppose that $\Omega$ is a bounded domain in $\mathbb{R}^{N} ; f$ satisfies $\left(f_{1}\right),\left(f_{2}\right)$, ( $\mathrm{f}_{3}$ ), $\left(\mathrm{f}_{4}\right),\left(\mathrm{f}_{5}\right)$; and let $u_{\varepsilon}$ be the least-energy solution of (1.1). Then, there is a $\varepsilon_{o}>0$ such that
(i) $u_{\varepsilon}$ attains only one local maximum at some $z_{\varepsilon} \in \Omega$ (hence global maximum $)$, for all $\varepsilon \in\left(0, \varepsilon_{o}\right]$;
(ii) $u_{\varepsilon}$ converges uniformly to zero over compact subsets of $\Omega \backslash\left\{z_{\varepsilon}\right\}$ as $\varepsilon \rightarrow 0$;
(iii) $\operatorname{dist}\left(z_{\mathcal{E}}, \partial \Omega\right) \rightarrow \max _{z \in \Omega} \operatorname{dist}(z, \partial \Omega)$.

This statement is analogous to the one given by Ni and Wei in [8], in the subcritical case

$$
\begin{gather*}
-\varepsilon^{2} \Delta u+u=h(u), \quad \text { in } \Omega, \\
u=0, \quad \text { on } \partial \Omega, \tag{1.8}
\end{gather*}
$$

where $h$ satisfies the following hypothesis:
(i) $\left(f_{1}\right),\left(f_{2}\right),\left(f_{4}\right)$, and $\left(f_{5}\right)$ hold;
(ii) the global problem

$$
\begin{equation*}
-\Delta u+u=h(u), \quad \text { in } \mathbb{R}^{N} \tag{1.9}
\end{equation*}
$$

has a unique positive solution in $H^{1}\left(\mathbb{R}^{N}\right)$;
(iii) this solution is nondegenerate in the sense that

$$
\begin{equation*}
-\Delta v+v=h^{\prime}(u) v, \quad \text { in } \mathbb{R}^{N} \tag{1.10}
\end{equation*}
$$

has no nontrivial spherically symmetric solution in $L^{2}\left(\mathbb{R}^{N}\right)$.
In [8], Ni and Wei also have described the asymptotic profile (in $\varepsilon$ ) of $u_{\varepsilon}$, giving a detailed description for $\varepsilon$ small. Here in the critical case, the solutions have the same profile.

In this work we will show that a ground state solution of the critical problem (1.1) is also solution of a subcritical problem (1.8) by showing that for small $\varepsilon$ we have a uniform bound for the $L^{\infty}$ norm of $u_{\varepsilon}$.

The difficulty here lies in finding an upper bound for $\left\|u_{\mathcal{\varepsilon}}\right\|_{L^{\infty}(\Omega)}$ by obtaining a bound for $u_{\varepsilon}$ in $L^{p}(\Omega)$ norm, for all $p \geq 2$. In the subcritical case this boundedness is obtained since the family $u_{\varepsilon}$ is bounded in $H^{1}(\Omega)$ but this argument does not work in the critical case. Here, we obtain an $L^{\infty}$-bound for $u_{\varepsilon}$ through the estimate below, which is based on Moser's iteration technique (see [11]) and is essentially due to Brézis and Kato [2].

Proposition 1.2. Let $\Lambda$ be an open subset and $q \in L^{N / 2}(\Lambda)$. Suppose that $g: \Lambda \times$ $\mathbb{R} \rightarrow \mathbb{R}$ is a Caratheodory function satisfying

$$
\begin{equation*}
|g(x, s)| \leq\left(q(x)+C_{g}\right)|s|, \quad \forall s \in \mathbb{R}, x \in \Lambda \text { and for some } C_{g}>0 \tag{1.11}
\end{equation*}
$$

Then, if $v \in H_{o}^{1}(\Lambda)$ is such that

$$
\begin{equation*}
-\Delta v=g(x, v), \quad \text { in } \Lambda \tag{1.12}
\end{equation*}
$$

we have $v \in L^{p}(\Lambda)$ for all $2 \leq p<\infty$. Moreover, there is a positive constant $C_{p}=$ $C\left(p, C_{g}, q\right)$ such that

$$
\begin{equation*}
\|v\|_{L^{2^{*}(p+1)}(\Lambda)} \leq C_{p}\|v\|_{L^{2(p+1)}(\Lambda)} \tag{1.13}
\end{equation*}
$$

Remark 1.3. The dependence on $q$ of $C_{p}$ can be given uniformly on a family of functions $\left\{q_{\varepsilon}\right\}_{\varepsilon>0}$ such that $q_{\varepsilon}$ converges in $L^{N / 2}$ (see the appendix).

We have organized this paper as follows: the next section contains the proof of Theorem 1.1. This proof consists in a series of lemmas which show the $L^{\infty}$ bound for $u_{\mathcal{\varepsilon}}$, where these functions are solutions of a class of subcritical problems (1.8). The third section is an appendix proving Proposition 1.2, for the sake of completeness.

## 2. Proof of Theorem 1.1

Before proving Theorem 1.1, let us fix some notation and preliminaries.
Remark 2.1. Throughout this section, we use the equivalent characterization of $c_{\varepsilon}$, which is more adequate to our purposes, given by

$$
\begin{equation*}
c_{\varepsilon}=\inf _{v \in H_{o}^{1}(\Omega) \backslash\{0\}} \max _{t \geq 0} J_{\varepsilon}(t v) . \tag{2.1}
\end{equation*}
$$

(see Willem [13, Theorem 4.2]).
We denote by $J: H^{1}\left(\mathbb{R}^{N}\right) \rightarrow \mathbb{R}$ the functional given by

$$
\begin{equation*}
J(u)=\frac{1}{2}\|u\|^{2}-\int_{\mathbb{R}^{N}}\left[F(u)+\frac{1}{2^{*}}\left(u_{+}\right)^{2^{*}}\right] d x \tag{2.2}
\end{equation*}
$$

where

$$
\begin{equation*}
\|u\|^{2}=\int_{\mathbb{R}^{N}}\left(|\nabla u|^{2}+u^{2}\right) d x \tag{2.3}
\end{equation*}
$$

associated with the problem

$$
\begin{equation*}
-\Delta u+u=f(u)+|u|^{2^{*}-2} u, \quad \text { in } \mathbb{R}^{N} \tag{2.4}
\end{equation*}
$$

It is known that under assumptions $\left(f_{1}\right),\left(f_{2}\right),\left(f_{3}\right),\left(f_{4}\right),\left(f_{5}\right)$, and (2.4) possesses a ground state solution $\omega$ in the level

$$
\begin{equation*}
c=J(\omega)=\inf _{v \in H^{1}\left(\mathbb{R}^{N}\right) \backslash\{0\}} \max _{t \geq 0} J(t v), \tag{2.5}
\end{equation*}
$$

(see [1]).
Remark 2.2. It is easy to check that for each nonzero $v$ in $H^{1}\left(\mathbb{R}^{N}\right)$, there is a unique $t_{o}=t(v)$ such that

$$
\begin{equation*}
J\left(t_{o} v\right)=\max _{t \geq 0} J(t v) \tag{2.6}
\end{equation*}
$$

Indeed, since

$$
\begin{equation*}
J(t v)=\frac{t^{2}}{2}\|v\|^{2}-\int_{\mathbb{R}^{N}}\left[F(t v)-\frac{2^{2^{*}}}{2^{*}}\left(v^{+}\right)^{2^{*}}\right] d x, \quad \text { for } t \geq 0 \tag{2.7}
\end{equation*}
$$

the maximum point $t_{o}$ of $J(t v)$ is given by

$$
\begin{equation*}
\|v\|^{2}=\int_{\mathbb{R}^{N}}\left[t_{o}^{-1} v f\left(t_{o} v\right)+t_{o}^{2^{*}-2}\left(v^{+}\right)^{2^{*}}\right] d x . \tag{2.8}
\end{equation*}
$$

We assume, without loss of generality that $0 \in \Omega$. Set $\Omega_{\varepsilon}=\left\{x \in \mathbb{R}^{N} ; \varepsilon x \in \Omega\right\}$.
The restriction of $J$ to $H_{o}^{1}\left(\Omega_{\varepsilon}\right)$ is the energy functional,

$$
\begin{equation*}
J(u)=\frac{1}{2} \int_{\Omega_{\varepsilon}}\left(|\nabla u|^{2}+u^{2}\right) d x-\int_{\Omega_{\varepsilon}}\left[F\left(u_{+}\right)+\frac{1}{2^{*}} u_{+}^{2^{*}}\right] d x, \quad u \in H_{o}^{1}\left(\Omega_{\varepsilon}\right), \tag{2.9}
\end{equation*}
$$

associated with the problem

$$
\begin{align*}
-\Delta u+u & =f(u)+u^{2^{*}-1} \quad \text { in } \Omega_{\varepsilon},  \tag{2.10}\\
u & =0 \quad \text { on } \partial \Omega_{\varepsilon} .
\end{align*}
$$

If $u_{\varepsilon}$ is a critical point of $J_{\varepsilon}$, the family

$$
\begin{equation*}
v_{\varepsilon}(x)=u_{\varepsilon}(z)=u_{\varepsilon}(\varepsilon x), \quad z=\varepsilon x \tag{2.11}
\end{equation*}
$$

is such that each $v_{\varepsilon}$ is a critical point of functional $J$ restricted to $H_{o}^{1}\left(\Omega_{\varepsilon}\right)$ at the level

$$
\begin{equation*}
b_{\varepsilon}=J\left(v_{\varepsilon}\right)=\inf _{v \in H_{o}\left(\Omega_{\varepsilon}\right) \backslash\{0\}} \max _{t \geq 0} J(t v) . \tag{2.12}
\end{equation*}
$$

It is easy to check that $b_{\varepsilon}=\varepsilon^{-N} \mathcal{C}_{\varepsilon}$ and from the definition of $c$ it follows that $b_{\varepsilon} \geq c$ for all $\varepsilon>0$.

We will start with the following property of $\left\{b_{\varepsilon}\right\}_{\varepsilon>0}$.
Lemma 2.3. For $\left\{b_{\varepsilon}\right\}_{\varepsilon>0}, \lim _{\varepsilon \rightarrow 0} b_{\varepsilon}=c$.
Proof. Fix $\omega$ a ground state solution of problem (2.4) and let $\psi_{\varepsilon}(x)=\varphi(\varepsilon x) \omega(x)$, where $\varphi$ is a $C^{1}$-function such that

$$
\varphi(x)= \begin{cases}1 & \text { if } x \in B_{1}  \tag{2.13}\\ 0 & \text { if } x \notin B_{2}\end{cases}
$$

$B_{1}=B_{\rho}(0), B_{2}=B_{2 \rho}(0) \subset \Omega$. Observe that $\psi_{\varepsilon} \rightarrow \omega$ in $H^{1}\left(\mathbb{R}^{N}\right)$ and the support of $\psi_{\varepsilon}$ is in $\Omega_{\varepsilon}$. By definition of $b_{\varepsilon}$, we have $t_{\varepsilon}>0$ such that

$$
\begin{equation*}
b_{\varepsilon} \leq \max _{t>0} J\left(t \psi_{\varepsilon}\right)=J\left(t_{\varepsilon} \psi_{\varepsilon}\right) \tag{2.14}
\end{equation*}
$$

From (2.8) and condition $\left(f_{3}\right)$ it follows that

$$
\begin{align*}
\left\|\psi_{\varepsilon}\right\|^{2} & =\int_{\mathbb{R}^{N}}\left[t_{\varepsilon}^{-1} \psi_{\varepsilon} f\left(t_{\varepsilon} \psi_{\varepsilon}\right)+t_{\varepsilon}^{2^{*}-2} \psi_{\varepsilon}^{2^{*}}\right] d x \\
& \geq \int_{\mathbb{R}^{N}}\left[\lambda t_{\varepsilon}^{q_{2}-1} \psi_{\varepsilon}^{q_{2}+1}+t_{\varepsilon}^{2^{*}-2} \psi_{\varepsilon}^{2^{*}}\right] d x \tag{2.15}
\end{align*}
$$

so that, $t_{\varepsilon}$ is bounded. Equality (2.15) and Remark 2.2 show that $t_{\varepsilon} \rightarrow t(\omega)=1$, as $\varepsilon \rightarrow 0$. Then we have $t_{\varepsilon} \psi_{\varepsilon} \rightarrow \omega$ in $H^{1}\left(\mathbb{R}^{N}\right)$ and

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} J\left(t_{\varepsilon} \psi_{\varepsilon}\right)=J(\omega)=c . \tag{2.16}
\end{equation*}
$$

Combining (2.14), (2.16), and the inequality $b_{\varepsilon} \geq c$, for all $\varepsilon>0$, we have proved this lemma.
Lemma 2.4. The inequality $c<(1 / N) S^{N / 2}$ holds, where $S$ is the best Sobolev constant for the embedding $D^{1,2}\left(\mathbb{R}^{N}\right) \hookrightarrow L^{2^{*}}\left(\mathbb{R}^{N}\right)$.
Proof. For each $h>0$, consider the function

$$
\begin{equation*}
\phi_{h}(x)=\frac{[N(N-2) h]^{(N-2) / 4}}{\left(h+|x|^{2}\right)^{(N-2) / 2}} \tag{2.17}
\end{equation*}
$$

We recall that $\phi_{h}$ satisfies the problem

$$
\begin{gather*}
-\Delta u=u^{2^{*}-1} \quad \text { in } \mathbb{R}^{N}, \\
u(x)>0, \quad \int_{\mathbb{R}^{N}}|\nabla u|^{2} d x<\infty,  \tag{2.18}\\
\int_{\mathbb{R}^{N}}\left|\nabla \phi_{h}\right|^{2} d x=\int_{\mathbb{R}^{N}} \phi_{h}^{2^{*}} d x=S^{N / 2} \quad \text { (see Talenti [12]). }
\end{gather*}
$$

Now, consider $\psi_{h}(x)=\varphi \phi_{h}(x) /\left\|\varphi \phi_{h}\right\|_{L^{2^{*}}\left(\mathbb{R}^{N}\right)}$, where $\varphi$ is the function defined in the proof of Lemma 2.3. From condition $\left(f_{3}\right)$ we have

$$
\begin{equation*}
J\left(t \psi_{h}\right) \leq \frac{t^{2}}{2} \int_{B_{2}}\left(\left|\nabla \psi_{h}\right|^{2}+\psi_{h}^{2}\right) d x-\frac{\lambda t^{q_{2}+1}}{q_{2}+1} \int_{B_{2}} \psi_{h}^{q_{2}+1} d x-\frac{t^{2^{*}}}{2^{*}} \tag{2.19}
\end{equation*}
$$

Using arguments as in [7], there exists $h>0$ such that

$$
\begin{equation*}
\max _{t \geq 0}\left\{\frac{t^{2}}{2} \int_{B_{2}}\left(\left|\nabla \psi_{h}\right|^{2}+\psi_{h}^{2}\right) d x-\frac{\lambda t^{q_{2}+1}}{q_{2}+1} \int_{B_{2}} \psi_{h}^{q_{2}+1} d x-\frac{t^{2^{*}}}{2^{*}}\right\}<\frac{1}{N} S^{N / 2} . \tag{2.20}
\end{equation*}
$$

Therefore, from (2.19) and (2.20) we have that

$$
\begin{equation*}
\max _{t \geq 0} J\left(t \psi_{h}\right)<\frac{1}{N} S^{N / 2} \tag{2.21}
\end{equation*}
$$

and the proof of the lemma is completed.
Notice that the same proof of Lemma 2.4 can be used to show that $b_{\varepsilon}<$ $(1 / N) S^{N / 2}$, for all $\varepsilon>0$. Using [3, Theorem 2.1], this inequality implies the existence of $v_{\varepsilon}$ and then the existence of $u_{\varepsilon}$.

Lemma 2.5. There are $\varepsilon_{0}>0$; a family $\left\{y_{\varepsilon}\right\}_{\left\{0<\varepsilon \leq \varepsilon_{0}\right\}} \subset \mathbb{R}^{N}, y_{\varepsilon} \in \Omega_{\varepsilon} ;$ constants $R>0$ and $\beta>0$ such that

$$
\begin{gather*}
\int_{B_{R}\left(y_{\varepsilon}\right)} v_{\varepsilon}^{2} d x \geq \beta>0, \quad \forall 0<\varepsilon \leq \varepsilon_{o}  \tag{2.22}\\
\lim _{\varepsilon \rightarrow 0} d\left(y_{\varepsilon}, \partial \Omega_{\varepsilon}\right)=\infty \tag{2.23}
\end{gather*}
$$

Proof. Start by showing that there is a family satisfying inequality (2.22). Arguing to the contrary, there is $\varepsilon_{n} \backslash 0$ such that for all $R>0$

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \sup _{x \in \mathbb{R}^{N}} \int_{B_{\mathbb{R}}(x)} v_{\varepsilon_{n}}^{2} d x=0 \tag{2.24}
\end{equation*}
$$

Using (Lions [6, Lemma I.1]) we have

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} v_{\varepsilon_{n}}^{q} d x=o_{n}(1), \quad \text { as } n \longrightarrow \infty, \forall 2<q<2^{*} \tag{2.25}
\end{equation*}
$$

and, from $\left(f_{1}\right)$ and $\left(f_{2}\right)$,

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} F\left(v_{\varepsilon_{n}}\right) d x=\int_{\mathbb{R}^{N}} v_{\varepsilon_{n}} f\left(v_{\varepsilon_{n}}\right) d x=o_{n}(1) \tag{2.26}
\end{equation*}
$$

Since $J^{\prime}\left(v_{\varepsilon_{n}}\right) \cdot v_{\varepsilon_{n}}=0$, we conclude from (2.26) that

$$
\begin{equation*}
\left\|v_{\varepsilon_{n}}\right\|^{2}=\int_{\mathbb{R}^{N}}{ }_{\varepsilon_{n}}^{2^{*}} d x+o_{n}(1) \tag{2.27}
\end{equation*}
$$

Let $\ell \geq 0$ be such that $\left\|v_{\varepsilon_{n}}\right\|^{2} \rightarrow \ell$. Passing to the limit in $J\left(v_{\varepsilon_{n}}\right)=b_{\varepsilon_{n}}$ and using (2.26) we have

$$
\begin{equation*}
\ell=N c \tag{2.28}
\end{equation*}
$$

and hence $\ell>0$. Now, using the definition of the constant $S$, we have

$$
\begin{equation*}
\left\|v_{\varepsilon_{n}}\right\|^{2} \geq S\left(\int_{\mathbb{R}^{N}} v_{\varepsilon_{n}}^{2^{*}} d x\right)^{2 / 2^{*}} \tag{2.29}
\end{equation*}
$$

Taking the limit in the above inequalities, as $n \rightarrow \infty$, we achieve that

$$
\begin{equation*}
\ell \geq S \ell^{2 / 2^{*}} \tag{2.30}
\end{equation*}
$$

and by (2.28), that

$$
\begin{equation*}
c \geq \frac{1}{N} S^{N / 2} \tag{2.31}
\end{equation*}
$$

which contradicts Lemma 2.4 and then (2.22) holds.
Finally, to establish (2.23), suppose the contrary. That is, there exist $\varepsilon_{n} \rightarrow 0$ and $R>0$ such that $\operatorname{dist}\left(y_{\varepsilon_{n}}, \partial \Omega_{\varepsilon_{n}}\right) \leq R$, hence $\operatorname{dist}\left(\varepsilon_{n} y_{\varepsilon_{n}}, \partial \Omega\right) \leq \varepsilon_{n} R$. Without loss of generality, we have $\varepsilon_{n} y_{\varepsilon_{n}} \rightarrow y_{o}$ for some $y_{o} \in \partial \Omega$. The arguments that follow can be found in [8].

Let $v$ be the unit interior normal to $\partial \Omega$ at $y_{o}$, and $\delta>0$ such that $B_{\delta}\left(y_{o}+\right.$ $\delta v) \subset \Omega$ and $B_{\delta}\left(y_{o}-\delta v\right) \cap \Omega=\varnothing$. Let $\Omega_{n}=\left\{x \in \mathbb{R}^{N}: y_{o}+\varepsilon_{n} x \in \Omega\right\}$ and $w_{n}(x)$ $=u_{\varepsilon_{n}}\left(y_{o}+\varepsilon_{n} x\right)$. This sequence $w_{n}$ is bounded in $H^{1}\left(\mathbb{R}^{N}\right),-\Delta w_{n}+w_{n}=f\left(w_{n}\right)+$ $w_{n}^{2^{*}-1}$ in $\Omega_{n}$,

$$
\begin{equation*}
\int_{B_{2 R}(0)} w_{n}^{2} d x \geq \int_{B_{R}\left(y_{\varepsilon_{n}}\right)} v_{\varepsilon_{n}}^{2} d x \geq \beta>0, \quad \forall n \tag{2.32}
\end{equation*}
$$

and we have that $w_{n}$ converges weakly to some $w$ in $H^{1}\left(\mathbb{R}^{N}\right)$.
Let $\mathbb{R}_{+, v}^{N}$ be the half space $\left\{x \in \mathbb{R}^{N}: x \cdot v>0\right\}$. Notice that $B_{\varepsilon_{n}^{-1}} \delta\left(\varepsilon_{n}^{-1} \delta v\right) \subset \Omega_{n}$ and $B_{\varepsilon_{n}^{-1}} \delta\left(-\varepsilon_{n}^{-1} \delta \nu\right) \cap \Omega_{n}=\varnothing$ and then we can prove that for all compacts $K_{+} \subset$ $\mathbb{R}_{+, v}^{N}$ and $K_{-} \subset \mathbb{R}_{-, v}^{N}=\mathbb{R}^{N} \backslash \overline{\mathbb{R}_{+, v}^{N}}$, we have $K_{+} \subset \Omega_{n}$ and $K_{-} \cap \Omega_{n}=\varnothing$, for $n$ large.

Then for each $\phi \in C_{o}^{\infty}\left(\mathbb{R}_{+, v}^{N}\right)$ such that $\operatorname{supp} \phi \subset \Omega_{n}$, we have

$$
\begin{equation*}
\int_{\mathbb{R}_{t, v}^{N}}\left(\nabla w_{n} \nabla \phi+w_{n} \phi\right) d x=\int_{\mathbb{R}_{+, v}^{N}}\left(f\left(w_{n}\right)+w_{n}^{2^{*}-1}\right) \phi d x . \tag{2.33}
\end{equation*}
$$

From (2.33), usual arguments show that $w \in H^{1}\left(\mathbb{R}^{N}\right) \cap C^{2}\left(\mathbb{R}_{+}^{N}\right)$ and satisfies $-\Delta w+w=f(w)+w^{2^{*}-1}$, in $\mathbb{R}_{+, v}^{N}$, and $w \equiv 0$ in $\mathbb{R}_{-, v}^{N}$. Theorem I.1, due to Esteban and Lions in [4], shows that $w \equiv 0$ which contradicts

$$
\begin{equation*}
\int_{B_{2 R}(0) \cap \mathbb{R}_{t, v}^{N}} w^{2} d x \geq \beta>0 . \tag{2.34}
\end{equation*}
$$

This completes the proof of the lemma.
Now we will consider the translation of $v_{\varepsilon}$, defined by $\omega_{\varepsilon}(x)=v_{\varepsilon}\left(x+y_{\varepsilon}\right)=$ $u_{\varepsilon}\left(\varepsilon y_{\varepsilon}+\varepsilon x\right)$ in $\widetilde{\Omega}_{\varepsilon}=\left\{x \in \mathbb{R}^{N} ; \varepsilon y_{\varepsilon}+\varepsilon x \in \Omega\right\}$ and $\omega_{\varepsilon}=0$ outside $\widetilde{\Omega}_{\varepsilon}$. From (2.23), any compact subset of $\mathbb{R}^{N}$ is contained in $\widetilde{\Omega}_{\varepsilon}$, for $\varepsilon$ sufficiently small.

From Lemma 2.5,

$$
\begin{equation*}
\int_{B_{R}(0)} \omega_{\varepsilon}^{2} d x \geq \beta>0, \quad \forall 0<\varepsilon \leq \varepsilon_{o} . \tag{2.35}
\end{equation*}
$$

Consider a sequence $\varepsilon_{n} \searrow 0$ and set $\widetilde{\Omega}_{n}=\widetilde{\Omega}_{\varepsilon_{n}}, \omega_{n}=\omega_{\varepsilon_{n}}, v_{n}=v_{\varepsilon_{n}}, y_{\varepsilon}=y_{\varepsilon_{n}}$.
We will prove that $\omega_{n}$ is bounded in the $L^{\infty}$ norm. In that case, $u_{\varepsilon}$ is also bounded in $L^{\infty}(\Omega)$ norm and the proof of Theorem 1.1 follows from the subcritical case, as Lemma 2.8 will show.

Since the sequence $\omega_{n}$ a translation of $v_{n}$, we have a uniform bound for $\left\|\omega_{n}\right\|$ and there is a $\omega_{o} \in H^{1}\left(\mathbb{R}^{N}\right)$ which is weak limit of $\omega_{n}$ in $H^{1}\left(\mathbb{R}^{N}\right)$. From (2.35) we have $\omega_{o} \neq 0$. We can write limit (2.23) in the following form

$$
\begin{equation*}
\lim _{n \rightarrow \infty} d\left(0, \partial \widetilde{\Omega}_{n}\right)=\infty \tag{2.36}
\end{equation*}
$$

Then for each $\phi \in C_{o}^{\infty}\left(\mathbb{R}^{N}\right)$ and large $n$ such that $\operatorname{supp} \phi \subset \widetilde{\Omega}_{n}$, we have

$$
\begin{equation*}
\int_{\mathbb{R}^{N}}\left(\nabla \omega_{n} \nabla \phi+\omega_{n} \phi\right) d x=\int_{\mathbb{R}^{N}}\left(f\left(\omega_{n}\right)+\omega_{n}^{2^{*}-1}\right) \phi d x, \quad \forall n \tag{2.37}
\end{equation*}
$$

From (2.37), usual arguments show that $\omega_{o}$ is a solution of problem (2.4), hence a critical point of $J$, and $J\left(\omega_{o}\right) \geq c$.

Lemma 2.6. The sequence $\omega_{n}$ converges to $\omega_{o}$ in $H^{1}\left(\mathbb{R}^{N}\right)$ and $J\left(\omega_{o}\right)=c$.
Proof. This fact comes from Lemma 2.5 and Fatou's lemma applied in the positive sequence $\omega_{n} f\left(\omega_{n}\right)-\theta F\left(\omega_{n}\right)$. Observe that

$$
\begin{align*}
b_{\varepsilon_{n}} & =J\left(v_{n}\right)-\frac{1}{\theta} J^{\prime}\left(v_{n}\right) v_{n} \\
& =\left(\frac{\theta-2}{2 \theta}\right)\left\|v_{n}\right\|^{2}+\frac{1}{\theta} \int_{\mathbb{R}^{N}}\left[v_{n} f\left(v_{n}\right)-\theta F\left(v_{n}\right)\right]+\left(\frac{2^{*}-\theta}{2^{*} \theta}\right) \int_{\mathbb{R}^{N}} v_{n}^{2^{*}}  \tag{2.38}\\
& =\left(\frac{\theta-2}{2 \theta}\right)\left\|\omega_{n}\right\|^{2}+\frac{1}{\theta} \int_{\mathbb{R}^{N}}\left[\omega_{n} f\left(\omega_{n}\right)-\theta F\left(\omega_{n}\right)\right]+\left(\frac{2^{*}-\theta}{2 * \theta}\right) \int_{\mathbb{R}^{N}} \omega_{n}^{2^{*}}
\end{align*}
$$

From (2.38)

$$
\begin{align*}
c & \leq J\left(\omega_{o}\right)=J\left(\omega_{o}\right)-\frac{1}{\theta} J^{\prime}\left(\omega_{o}\right) \omega_{o} \\
& =\left(\frac{\theta-2}{2 \theta}\right)\left\|\omega_{o}\right\|^{2}+\frac{1}{\theta} \int_{\mathbb{R}^{N}}\left[\omega_{o} f\left(\omega_{o}\right)-\theta F\left(\omega_{o}\right)\right]+\left(\frac{2^{*}-\theta}{2^{*} \theta}\right) \int_{\mathbb{R}^{N}} \omega_{o}^{2^{*}} \\
& \leq \liminf \left(\frac{\theta-2}{2 \theta}\right)\left\|\omega_{n}\right\|^{2}+\frac{1}{\theta} \int_{\mathbb{R}^{N}}\left[\omega_{n} f\left(\omega_{n}\right)-\theta F\left(\omega_{n}\right)\right]+\left(\frac{2^{*}-\theta}{2^{*} \theta}\right) \int_{\mathbb{R}^{N}} \omega_{n}^{2^{*}} \\
& =\lim _{n \rightarrow \infty} b_{\varepsilon_{n}}=c . \tag{2.39}
\end{align*}
$$

We have proved that $J\left(\omega_{o}\right)=c$ and then (2.39) becomes an equality.

Combining (2.39) with the three following inequalities:

$$
\begin{align*}
\left\|\omega_{o}\right\|^{2} & \leq \liminf \left\|\omega_{n}\right\|^{2}, \\
\int_{\mathbb{R}^{N}}\left[\omega_{o} f\left(\omega_{o}\right)-\theta F\left(\omega_{o}\right)\right] d x & \leq \liminf \int_{\mathbb{R}^{N}}\left[\omega_{n} f\left(\omega_{n}\right)-\theta F\left(\omega_{n}\right)\right] d x,  \tag{2.40}\\
\int_{\mathbb{R}^{N}} \omega_{o}^{2^{*}} d x & \leq \liminf \int_{\mathbb{R}^{N^{2}}} \omega_{n}^{2^{*}} d x
\end{align*}
$$

we conclude that $\left\|\omega_{n}\right\| \rightarrow\left\|\omega_{o}\right\|$ and then $\omega_{n} \rightarrow \omega_{o}$ in $H^{1}\left(\mathbb{R}^{N}\right)$.
We are ready to conclude the proof of our main result. From Proposition 1.2 and Remark 1.3 with $q(x)=\omega_{n}^{2^{*}-2} \in L^{N / 2} ; g(x, s)=f(s)+s^{2^{*}}-s$, we have $\omega_{n} \in$ $L^{t}$ for all $t \geq 2$ and

$$
\begin{equation*}
\left\|\omega_{n}\right\|_{L^{t}} \leq C_{t} \tag{2.41}
\end{equation*}
$$

where $C_{t}$ does not depend on $n$.
Now we will make use of a very particular version of [5, Theorem 8.17], due to Trudinger.
Proposition 2.7. Suppose that $t>N, g \in L^{t / 2}(\Lambda)$, and $u \in H_{o}^{1}(\Lambda)$ satisfies (in the weak sense)

$$
\begin{equation*}
-\Delta u+u \leq \tilde{g}(x) \tag{2.42}
\end{equation*}
$$

where $\Lambda$ is an open subset of $\mathbb{R}^{N}$. Then for any ball $B_{2 R}(y) \subset \Lambda$,

$$
\begin{equation*}
\sup _{B_{R}(y)} u \leq C\left(\left\|u^{+}\right\|_{L^{2}\left(B_{2 R}(y)\right)}+\|g\|_{L^{t / 2}\left(B_{2 R}(y)\right)}\right) \tag{2.43}
\end{equation*}
$$

where $C$ depends on $N, t$, and $R$.
We know that each $\omega_{n}$ satisfies

$$
\begin{equation*}
-\Delta \omega_{n}+\omega_{n}=\omega_{n}^{2^{*}-1}+f\left(\omega_{n}\right), \quad \text { in } \widetilde{\Omega}_{n} \tag{2.44}
\end{equation*}
$$

and this implies that

$$
\begin{equation*}
-\Delta \omega_{n}+\omega_{n} \leq g_{n}(x)=\omega_{n}^{2^{*}-1}+f\left(\omega_{n}\right), \quad \text { in } \mathbb{R}^{N} \tag{2.45}
\end{equation*}
$$

in the weak sense.
Since (2.41) holds, $\left\|g_{n}\right\|_{L^{t}}$ is bounded from above for some $t>N$. Using Proposition 2.7 in (2.45) we have

$$
\begin{equation*}
\sup _{B_{1}(y)} \omega_{n} \leq C\left(\left\|\omega_{n}\right\|_{L^{2}\left(B_{2 R}(y)\right)}+\left\|g_{n}\right\|_{L^{t}\left(B_{2 R}(y)\right)}\right) \tag{2.46}
\end{equation*}
$$

for all $y \in \mathbb{R}^{N}$, which implies that there is a constant $a>0$, independent of $n$, such that

$$
\begin{equation*}
\omega_{n}(x) \leq a, \quad \forall x \in \mathbb{R}^{N} \tag{2.47}
\end{equation*}
$$

It follows that there is a $\varepsilon_{o}>0$ such that

$$
\begin{equation*}
u_{\varepsilon}(z) \leq a, \quad \forall z \in \Omega, \quad \forall \varepsilon<\varepsilon_{o} \tag{2.48}
\end{equation*}
$$

To conclude the proof observe that $\mathcal{u}_{\varepsilon}$ becomes a solution of the subcritical case (1.8) with $h$ given by

$$
h(s)= \begin{cases}f(s)+s^{2^{*}-1}, & \text { if } s \leq a  \tag{2.49}\\ f(s)+\frac{\left(2^{*}-1\right)}{(\theta-1)} a^{2^{*}-\theta} s^{\theta-1}-\frac{\left(2^{*}-\theta\right)}{(\theta-1)} a^{2^{*}-1}, & \text { if } s>a\end{cases}
$$

where $\theta>2$ is that one fixed in condition $\left(f_{4}\right)$. It is easy to check that $h$ is a $C^{1, \alpha}$ function, $h$ and $H(s)=\int_{o}^{s} h(\tau) d \tau$ satisfy $\left(\mathrm{f}_{1}\right),\left(\mathrm{f}_{2}\right),\left(\mathrm{f}_{3}\right),\left(\mathrm{f}_{4}\right)$, and $\left(\mathrm{f}_{5}\right)$. Let $\tilde{J}_{\varepsilon}$ be the $C^{1}$-functional on $H_{o}^{1}(\Omega)$ given by

$$
\begin{equation*}
\tilde{J}_{\varepsilon}(u)=\frac{1}{2} \int_{\Omega}\left(\varepsilon^{2}|\nabla u|^{2}+u^{2}\right) d z-\int_{\Omega} H(u) d z . \tag{2.50}
\end{equation*}
$$

Since $f(s)+s^{2^{*}-1} \geq h(s)$ for all $s>0$, we have that

$$
\begin{equation*}
J_{\varepsilon}(u) \leq \tilde{J}_{\varepsilon}(u), \quad \forall u \in H_{o}^{1}(\Omega), \tag{2.51}
\end{equation*}
$$

$J_{\varepsilon}\left(u_{\varepsilon}\right)=\tilde{J}_{\varepsilon}\left(u_{\varepsilon}\right), J_{\varepsilon}^{\prime}\left(u_{\varepsilon}\right)=\tilde{J}_{\varepsilon}^{\prime}\left(u_{\varepsilon}\right)=0$. We conclude that $u_{\varepsilon}$ is a least-energy solution of the subcritical problem (1.8).
Lemma 2.8. (i) If $\tilde{\mathcal{C}}_{\varepsilon}$ is the minimax level of $\tilde{J}_{\varepsilon}$, then $\tilde{\mathcal{c}}_{\varepsilon}=c_{\varepsilon}$;
(ii) each $u_{\varepsilon}$ is a critical point of $\tilde{J}_{\varepsilon}$ in the minimax level and satisfies (1.8).

Since global problem (1.9) has a unique nondegenerate positive solution (cf. [9, 10]), Theorem 1.1 comes from [8, Theorem 2.2] applied to the functional $\tilde{J}_{\varepsilon}$, and the asymptotic profile comes from [8, Theorem 2.3].

## Appendix

Let $\Lambda$ be some general domain in $\mathbb{R}^{N}$ (bounded or unbounded). We will start with the following lemma due to Brézis and Kato [2].

Lemma A.1. Let $q \in L^{N / 2}(\Lambda)$ be a nonnegative function. Then, for every $\varepsilon>0$, there is a constant $\sigma_{\varepsilon}=\sigma(\varepsilon, q)>0$ such that

$$
\begin{equation*}
\int_{\Lambda} q(x) u^{2} d x \leq \varepsilon \int_{\Lambda}|\nabla u|^{2} d x+\sigma_{\varepsilon} \int_{\Lambda} u^{2} d x, \quad \forall u \in H_{o}^{1}(\Lambda) \tag{A.1}
\end{equation*}
$$

Remark A.2. If $q_{k} \rightarrow q$ in $L^{N / 2}(\Lambda)$, we can choose a constant $\sigma_{\varepsilon}$ independent of $k$. That is, $\sigma\left(\varepsilon, q_{k}\right)=\sigma_{\varepsilon}$ and

$$
\begin{equation*}
\int_{\Lambda} q_{k}(x) u^{2} d x \leq \varepsilon \int_{\Lambda}|\nabla u|^{2} d x+\sigma_{\varepsilon} \int_{\Lambda} u^{2} d x, \quad \forall u \in H_{o}^{1}(\Lambda), k \in \mathbb{N} . \tag{A.2}
\end{equation*}
$$

Proof. Let $\sigma_{\varepsilon}=\sigma(\varepsilon, q)>0$ be such that

$$
\begin{equation*}
\|q\|_{L^{N / 2}\left(\left\{q \geq \sigma_{\varepsilon}\right\}\right)} \leq \varepsilon S, \tag{A.3}
\end{equation*}
$$

where $S$ is a best constant in the Sobolev immersion $H_{o}^{1}(\Lambda) \hookrightarrow L^{2^{*}}(\Lambda)$, where $2^{*}=2 N /(N-2)$. For all $u \in H_{o}^{1}(\Lambda)$, we have

$$
\begin{align*}
\int_{\Lambda} q(x) u^{2} d x & =\int_{\left\{q \geq \sigma_{\varepsilon}\right\}} q(x) u^{2} d x+\int_{\left\{q \leq \sigma_{\varepsilon}\right\}} q(x) u^{2} d x \\
& \leq \sigma_{\varepsilon} \int_{\left\{q \leq \sigma_{\varepsilon}\right\}} u^{2} d x+\int_{\left\{q \geq \sigma_{\varepsilon}\right\}} q(x) u^{2} d x  \tag{A.4}\\
& \leq \sigma_{\varepsilon} \int_{\Lambda} u^{2} d x+\|q\|_{L^{N / 2}\left(\left\{q \geq \sigma_{\varepsilon}\right\}\right)}\|u\|_{L^{*}\left(\left\{q \geq \sigma_{\varepsilon}\right\}\right)}^{2} .
\end{align*}
$$

Inequality (A.1) follows from Sobolev estimate and the choice of $\sigma_{\varepsilon}$.
Remark 1.3 follows from the proof of Lemma A. 1 and the inequality

$$
\begin{equation*}
\int_{\Lambda} q_{k}(x) u^{2} d x \leq \int_{\Lambda} q(x) u^{2} d x+\left\|q_{k}-q\right\|_{L^{N / 2}(\Lambda)}\|u\|_{L^{*}(\Lambda)^{*}}^{2} \tag{A.5}
\end{equation*}
$$

Proof of Proposition 1.2. For any $n \in \mathbb{N}$ and $p>0$, consider $A_{n}=\left\{x \in \Lambda:|v|^{p} \leq\right.$ $n\}, B_{n}=\Lambda \backslash A_{n}$, and define $v_{n}$ by

$$
\begin{equation*}
v_{n}=v|v|^{2 p} \quad \text { in } A_{n}, \quad v_{n}=n^{2} v \quad \text { in } B_{n} . \tag{A.6}
\end{equation*}
$$

Observe that $v_{n} \in H_{o}^{1}(\Lambda), v_{n} \leq|v|^{2 p+1}$ and

$$
\begin{equation*}
\nabla v_{n}=(2 p+1)|v|^{2 p} \nabla v \quad \text { in } A_{n}, \quad \nabla v_{n}=n^{2} \nabla v \quad \text { in } B_{n} . \tag{A.7}
\end{equation*}
$$

So, using $v_{n}$ as a test function

$$
\begin{equation*}
\int_{\Lambda} \nabla v \nabla v_{n} d x=\int_{\Lambda} g(x, v) v_{n} d x \tag{A.8}
\end{equation*}
$$

Using (A.7), we have

$$
\begin{align*}
& (2 p+1) \int_{A_{n}}|v|^{2 p}|\nabla v|^{2} d x+n^{2} \int_{B_{n}}|\nabla v|^{2} d x  \tag{A.9}\\
& \quad \leq \int_{\Lambda}\left|g(x, v) v_{n}\right| d x \leq \int_{\Lambda}\left(q(x)+C_{g}\right)\left|v v_{n}\right| d x .
\end{align*}
$$

Now consider

$$
\begin{equation*}
\omega_{n}=v|v|^{p} \quad \text { in } A_{n}, \quad \omega_{n}=n v \quad \text { in } B_{n} \tag{A.10}
\end{equation*}
$$

Notice that $\omega_{n}^{2}=v v_{n} \leq|v|^{2(p+1)}$ and

$$
\begin{equation*}
\nabla \omega_{n}=(p+1)|v|^{p} \nabla v \quad \text { in } A_{n}, \quad \nabla v_{n}=n \nabla v \quad \text { in } B_{n} . \tag{A.11}
\end{equation*}
$$

Therefore,

$$
\begin{equation*}
\int_{\Lambda}\left|\nabla \omega_{n}\right|^{2} d x=(p+1)^{2} \int_{A_{n}}|v|^{2 p}|\nabla v|^{2} d x+n^{2} \int_{B_{n}}|\nabla v|^{2} d x . \tag{A.12}
\end{equation*}
$$

Combining (A.9) and (A.12), we obtain

$$
\begin{equation*}
\frac{2 p+1}{(p+1)^{2}} \int_{\Lambda}\left|\nabla \omega_{n}\right|^{2} d x \leq \int_{\Lambda}\left(q(x)+C_{g}\right) \omega_{n}^{2} d x \tag{A.13}
\end{equation*}
$$

Let $\sigma_{p}$ be given by Lemma A. 1 with $\varepsilon=(2 p+1) / 2(p+1)^{2}$. Then

$$
\begin{equation*}
\int_{\Lambda}\left|\nabla \omega_{n}\right|^{2} d x \leq \tilde{C}_{p} \int_{\Lambda} \omega_{n}^{2} d x \tag{A.14}
\end{equation*}
$$

where $\tilde{C}_{n}=\left(2(p+1)^{2} /(2 p+1)\right)\left(C_{g}+\sigma_{p}\right)$. Suppose that $v \in L^{2(p+1)}(\Lambda)$ for some $p \geq 2$. Applying Sobolev immersion in inequality (A.14) we have

$$
\begin{equation*}
\left[\int_{A_{n}} \omega_{n}^{2^{*}} d x\right]^{2 / 2^{*}} \leq\left[\int_{\Lambda} \omega_{n}^{2^{*}} d x\right]^{2 / 2^{*}} \leq S \tilde{C}_{p} \int_{\Lambda}|v|^{2(p+1)} d x \tag{A.15}
\end{equation*}
$$

that is,

$$
\begin{equation*}
\left[\int_{A_{n}}|v|^{2^{*}(p+1)} d x\right]^{2 / 2^{*}} d x \leq C_{p} \int_{\Lambda}|v|^{2(p+1)} d x \tag{A.16}
\end{equation*}
$$

where

$$
\begin{equation*}
C_{p}=\frac{2(p+1)^{2}}{2 p+1} S\left(C_{g}+\sigma_{p}\right) \tag{A.17}
\end{equation*}
$$

Now, passing to the limit in (A.16) we have $v \in L^{2^{*}(p+1)}(\Lambda)$ and

$$
\begin{equation*}
\|v\|_{L^{2^{*}(p+1)}(\Lambda)} \leq C_{p}\|v\|_{L^{2(p+1)}(\Lambda)} . \tag{A.18}
\end{equation*}
$$

The proof follows from the following iteration argument: let $p_{1}$ a positive such that $2\left(p_{1}+1\right)=2^{*}$. It is easy to see that $0<p_{1}$ and $v \in L^{2\left(p_{1}+1\right)}(\Lambda)$. Using inequality (A.18) we have

$$
\begin{equation*}
v \in L^{2^{*}\left(p_{1}+1\right)}(\Lambda) \tag{A.19}
\end{equation*}
$$

Now choose $p_{2}$ such that $2\left(p_{2}+1\right)=2^{*}\left(p_{1}+1\right)$. It is easy to see that $0<$ $p_{1}<p_{2}$ and $v \in L^{p_{2}+1}(\Lambda)$. Using inequality (A.18) we have

$$
\begin{equation*}
v \in L^{2^{*}\left(p_{2}+1\right)}(\Lambda) \tag{A.20}
\end{equation*}
$$

Continuing with this iteration we obtain an increasing sequence $p_{k}$ given by $2\left(p_{k+1}+1\right)=2^{*}\left(p_{k}+1\right)$ such that $v \in L^{2\left(p_{k+1}+1\right)}(\Lambda)$ for all $k \in \mathbb{N}$. From

$$
\begin{equation*}
p_{k+1}+1=\frac{N}{N-2}\left(p_{k}+1\right) \tag{A.21}
\end{equation*}
$$

it follows that

$$
\begin{equation*}
p_{k+1}+1=\left[\frac{N}{N-2}\right]^{k} 2^{*} \tag{A.22}
\end{equation*}
$$

This shows that $p_{k}$ goes to $\infty$ and therefore,

$$
\begin{equation*}
v \in L^{p}(\Lambda), \quad \forall p \geq 2 \tag{A.23}
\end{equation*}
$$

Remark A.3. Proposition 1.2 is valid for positive subsolutions of problem (1.12) as we can check in its proof.

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