## Research Article

# On the Study of Local Solutions for a Generalized Camassa-Holm Equation 

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The pseudoparabolic regularization technique is employed to study the local well-posedness of strong solutions for a nonlinear dispersive model, which includes the famous Camassa-Holm equation. The local well-posedness is established in the Sobolev space $H^{s}(R)$ with $s>3 / 2$ via a limiting procedure.

## 1. Introduction

In recent years, extensive research has been carried out worldwide to study highly nonlinear equations including the Camassa-Holm (CH) equation and its various generalizations [1-6]. It is shown in [7-9] that the inverse spectral or scattering approach is a powerful technique to handle the Camassa-Holm equation and analyze its dynamics. It is pointed out in [10-12] that the CH equation gives rise to geodesic flow of a certain invariant metric on the BottVirasoro group, and this geometric illustration leads to a proof that the Least Action Principle holds. Li and Olver [13] established the local well-posedness to the CH model in the Sobolev space $H^{s}(R)$ with $s>3 / 2$ and gave conditions on the initial data that lead to finite time blow-up of certain solutions. Constantin and Escher [14] proved that the blow-up occurs in the form of breaking waves, namely, the solution remains bounded but its slope becomes unbounded in finite time. Hakkaev and Kirchev [15] investigated a generalized form of the Camassa-Holm equation with high order nonlinear terms and obtained the orbit stability of the traveling wave solutions under certain assumptions. Lai and Wu [16] discussed a generalized Camassa-Holm model and acquired its local existence and uniqueness. Recently, Li et al. [17] investigated the generalized Camassa-Holm equation

$$
\begin{equation*}
u_{t}-u_{t x x}+k u^{m} u_{x}+(m+3) u^{m+1} u_{x}=(m+2) u^{m} u_{x} u_{x x}+u^{m+1} u_{x x x}, \tag{1.1}
\end{equation*}
$$

where $m \geq 0$ is a natural number and $k \geq 0$. The authors in [17] assume that the initial value satisfies the sign condition and establish the global existence of solutions for (1.1).

In this paper, we will study the following generalization of (1.1):

$$
\begin{equation*}
u_{t}-u_{t x x}+k u^{m} u_{x}+(m+3) u^{m+1} u_{x}=(m+2) u^{m} u_{x} u_{x x}+u^{m+1} u_{x x x}+\lambda\left(u-u_{x x}\right), \tag{1.2}
\end{equation*}
$$

where $m \geq 0$ is a natural number, $k \geq 0$, and $\lambda$ is a constant.
The objective of this paper is to study the local well-posedness of (1.2). Its local wellposedness of strong solutions in the Sobolev space $H^{s}(R)$ with $s>3 / 2$ is investigated by using the pseudoparabolic regularization method. Comparing with the work by Li et al. [17], (1.2) considered in this paper possesses a conservation law different to that in [17] (see Lemma 3.2 in Section 3). Also (1.2) contains a dissipative term $\lambda\left(u-u_{x x}\right)$, which causes difficulty to establish its local and global existence in the Sobolev space. It should be mentioned that the existence and uniqueness of local strong solutions for the generalized nonlinear Camassa-Holm models like (1.2) have never been investigated in the literatures.

The organization of this work is as follows. The main result is given in Section 2. Section 3 establishes several lemmas, and the last section gives the proof of the main result.

## 2. Main Result

Firstly, we introduce several notations.
$L^{p}=L^{p}(R)(1 \leq p<+\infty)$ is the space of all measurable functions $h$ such that $\|h\|_{L^{p}}^{p}=\int_{R}|h(t, x)|^{p} d x<\infty$. We define $L^{\infty}=L^{\infty}(R)$ with the standard norm $\|h\|_{L^{\infty}}=$ $\inf _{m(e)=0} \sup _{x \in R \backslash e}|h(t, x)|$. For any real number $s, H^{s}=H^{s}(R)$ denotes the Sobolev space with the norm defined by

$$
\begin{equation*}
\|h\|_{H^{s}}=\left(\int_{R}\left(1+|\xi|^{2}\right)^{s}|\widehat{h}(t, \xi)|^{2} d \xi\right)^{1 / 2}<\infty \tag{2.1}
\end{equation*}
$$

where $\widehat{h}(t, \xi)=\int_{R} e^{-i x \xi} h(t, x) d x$.
For $T>0$ and nonnegative number $s, C\left([0, T) ; H^{s}(R)\right)$ denotes the Frechet space of all continuous $H^{s}$-valued functions on $[0, T)$. We set $\Lambda=\left(1-\partial_{x}^{2}\right)^{1 / 2}$. For simplicity, throughout this paper, we let $c$ denote any positive constant that is independent of parameter $\varepsilon$.

We consider the Cauchy problem of (1.2)

$$
\begin{gather*}
u_{t}-u_{t x x}=-\frac{k}{m+1}\left(u^{m+1}\right)_{x}-\frac{m+3}{m+2}\left(u^{m+2}\right)_{x}+\frac{1}{m+2} \partial_{x}^{3}\left(u^{m+2}\right) \\
-(m+1) \partial_{x}\left(u^{m} u_{x}^{2}\right)+u^{m} u_{x} u_{x x}+\lambda\left(u-u_{x x}\right), \quad k \geq 0, m \geq 0  \tag{2.2}\\
u(0, x)=u_{0}(x)
\end{gather*}
$$

Now, we give our main results for problem of (2.2).

Theorem 2.1. Suppose that the initial function $u_{0}(x)$ belongs to the Sobolev space $H^{s}(R)$ with $s>$ $3 / 2$ and $\lambda$ is a constant. Then, there is a $T>0$, which depends on $\left\|u_{0}\right\|_{H^{s}}$, such that problem (2.2) has a unique solution $u(t, x)$ satisfying

$$
\begin{equation*}
u(t, x) \in C\left([0, T) ; H^{s}(R)\right) \bigcap C^{1}\left([0, T) ; H^{s-1}(R)\right) . \tag{2.3}
\end{equation*}
$$

## 3. Local Well-Posedness

In order to prove Theorem 2.1, we consider the associated regularized problem

$$
\begin{gather*}
u_{t}-u_{t x x}+\varepsilon u_{t x x x x}=-\frac{k}{m+1}\left(u^{m+1}\right)_{x}-\frac{m+3}{m+2}\left(u^{m+2}\right)_{x}+\frac{1}{m+2} \partial_{x}^{3}\left(u^{m+2}\right) \\
-(m+1) \partial_{x}\left(u^{m} u_{x}^{2}\right)+u^{m} u_{x} u_{x x}+\lambda\left(u-u_{x x}\right),  \tag{3.1}\\
u(0, x)=u_{0}(x),
\end{gather*}
$$

where the parameter $\varepsilon$ satisfies $0<\varepsilon<1 / 4$.
Lemma 3.1. For $s \geq 1$ and $f(x) \in H^{s}(R)$ and letting $k_{1}>0$ be an integer such that $k_{1} \leq s-1$, $f, f^{\prime}, \ldots, f^{k_{1}}$ are uniformly continuous bounded functions that converge to 0 at $x= \pm \infty$.

The proof of Lemma 3.1 was stated on page 559 by Bona and Smith [18].
Lemma 3.2. If $u(t, x) \in H^{s}(s>7 / 2)$ is a solution to problem (3.1), it holds that

$$
\begin{equation*}
\int_{R}\left(u^{2}+u_{x}^{2}+\varepsilon u_{x x}^{2}\right) d x=\int_{R}\left(u_{0}^{2}+u_{0 x}^{2}+\varepsilon u_{0 x x}^{2}\right) d x+2 \lambda \int_{0}^{t} \int_{R}\left(u^{2}+u_{x}^{2}\right) d x \tag{3.2}
\end{equation*}
$$

Proof. Using Lemma 3.1, we have $u(t, \pm \infty)=u_{x}(t, \pm \infty)=u_{x x}(t, \pm \infty)=u_{x x x}(t, \pm \infty)=0$. The integration by parts results in

$$
\begin{align*}
\int_{R} u^{m+2} u_{x x x} d x & =\int_{R} u^{m+2} d u_{x x}=\left.u^{m+2} u_{x x}\right|_{-\infty} ^{+\infty}-(m+2) \int_{R} u^{m+1} u_{x} u_{x x} d x \\
& =-(m+2) \int_{R} u^{m+1} u_{x} u_{x x} d x \tag{3.3}
\end{align*}
$$

Direct calculation and integration by parts give rise to

$$
\begin{align*}
& \frac{1}{2} \frac{d}{d t} \int_{R}\left(u^{2}+u_{x}^{2}+\varepsilon u_{x x}^{2}\right) d x \\
& = \\
& =\int_{R}\left(u u_{t}+u_{x} u_{t x}+\varepsilon u_{x x} u_{t x x}\right) d x \\
& =\int_{R} u\left(u_{t}-u_{t x x}+\varepsilon u_{t x x x x}\right) d x  \tag{3.4}\\
& =\int_{R} u\left[-\frac{k}{m+1}\left(u^{m+1}\right)_{x}-\frac{m+3}{m+2}\left(u^{m+2}\right)_{x}+\frac{1}{m+2} \partial_{x}^{3}\left(u^{m+2}\right)\right. \\
& \quad=\int_{R} u\left[(m+2) u^{m} u_{x} u_{x x}+u^{m+1} u_{x x x}\right] d x+\lambda \int_{R}\left(u^{2}-u u_{x x}\right) d x \\
& \quad=\int_{R}\left[(m+2) u^{m+1} u_{x} u_{x x}+u^{m+2} u_{x x x}\right] d x+\lambda \int_{R}\left(u^{2}+u_{x}^{2}\right) d x \\
& \quad= \\
& \quad \lambda \int_{R}\left(u^{2}+u_{x}^{2}\right) d x,
\end{align*}
$$

in which we have used (3.3). From (3.4), we obtain the conservation law (3.2).
Lemma 3.3. Let $s \geq 7 / 2$. The function $u(t, x)$ is a solution of problem (3.1) and the initial value $u_{0}(x) \in H^{s}$. Then, the following inequality holds:

$$
\begin{align*}
& \|u\|_{H^{1}}^{2} \leq \int_{R}\left(u_{0}^{2}+u_{0 x}^{2}+\varepsilon u_{0 x x}^{2}\right) d x, \quad \text { if } \lambda \leq 0  \tag{3.5}\\
& \|u\|_{H^{1}}^{2} \leq e^{2 \lambda t} \int_{R}\left(u_{0}^{2}+u_{0 x}^{2}+\varepsilon u_{0 x x}^{2}\right) d x, \quad \text { if } \lambda>0
\end{align*}
$$

For $q \in(0, s-1]$, there is a constant $c$ independent of $\varepsilon$ such that

$$
\begin{align*}
\int_{R}\left(\Lambda^{q+1} u\right)^{2} d x \leq & \int_{R}\left[\left(\Lambda^{q+1} u_{0}\right)^{2}+\varepsilon\left(\Lambda^{q} u_{0 x x}\right)^{2}\right] d x \\
& +c \int_{0}^{t}\|u\|_{H^{q+1}}^{2}\left(|\lambda|+\left(\|u\|_{L^{\infty}}^{m-1}+\|u\|_{L^{\infty}}^{m}\right)\left\|u_{x}\right\|_{L^{\infty}}+\|u\|_{L^{\infty}}^{m-1}\left\|u_{x}\right\|_{L^{\infty}}^{2}\right) d \tau \tag{3.6}
\end{align*}
$$

For $q \in[0, s-1]$, there is a constant $c$ independent of $\varepsilon$ such that

$$
\begin{equation*}
(1-2 \varepsilon)\left\|u_{t}\right\|_{H^{q}} \leq c\|u\|_{H^{q+1}}\left(|\lambda|+\left(\|u\|_{L^{\infty}}^{m-1}+\|u\|_{L^{\infty}}^{m}\right)\|u\|_{H^{1}}+\|u\|_{L^{\infty}}^{m}\left\|u_{x}\right\|_{L^{\infty}}+\|u\|_{L^{\infty}}^{m-1}\left\|u_{x}\right\|_{L^{\infty}}^{2}\right) \tag{3.7}
\end{equation*}
$$

The proof of this lemma is similar to that of Lemma 3.5 in [17]. Here we omit it.
Lemma 3.4. Let $r$ and $q$ be real numbers such that $-r<q \leq r$. Then,

$$
\begin{gather*}
\|u v\|_{H^{q}} \leq c\|u\|_{H^{r}}\|v\|_{H^{q}}, \quad \text { if } r>\frac{1}{2}  \tag{3.8}\\
\|u v\|_{H^{r+q-1 / 2}} \leq c\|u\|_{H^{r}}\|v\|_{H^{q}}, \quad \text { if } r<\frac{1}{2} .
\end{gather*}
$$

This lemma can be found in [19] or [20].
Lemma 3.5. Let $u_{0}(x) \in H^{s}(R)$ with $s>3 / 2$. Then, the Cauchy problem (3.1) has a unique solution $u(t, x) \in C\left([0, T] ; H^{s}(R)\right)$, where $T>0$ depends on $\left\|u_{0}\right\|_{H^{s}(R)}$. If $s \geq 7 / 2$, the solution $u \in C\left([0,+\infty) ; H^{s}\right)$ exists for all time.

Proof. Letting $D=\left(1-\partial_{x}^{2}+\varepsilon \partial_{x}^{4}\right)^{-1}$, we know that $D: H^{s} \rightarrow H^{s+4}$ is a bounded linear operator. Applying the operator $D$ on both sides of the first equation of system (3.1) and then integrating the resultant equation with respect to $t$ over the interval $(0, t)$, we get

$$
\begin{gather*}
u(t, x)=u_{0}(x)+\int_{0}^{t} D\left[-\frac{k}{m+1}\left(u^{m+1}\right)_{x}-\frac{m+3}{m+2}\left(u^{m+2}\right)_{x}+\frac{1}{m+2} \partial_{x}^{3}\left(u^{m+2}\right)\right.  \tag{3.9}\\
\left.-(m+1) \partial_{x}\left(u^{m} u_{x}^{2}\right)+u^{m} u_{x} u_{x x}+\lambda\left(u-u_{x x}\right)\right] d t
\end{gather*}
$$

Suppose that both $u$ and $v$ are in the closed ball $B_{M_{0}}(0)$ of radius $M_{0}$ about the zero function in $C\left([0, T] ; H^{s}(R)\right)$ and $A$ is the operator in the right-hand side of (3.9). For any fixed $t \in[0, T]$, we obtain

$$
\begin{align*}
& \| \int_{0}^{t} D\left[-\frac{k}{m+1}\left(u^{m+1}\right)_{x}-\frac{m+3}{m+2}\left(u^{m+2}\right)_{x}+\frac{1}{m+2} \partial_{x}^{3}\left(u^{m+2}\right)\right. \\
& \left.\quad-(m+1) \partial_{x}\left(u^{m} u_{x}^{2}\right)+u^{m} u_{x} u_{x x}+\lambda\left(u-u_{x x}\right)\right] d t \\
& -\int_{0}^{t} D\left[-\frac{k}{m+1}\left(v^{m+1}\right)_{x}-\frac{m+3}{m+2}\left(v^{m+2}\right)_{x}+\frac{1}{m+2} \partial_{x}^{3}\left(v^{m+2}\right)\right. \\
& \left.\quad-(m+1) \partial_{x}\left(v^{m} v_{x}^{2}\right)+v^{m} v_{x} v_{x x}+\lambda\left(v-v_{x x}\right)\right] d t \|_{H^{s}}  \tag{3.10}\\
& \leq T C_{1}\left(\sup _{0 \leq t \leq T}\|u-v\|_{H^{s}}+\sup _{0 \leq t \leq T}\left\|u^{m+1}-v^{m+1}\right\|_{H^{s}}+\sup _{0 \leq t \leq T}\left\|u^{m+2}-v^{m+2}\right\|_{H^{s}}\right. \\
& \left.\quad+\sup _{0 \leq t \leq T}\left\|D \partial_{x}\left[u^{m} u_{x}^{2}-v^{m} v_{x}^{2}\right]\right\|_{H^{s}}+\sup _{0 \leq t \leq T}\left\|D\left[u^{m} u_{x} u_{x x}-v^{m} v_{x} v_{x x}\right]\right\|_{H^{s}}\right),
\end{align*}
$$

where $C_{1}$ may depend on $\varepsilon$. The algebraic property of $H^{s_{0}}(R)$ with $s_{0}>1 / 2$ derives

$$
\begin{align*}
&\left\|u^{m+2}-v^{m+2}\right\|_{H^{s}}=\left\|(u-v)\left(u^{m+1}+u^{m} v+\cdots+u v^{m}+v^{m+1}\right)\right\|_{H^{s}} \\
& \leq\|(u-v)\|_{H^{s}} \sum_{j=0}^{m+1}\|u\|_{H^{s}}^{m+1-j}\|v\|_{H^{s}}^{j} \\
& \leq M_{0}^{m+1}\|u-v\|_{H^{s}} \\
&\left\|u^{m+1}-v^{m+1}\right\|_{H^{s}} \leq M_{0}^{m}\|u-v\|_{H^{s}}  \tag{3.11}\\
&\left\|D \partial_{x}\left(u^{m} u_{x}^{2}-v^{m} v_{x}^{2}\right)\right\|_{H^{s}} \leq\left\|D \partial_{x}\left[u^{m}\left(u_{x}^{2}-v_{x}^{2}\right)\right]\right\|_{H^{s}}+\left\|D \partial_{x}\left[v_{x}^{2}\left(u^{m}-v^{m}\right)\right]\right\|_{H^{s}} \\
& \leq C\left(\left\|u^{m}\left(u_{x}^{2}-v_{x}^{2}\right)\right\|_{H^{s-1}}+\left\|v_{x}^{2}\left(u^{m}-v^{m}\right)\right\|_{H^{s-1}}\right) \\
& \leq C M_{0}^{m+1}\|u-v\|_{H^{s} .}
\end{align*}
$$

Using the first inequality of Lemma 3.4 gives rise to

$$
\begin{aligned}
\left\|D\left[u^{m} u_{x} u_{x x}-v^{m} v_{x} v_{x x}\right]\right\|_{H^{s}} & =\left\|\frac{1}{2} D\left[u^{m}\left(u_{x}^{2}\right)_{x}-v^{m}\left(v_{x}^{2}\right)_{x}\right]\right\|_{H^{s}} \\
& \leq \frac{1}{2}\left(\left\|D\left[u^{m}\left(u_{x}^{2}-v_{x}^{2}\right)_{x}\right]\right\|_{H^{s}}+\left\|D\left[\left(v_{x}^{2}\right)_{x}\left(u^{m}-v^{m}\right)\right]\right\|_{H^{s}}\right) \\
& \leq C\left(\left\|u^{m}\left(u_{x}^{2}-v_{x}^{2}\right)_{x}\right\|_{H^{s-2}}+\left\|\left(v_{x}^{2}\right)_{x}\left(u^{m}-v^{m}\right)\right\|_{H^{s-2}}\right) \\
& \leq C\left(\left\|u^{m}\right\|_{H^{s}}\left\|u_{x}^{2}-v_{x}^{2}\right\|_{H^{s-1}}+\left\|v_{x}^{2}\right\|_{H^{s-1}}\left\|u^{m}-v^{m}\right\|_{H^{s}}\right) \\
& \leq C M_{0}^{m+1}\|u-v\|_{H^{s}}
\end{aligned}
$$

where C may depend on $\varepsilon$. From (3.11)-(3.12), we obtain

$$
\begin{equation*}
\|A u-A v\|_{H^{s}} \leq \theta\|u-v\|_{H^{s}} \tag{3.13}
\end{equation*}
$$

where $\theta=T C_{2}\left(M_{0}^{m}+M_{0}^{m+1}\right)$ and $C_{2}$ is independent of $0<t<T$. Choosing $T$ sufficiently small such that $\theta<1$, we know that $A$ is a contraction. Similarly, it follows from (3.10) that

$$
\begin{equation*}
\|A u\|_{H^{s}} \leq\left\|u_{0}\right\|_{H^{s}}+\theta\|u\|_{H^{s}} \tag{3.14}
\end{equation*}
$$

Choosing $T$ sufficiently small such that $\theta M_{0}+\left\|u_{0}\right\|_{H^{s}}<M_{0}$, we deduce that $A$ maps $B_{M_{0}}(0)$ to itself. It follows from the contraction-mapping principle that the mapping $A$ has a unique fixed point $u$ in $B_{M_{0}}(0)$. It completes the proof.

From the above and Lemma 3.2, we have

$$
\begin{equation*}
\int_{R}\left(u^{2}+u_{x}^{2}+\varepsilon u_{x x}^{2}\right) d x \leq e^{2|\lambda| t} \int_{R}\left(u_{0}^{2}+u_{0 x}^{2}+\varepsilon u_{0 x x}^{2}\right) d x \tag{3.15}
\end{equation*}
$$

Therefore,

$$
\begin{equation*}
\left\|u_{x}\right\|_{L^{\infty}} \leq C_{\varepsilon} e^{2|\lambda| t} \int_{R}\left(u_{0}^{2}+u_{0 x}^{2}+\varepsilon u_{0 x x}^{2}\right) d x \tag{3.16}
\end{equation*}
$$

which together with Lemma 3.3 completes the proof of the global existence.
Setting $\phi_{\varepsilon}(x)=\varepsilon^{-1 / 4} \phi\left(\varepsilon^{-1 / 4} x\right)$ with $0<\varepsilon<1 / 4$ and $u_{\varepsilon 0}=\phi_{\varepsilon} \star u_{0}$, we know that $u_{\varepsilon 0} \in C^{\infty}$ for any $u_{0} \in H^{s}, s>0$. From Lemma 3.5, it derives that the Cauchy problem

$$
\begin{gather*}
u_{t}-u_{t x x}+\varepsilon u_{t x x x x}=-\frac{k}{m+1}\left(u^{m+1}\right)_{x}-\frac{m+3}{m+2}\left(u^{m+2}\right)_{x}+\frac{1}{m+2} \partial_{x}^{3}\left(u^{m+2}\right) \\
-(m+1) \partial_{x}\left(u^{m} u_{x}^{2}\right)+u^{m} u_{x} u_{x x}+\lambda\left(u-u_{x x}\right),  \tag{3.17}\\
u(0, x)=u_{\varepsilon 0}(x), \quad x \in R,
\end{gather*}
$$

has a unique solution $u_{\varepsilon}(t, x) \in C^{\infty}\left([0, \infty) ; H^{\infty}\right)$.
Furthermore, we have the following.
Lemma 3.6. For $s>0, u_{0} \in H^{s}$, it holds that

$$
\begin{gather*}
\left\|u_{\varepsilon 0 x}\right\|_{L^{\infty}} \leq c\left\|u_{0 x}\right\|_{L^{\infty}},  \tag{3.18}\\
\left\|u_{\varepsilon 0}\right\|_{H^{q}} \leq c, \quad \text { if } q \leq s,  \tag{3.19}\\
\left\|u_{\varepsilon 0}\right\|_{H^{q}} \leq c \varepsilon^{(s-q) / 4}, \quad \text { if } q>s,  \tag{3.20}\\
\left\|u_{\varepsilon 0}-u_{0}\right\|_{H^{q}} \leq c \varepsilon^{(s-q) / 4}, \quad \text { if } q \leq s,  \tag{3.21}\\
\left\|u_{\varepsilon 0}-u_{0}\right\|_{H^{s}}=o(1), \tag{3.22}
\end{gather*}
$$

where $c$ is a constant independent of $\varepsilon$.
The proof of Lemma 3.6 can be found in [16].

Remark 3.7. For $s \geq 1$, using $\left\|u_{\varepsilon}\right\|_{L^{\infty}} \leq c\left\|u_{\varepsilon}\right\|_{H^{1 / 2+}} \leq c\left\|u_{\varepsilon}\right\|_{H^{1}},\left\|u_{\varepsilon}\right\|_{H^{1}}^{2} \leq c \int_{R}\left(u_{\varepsilon}^{2}+u_{\varepsilon x}^{2}\right) d x$, (3.5), (3.19), and (3.20), we know that,

$$
\begin{align*}
\left\|u_{\varepsilon}\right\|_{L^{\infty}}^{2} \leq c\left\|u_{\varepsilon}\right\|_{H^{1}}^{2} & \leq c e^{2|\lambda| t} \int_{R}\left(u_{\varepsilon 0}^{2}+u_{\varepsilon 0 x}^{2}+\varepsilon u_{\varepsilon 0 x x}^{2}\right) d x \\
& \leq c e^{2|\lambda| t}\left(\left\|u_{\varepsilon 0}\right\|_{H^{1}}^{2}+\varepsilon\left\|u_{\varepsilon 0}\right\|_{H^{2}}^{2}\right)  \tag{3.23}\\
& \leq c e^{|2 \lambda| t}\left(c+c \varepsilon \times \varepsilon^{(s-2) / 2}\right) \\
& \leq c_{0} e^{2|\lambda| t}
\end{align*}
$$

where $c_{0}$ is independent of $\varepsilon$ and $t$.
Lemma 3.8. Suppose $u_{0}(x) \in H^{s}(R)$ with $s \geq 1$ such that $\left\|u_{0 x}\right\|_{L^{\infty}}<\infty$. Let $u_{\varepsilon 0}$ be defined as in system (3.17). Then, there exist two positive constants $T$ and $c$, which are independent of $\varepsilon$, such that the solution $u_{\varepsilon}$ of problem (3.17) satisfies $\left\|u_{\varepsilon x}\right\|_{L^{\infty}} \leq c$ for any $t \in[0, T)$.

Here we omit the proof of Lemma 3.8 since it is similar to Lemma 3.9 presented in [17].

Lemma 3.9 (see Li and Olver [13]). If $u$ and $f$ are functions in $H^{q+1} \cap\left\{\left\|u_{x}\right\|_{L^{\infty}}<\infty\right\}$, then

$$
\left|\int_{R} \Lambda^{q} u \Lambda^{q}(u f)_{x} d x\right| \leq\left\{\begin{array}{c}
c_{q}\|f\|_{H^{q+1}}\|u\|_{H^{q}}^{2}, \quad q \in\left(\frac{1}{2}, 1\right]  \tag{3.24}\\
c_{q}\left(\|f\|_{H^{q+1}}\|u\|_{H^{q}}\|u\|_{L_{\infty}}\right. \\
\left.+\left\|f_{x}\right\|_{L_{\infty}}\|u\|_{H^{q}}^{2}+\|f\|_{H^{q}}\|u\|_{H^{q}}\left\|u_{x}\right\|_{L_{\infty}}\right) \\
q \in(0, \infty)
\end{array}\right.
$$

Lemma 3.10 (see Lai and $\mathrm{Wu}[16]$ ). For $u, v \in H^{s}(R)$ with $s>3 / 2, w=u-v, q>1 / 2$, and $a$ natural number $n$, it holds that

$$
\begin{equation*}
\left|\int_{R} \Lambda^{s} w \Lambda^{s}\left(u^{n+1}-v^{n+1}\right)_{x} d x\right| \leq c\left(\|w\|_{H^{s}}\|w\|_{H^{q}}\|v\|_{H^{s+1}}+\|w\|_{H^{s}}^{2}\right) \tag{3.25}
\end{equation*}
$$

Lemma 3.11 (see Lai and $W u[16]$ ). If $1 / 2<q<\min \{1, s-1\}$ and $s>3 / 2$, then for any functions $w, f$ defined on $R$, it holds that

$$
\begin{align*}
& \left|\int_{R} \Lambda^{q} w \Lambda^{q-2}(w f)_{x} d x\right| \leq c\|w\|_{H^{q}}^{2}\|f\|_{H^{q}}  \tag{3.26}\\
& \left|\int_{R} \Lambda^{q} w \Lambda^{q-2}\left(w_{x} f_{x}\right)_{x} d x\right| \leq c\|w\|_{H^{q}}^{2}\|f\|_{H^{s}} \tag{3.27}
\end{align*}
$$

Lemma 3.12. For problem (3.17), $s>3 / 2$, and $u_{0} \in H^{s}(R)$, there exist two positive constants $c$ and $M$, which are independent of $\varepsilon$, such that the following inequalities hold for any sufficiently small $\varepsilon$ and $t \in[0, T)$ :

$$
\begin{align*}
\left\|u_{\varepsilon}\right\|_{H^{s}} & \leq M e^{c t}, \\
\left\|u_{\varepsilon}\right\|_{H^{s+k_{1}}} & \leq \varepsilon^{-k_{1} / 4} M e^{c t}, \quad k_{1}>0  \tag{3.28}\\
\left\|u_{\varepsilon t}\right\|_{H^{s+k_{1}}} & \leq \varepsilon^{-\left(k_{1}+1\right) / 4} M e^{c t}, \quad k_{1}>-1 .
\end{align*}
$$

Slightly modifying the methods presented in [16] can complete the proof of Lemma 3.12.

Our next step is to demonstrate that $u_{\varepsilon}$ is a Cauchy sequence. Let $u_{\varepsilon}$ and $u_{\delta}$ be solutions of problem (3.17), corresponding to the parameters $\varepsilon$ and $\delta$, respectively, with $0<\varepsilon<\delta<1 / 4$, and let $w=u_{\varepsilon}-u_{\delta}$. Then, $w$ satisfies the problem

$$
\begin{align*}
& (1-\varepsilon) w_{t}-\varepsilon w_{x x t}+(\delta-\varepsilon)\left(u_{\delta t}+u_{\delta x x t}\right) \\
& =\left(1-\partial_{x}^{2}\right)^{-1}\left[-\varepsilon w_{t}+(\delta-\varepsilon) u_{\delta t}-\frac{k}{m+1} \partial_{x}\left(u_{\varepsilon}^{m+1}-u_{\delta}^{m+1}\right)-\partial_{x}\left(u_{\varepsilon}^{m+2}-u_{\delta}^{m+2}\right)\right. \\
&  \tag{3.29}\\
& -\partial_{x}\left[\partial_{x}\left(u_{\varepsilon}^{m+1}\right) \partial_{x} w+\partial_{x}\left(u_{\varepsilon}^{m+1}-u_{\varepsilon}^{m+1}\right) \partial_{x} u_{\delta}\right] \\
& \left.+\left[u_{\varepsilon}^{m} u_{\varepsilon x} u_{\varepsilon x x}-u_{\delta}^{m} u_{\delta x} u_{\delta x x}\right]\right]-\frac{1}{m+2} \partial_{x}\left(u_{\varepsilon}^{m+2}-u_{\delta}^{m+2}\right)+\lambda w  \tag{3.30}\\
& w(x, 0)=w_{0}(x)=u_{\varepsilon 0}(x)-u_{\delta 0}(x)
\end{align*}
$$

Lemma 3.13. For $s>3 / 2, u_{0} \in H^{s}(R)$, there exists $T>0$ such that the solution $u_{\varepsilon}$ of (3.17) is a Cauchy sequence in $C\left([0, T] ; H^{s}(R)\right) \bigcap C^{1}\left([0, T] ; H^{s-1}(R)\right)$.

Proof. For $q$ with $1 / 2<q<\min \{1, s-1\}$, multiplying both sides of (3.29) by $\Lambda^{q} w \Lambda^{q}$ and then integrating with respect to $x$ give rise to

$$
\begin{aligned}
\frac{1}{2} \frac{d}{d t} \int_{R} & {\left[(1-\varepsilon)\left(\Lambda^{q} w\right)^{2}+\varepsilon\left(\Lambda^{q} w_{x}\right)^{2}\right] d x } \\
= & (\varepsilon-\delta) \int_{R}\left(\Lambda^{q} w\right)\left[\left(\Lambda^{q} u_{\delta t}\right)+\left(\Lambda^{q} u_{\delta x x t}\right)\right] d x-\varepsilon \int_{R} \Lambda^{q} w \Lambda^{q-2} w_{t} d x \\
& +(\delta-\varepsilon) \int_{R} \Lambda^{q} w \Lambda^{q-2} u_{\delta t} d x-\frac{1}{m+2} \int_{R}\left(\Lambda^{q} w\right) \Lambda^{q}\left(u_{\varepsilon}^{m+2}-u_{\delta}^{m+2}\right)_{x} d x
\end{aligned}
$$

$$
\begin{align*}
& -\frac{k}{m+1} \int_{R} \Lambda^{q} w \Lambda^{q-2}\left(u_{\varepsilon}^{m+1}-u_{\delta}^{m+1}\right)_{x} d x-\int_{R} \Lambda^{q} w \Lambda^{q-2}\left(u_{\varepsilon}^{m+2}-u_{\delta}^{m+2}\right)_{x} d x \\
& -\int_{R} \Lambda^{q} w \Lambda^{q-2}\left[\partial_{x}\left(u_{\varepsilon}^{m+1}\right) \partial_{x} w\right]_{x} d x-\int_{R} \Lambda^{q} w \Lambda^{q-2}\left[\partial_{x}\left(u_{\varepsilon}^{m+1}-u_{\delta}^{m+1}\right) \partial_{x} u_{\delta}\right]_{x} d x \\
& +\int_{R} \Lambda^{q} w \Lambda^{q-2}\left[u_{\varepsilon}^{m} u_{\varepsilon x} u_{\varepsilon x x}-u_{\delta}^{m} u_{\delta x} u_{\delta x x}\right] d x+\lambda \int_{R} \Lambda^{q} w \Lambda^{q} w d x \tag{3.31}
\end{align*}
$$

It follows from the Schwarz inequality that

$$
\begin{align*}
& \frac{d}{d t} \int\left[(1-\varepsilon)\left(\Lambda^{q} w\right)^{2}+\varepsilon\left(\Lambda^{q} w_{x}\right)^{2}\right] d x \\
& \quad \leq c\left\{\left\|\Lambda^{q} w\right\|_{L^{2}}\left[(\delta-\varepsilon)\left(\left\|\Lambda^{q} u_{\delta t}\right\|_{L^{2}}+\left\|\Lambda^{q} u_{\delta x x t}\right\|_{L^{2}}\right)+\varepsilon\left\|\Lambda^{q-2} w_{t}\right\|_{L^{2}}+(\delta-\varepsilon)\left\|\Lambda^{q-2} u_{\delta t}\right\|_{L^{2}}\right]\right. \\
& \\
& +|\lambda| \int_{R}\left(\Lambda^{q} w\right)^{2} d x+\left|\int_{R} \Lambda^{q} w \Lambda^{q}\left(u_{\varepsilon}^{m+2}-u_{\delta}^{m+2}\right)_{x} d x\right| \\
& \\
&  \tag{3.32}\\
& \left|\iint^{q} w \Lambda^{q-2}\left(u_{\varepsilon}^{m+1}-u_{\delta}^{m+1}\right)_{x} d x\right|+\left|\int \Lambda^{q} w \Lambda^{q-2}\left(u_{\varepsilon}^{m+2}-u_{\delta}^{m+2}\right)_{x} d x\right| \\
& \\
& +\left|\int_{R} \Lambda^{q} w \Lambda^{q-2}\left[\partial_{x}\left(u_{\varepsilon}^{m+1}\right) \partial_{x} w\right]_{x} d x\right|+\left|\int_{R} \Lambda^{q} w \Lambda^{q-2}\left[\partial_{x}\left(u_{\varepsilon}^{m+1}-u_{\delta}^{m+1}\right) \partial_{x} u_{\delta}\right]_{x} d x\right| \\
& \\
& \left.\quad+\left|\int_{R} \Lambda^{q} w \Lambda^{q-2}\left[u_{\varepsilon}^{m} u_{\varepsilon x} u_{\varepsilon x x}-u_{\delta}^{m} u_{\delta x} u_{\delta x x}\right] d x\right|\right\}
\end{align*}
$$

Using the first inequality in Lemma 3.9, we have

$$
\begin{align*}
\left|\int_{R} \Lambda^{q} w \Lambda^{q}\left(u_{\varepsilon}^{m+2}-u_{\delta}^{m+2}\right)_{x} d x\right| & =\left|\int_{R} \Lambda^{q} w \Lambda^{q}\left(w g_{m+1}\right)_{x} d x\right|  \tag{3.33}\\
& \leq c\|w\|_{H^{q}}^{2}\left\|g_{m+1}\right\|_{H^{q+1}}
\end{align*}
$$

where $g_{m+1}=\sum_{j=0}^{m+1} u_{\varepsilon}^{m+1-j} u_{\delta}^{j}$. For the last three terms in (3.32), using Lemmas 3.4 and 3.12, $1 / 2<q<\min \{1, s-1\}, s>3 / 2$, the algebra property of $H^{s_{0}}$ with $s_{0}>1 / 2$, and (3.23),
we have

$$
\begin{gather*}
\left|\int_{R} \Lambda^{q} w \Lambda^{q-2}\left(\partial_{x}\left(u_{\varepsilon}^{m+1}\right) \partial_{x} w\right)_{x} d x\right| \leq c\|w\|_{H^{q}}^{2}\left\|u_{\varepsilon}\right\|_{H^{s}}^{m+1},  \tag{3.34}\\
\left|\int_{R} \Lambda^{q} w \Lambda^{q-2}\left(\partial_{x}\left(u_{\varepsilon}^{m+1}-u_{\delta}^{m+1}\right) \partial_{x} u_{\delta}\right)_{x} d x\right| \\
\leq c\|w\|_{H^{q}}\left\|u_{\delta}\right\|_{H^{s}}\left\|u_{\varepsilon}^{m+1}-u_{\delta}^{m+1}\right\|_{H^{q}}  \tag{3.35}\\
\leq c\|w\|_{H^{q}}^{2}\left\|u_{\delta}\right\|_{H^{s}}, \\
\left|\int_{R} \Lambda^{q} w \Lambda^{q-2}\left[u_{\varepsilon}^{m} u_{\varepsilon x} u_{\varepsilon x x}-u_{\delta}^{m} u_{\delta x} u_{\delta x x}\right] d x\right| \\
\left.\leq c\|w\|_{H^{q}} \|\left(u_{\varepsilon}^{m}-u_{\delta}^{m}\right)\left(u_{\varepsilon x}^{2}\right)\right)_{x}+u_{\delta}^{m}\left[u_{\varepsilon x}^{2}-u_{\delta x}^{2}\right]_{x} \|_{H^{q-2}} \\
\left.\leq c\|w\|_{H^{q}}\left(\left\|\left(u_{\varepsilon}^{m}-u_{\delta}^{m}\right)\left(u_{\varepsilon x}^{2}\right)_{x}\right\|_{H^{q-1}}+\| u_{\delta}^{m}\left[u_{\varepsilon x}^{2}-u_{\delta x}^{2}\right]\right]_{x} \|_{H^{q-2}}\right)  \tag{3.36}\\
\leq c\|w\|_{H^{q}}\left(\left\|u_{\varepsilon}^{m}-u_{\delta}^{m}\right\|_{H^{q}}\left\|\left(u_{\varepsilon x}^{2}\right)_{x}\right\|_{H^{q-1}}+\left\|u_{\delta}^{m}\right\|_{H^{s}}\left\|\left[u_{\varepsilon x}^{2}-u_{\delta x}^{2}\right]_{x}\right\|_{H^{q-2}}\right) \\
\leq c\|w\|_{H^{q}}\left(\|w\|_{H^{q}}\left\|g_{g^{m-1}}\right\|_{H^{q}}\|u\|_{H^{s}}^{2}+\left\|u_{\delta}^{m}\right\|_{H^{s}}\left\|u_{\varepsilon x}+u_{\delta x}\right\|_{H^{q}}\|w\|_{H^{q}}\right) \\
\leq c\|w\|_{H^{q}}^{2}\left(\left\|g_{m-1}\right\|_{H^{q}}\|u\|_{H^{s}}^{2}+\left\|u_{\delta}^{m}\right\|_{H^{s}}\left\|u_{\varepsilon x}+u_{\delta x}\right\|_{H^{q}}\right) .
\end{gather*}
$$

Using (3.26), we derive that the inequality

$$
\begin{align*}
\left|\int_{R} \Lambda^{q} w \Lambda^{q-2}\left(u_{\varepsilon}^{m+2}-u_{\delta}^{m+2}\right)_{x} d x\right| & =\left|\int_{R} \Lambda^{q} w \Lambda^{q-2}\left(w g_{m+1}\right)_{x} d x\right|  \tag{3.37}\\
& \leq c\left\|g_{m+1}\right\|_{H^{q}}\|w\|_{H^{q}}^{2}
\end{align*}
$$

holds for some constant $c$, where $g_{m+1}=\sum_{j=0}^{m+1} u_{\varepsilon}^{m+1-j} u_{\delta}^{j}$. Using the algebra property of $H^{q}$ with $q>1 / 2, q+1<s$ and Lemma 3.11, we have $\left\|g_{m}\right\|_{H^{q+1}} \leq c$ for $t \in(0, \widetilde{T}]$. Then, it follows from (3.28) and (3.33)-(3.37) that there is a constant $c$ depending on $\tilde{T}$ such that the estimate

$$
\begin{equation*}
\frac{d}{d t} \int_{R}\left[(1-\varepsilon)\left(\Lambda^{q} w\right)^{2}+\varepsilon\left(\Lambda^{q} w_{x}\right)^{2}\right] d x \leq c\left(\delta^{r}\|w\|_{H^{q}}+\|w\|_{H^{q}}^{2}\right) \tag{3.38}
\end{equation*}
$$

holds for any $t \in[0, \widetilde{T})$, where $\gamma=1$ if $s \geq 3+q$ and $\gamma=(1+s-q) / 4$ if $s<3+q$. Integrating (3.38) with respect to $t$, one obtains the estimate

$$
\begin{align*}
\frac{1}{2}\|w\|_{H^{q}}^{2} & =\frac{1}{2} \int_{R}\left(\Lambda^{q} w\right)^{2} d x \\
& \leq \int_{R}\left[(1-\varepsilon)\left(\Lambda^{q} w\right)^{2}+\varepsilon\left(\Lambda^{q} w\right)^{2}\right] d x  \tag{3.39}\\
& \leq \int_{R}\left[\left(\Lambda^{q} w_{0}\right)^{2}+\varepsilon\left(\Lambda^{q} w_{0 x}\right)^{2}\right] d x+c \int_{0}^{t}\left(\delta^{r}\|w\|_{H^{q}}+\|w\|_{H^{q}}^{2}\right) d \tau
\end{align*}
$$

Applying the Gronwall inequality and using (3.20) and (3.22) yield

$$
\begin{equation*}
\|u\|_{H^{q}} \leq c \delta^{(s-q) / 4} e^{c t}+\delta^{\gamma}\left(e^{c t}-1\right) \tag{3.40}
\end{equation*}
$$

for any $t \in[0, \widetilde{T})$.
Multiplying both sides of (3.29) by $\Lambda^{s} w \Lambda^{s}$ and integrating the resultant equation with respect to $x$, one obtains

$$
\begin{align*}
& \frac{1}{2} \frac{d}{d t} \int_{R}\left[(1-\varepsilon)\left(\Lambda^{s} w\right)^{2}+\varepsilon\left(\Lambda^{s} w_{x}\right)^{2}\right] d x \\
&=(\varepsilon-\delta) \int_{R}\left(\Lambda^{s} w\right)\left[\left(\Lambda^{s} u_{\delta t}\right)+\left(\Lambda^{s} u_{\delta x x t}\right)\right] d x-\varepsilon \int_{R} \Lambda^{s} w \Lambda^{s-2} w_{t} d x \\
&+(\delta-\varepsilon) \int_{R} \Lambda^{s} w \Lambda^{s-2} u_{\delta t} d x-\frac{k}{m+1} \int_{R}\left(\Lambda^{s} w\right) \Lambda^{s}\left(u_{\varepsilon}^{m+1}-u_{\delta}^{m+1}\right)_{x} d x  \tag{3.41}\\
&-\frac{1}{m+2} \int_{R}\left(\Lambda^{s} w\right) \Lambda^{s}\left(u_{\varepsilon}^{m+2}-u_{\delta}^{m+2}\right)_{x} d x-\int_{R} \Lambda^{s} w \Lambda^{s-2}\left(u_{\varepsilon}^{m+2}-u_{\delta}^{m+2}\right)_{x} d x \\
&-\int_{R} \Lambda^{s} w \Lambda^{s-2}\left[\partial_{x}\left(u_{\varepsilon}^{m+1}\right) \partial_{x} w\right]_{x} d x-\int_{R} \Lambda^{s} w \Lambda^{s-2}\left[\partial_{x}\left(u_{\varepsilon}^{m+1}-u_{\delta}^{m+1}\right) \partial_{x} u_{\delta}\right]_{x} d x \\
&+\int_{R} \Lambda^{s} w \Lambda^{s-2}\left[u_{\varepsilon}^{m} u_{\varepsilon x} u_{\varepsilon x x}-u_{\delta}^{m} u_{\delta x} u_{\delta x x}\right] d x+\lambda \int_{R}\left(\Lambda^{s} w\right)^{2} d x
\end{align*}
$$

From Lemma 3.12, we have

$$
\begin{equation*}
\left|\int_{R} \Lambda^{s} w \Lambda^{s-2}\left(u_{\varepsilon}^{m+2}-u_{\delta}^{m+2}\right)_{x} d x\right| \leq c_{3}\left\|g_{m+1}\right\|_{H^{s}}\|w\|_{H^{s}}^{2} \tag{3.42}
\end{equation*}
$$

From Lemma 3.10, it holds that

$$
\begin{equation*}
\left|\int_{R} \Lambda^{s} w \Lambda^{s}\left(u_{\varepsilon}^{m+2}-u_{\delta}^{m+2}\right)_{x} d x\right| \leq c\left(\|w\|_{H^{s}}\|w\|_{H^{q}}\left\|u_{\delta}\right\|_{H^{s+1}}+\|w\|_{H^{s}}^{2}\right) \tag{3.43}
\end{equation*}
$$

Using the Cauchy-Schwartz inequality and the algebra property of $H^{s_{0}}$ with $s_{0}>1 / 2$, for $s>3 / 2$, we have

$$
\begin{align*}
& \left|\int_{R} \Lambda^{s} w \Lambda^{s-2}\left[\partial_{x}\left(u_{\varepsilon}^{m+1}\right) \partial_{x} w\right]_{x} d x\right| \\
& =\left|\int_{R} \Lambda^{q} w \Lambda^{s-2}\left[\partial_{x}\left(u_{\varepsilon}^{m+1}\right) \partial_{x} w\right]_{x} d x\right| \\
& \leq c\left\|\Lambda^{s} w\right\|_{L^{2}}\left\|\Lambda^{s-2}\left[\partial_{x}\left(u_{\varepsilon}^{m+1}\right) \partial_{x} w\right]_{x}\right\|_{L^{2}}  \tag{3.44}\\
& \leq c\|w\|_{H^{q}}\left\|\partial_{x}\left(u_{\varepsilon}^{m+1}\right) \partial_{x} w\right\|_{H^{s-1}} \\
& \leq c\left\|u_{\varepsilon_{H^{s}}}^{m+1}\right\|\|w\|_{H^{s}}^{2}, \\
& \left|\int_{R} \Lambda^{s} w \Lambda^{s-2}\left[\partial_{x}\left(u_{\varepsilon}^{m+1}-u_{\delta}^{m+1}\right) \partial_{x} u_{\delta}\right]_{x} d x\right| \\
& \leq c\|w\|_{H^{s}}\left\|\Lambda^{s-2}\left[\partial_{x}\left(u_{\varepsilon}^{m+1}-u_{\delta}^{m+1}\right) \partial_{x} u_{\delta}\right]_{x}\right\|_{L^{2}}  \tag{3.45}\\
& \leq c\left\|u_{\delta}\right\|_{H^{s}}\left\|g_{m}\right\|_{H^{s}}\|w\|_{H^{s}}^{2}, \\
& \left|\int_{R} \Lambda^{s} w \Lambda^{s-2}\left[u_{\varepsilon}^{m} u_{\varepsilon x} u_{\varepsilon x x}-u_{\delta}^{m} u_{\delta x} u_{\delta x x}\right] d x\right| \\
& \leq c\|w\|_{H^{s}}\left(\left\|\left(u_{\varepsilon}^{m}-u_{\delta}^{m}\right)\left(u_{\varepsilon x}^{2}\right)_{x}\right\|_{H^{s-2}}+\left\|u_{\delta}^{m}\left[u_{\varepsilon x}^{2}-u_{\delta x}^{2}\right]_{x}\right\|_{H^{s-2}}\right)  \tag{3.46}\\
& \leq c\|w\|_{H^{s}}\left(\left\|u_{\varepsilon}^{m}-u_{\delta}^{m}\right\|_{H^{s}}\left\|\left(u_{\varepsilon x}^{2}\right)_{x}\right\|_{H^{s-2}}+\left\|u_{\delta}^{m}\right\|_{H^{s}}\left\|\left[u_{\varepsilon x}^{2}-u_{\delta x}^{2}\right]_{x}\right\|_{H^{s-2}}\right) \\
& \leq c\|w\|_{H^{s}}^{2},
\end{align*}
$$

in which we have used Lemma 3.4 and the bounded property of $\left\|u_{\varepsilon}\right\|_{H^{s}}$ and $\left\|u_{\delta}\right\|_{H^{s}}$ (see Remark 3.7). It follows from (3.41)-(3.46) and the inequalities (3.28) and (3.40) that there exists a constant $c$ depending on $m$ such that

$$
\begin{align*}
& \frac{d}{d t} \int_{R}\left[(1-\varepsilon)\left(\Lambda^{s} w\right)^{2}+\varepsilon\left(\Lambda^{s} w_{x}\right)^{2}\right] d x \\
& \quad \leq 2 \delta\left(\left\|u_{\delta t}\right\|_{H^{s}}+\left\|u_{\delta x x t}\right\|_{H^{s}}+\left\|\Lambda^{s-2} w_{t}\right\|_{L^{2}}+\left\|\Lambda^{s-2} u_{\delta t}\right\|\right)\|w\|_{H^{s}}  \tag{3.47}\\
& \quad+c\left(\|w\|_{H^{s}}^{2}+\|w\|_{H^{q}}\|w\|_{H^{s}}\left\|u_{\delta}\right\|_{H^{s+1}}\right) \\
& \quad \leq c\left(\delta^{r_{1}}\|w\|_{H^{s}}+\|w\|_{H^{s}}^{2}\right)
\end{align*}
$$

where $\gamma_{1}=\min (1 / 4,(s-q-1) / 4)>0$. Integrating (3.47) with respect to $t$ leads to the estimate

$$
\begin{align*}
\frac{1}{2}\|w\|_{H^{s}}^{2} & \leq \int_{R}\left[(1-\varepsilon)\left(\Lambda^{s} w\right)^{2}+\varepsilon\left(\Lambda^{s} w_{x}\right)^{2}\right] d x  \tag{3.48}\\
& \leq \int_{R}\left[\left(\Lambda^{s} w_{0}\right)^{2}+\varepsilon\left(\Lambda^{s} w_{0 x}\right)^{2}\right] d x+c \int_{0}^{t}\left(\delta^{r_{1}}\|w\|_{H^{s}}+\|w\|_{H^{s}}^{2}\right) d \tau
\end{align*}
$$

It follows from the Gronwall inequality and (3.48) that

$$
\begin{align*}
\|w\|_{H^{s}} & \leq\left(2 \int_{R}\left[\left(\Lambda^{s} w_{0}\right)^{2}+\varepsilon\left(\Lambda^{s} w_{0 x}\right)^{2}\right] d x\right)^{1 / 2} e^{c t}+\delta^{r_{1}}\left(e^{c t}-1\right)  \tag{3.49}\\
& \leq c_{1}\left(\left\|w_{0}\right\|_{H^{s}}+\delta^{3 / 4}\right) e^{c t}+\delta^{r_{1}}\left(e^{c t}-1\right)
\end{align*}
$$

where $c_{1}$ is independent of $\varepsilon$ and $\delta$.
Then, (3.22) and the above inequality show that

$$
\begin{equation*}
\|w\|_{H^{s}} \longrightarrow 0 \quad \text { as } \varepsilon \longrightarrow 0, \quad \delta \longrightarrow 0 \tag{3.50}
\end{equation*}
$$

Next, we consider the convergence of the sequence $\left\{u_{\varepsilon t}\right\}$. Multiplying both sides of (3.29) by $\Lambda^{s-1} w_{t} \Lambda^{s-1}$ and integrating the resultant equation with respect to $x$, we obtain

$$
\begin{align*}
&(1-\varepsilon)\left\|w_{t}\right\|_{H^{s-1}}^{2}+\frac{1}{m+2} \int_{R}\left(\Lambda^{s-1} w_{t}\right) \Lambda^{s-1}\left(u_{\varepsilon}^{m+2}-u_{\delta}^{m+2}\right)_{x} d x \\
&+\int_{R}\left[-\varepsilon\left(\Lambda^{s-1} w_{t}\right)\right.\left.\left(\Lambda^{s-1} w_{x x t}\right)+(\delta-\varepsilon)\left(\Lambda^{s-1} w_{t}\right) \Lambda^{s-1}\left(u_{\delta t}+u_{\delta x x t}\right)\right] d x \\
&=\int_{R}\left(\Lambda^{s-1} w_{t}\right) \Lambda^{s-3}\left[-\varepsilon w_{t}+(\delta-\varepsilon) u_{\delta t}-\frac{k}{m+1} \partial_{x}\left(u_{\varepsilon}^{m+1}-u_{\delta}^{m+1}\right)-\partial_{x}\left(u_{\varepsilon}^{m+2}-u_{\delta}^{m+2}\right)\right. \\
& \quad \partial_{x}\left[\partial_{x}\left(u_{\varepsilon}^{m+1}\right) \partial_{x} w+\partial_{x}\left(u_{\varepsilon}^{m+1}-u_{\varepsilon}^{m+1}\right) \partial_{x} u_{\delta}\right] \\
&\left.+\left[u_{\varepsilon}^{m} u_{\varepsilon x} u_{\varepsilon x x}-u_{\delta}^{m} u_{\delta x} u_{\delta x x}\right]\right] d x+\lambda \int_{R} \Lambda^{s-1} w_{t} \Lambda^{s-1} w d x \tag{3.51}
\end{align*}
$$

It follows from inequalities (3.28) and the Schwartz inequality that there is a constant $c$ depending on $\tilde{T}$ and $m$ such that

$$
\begin{equation*}
(1-\varepsilon)\left\|w_{t}\right\|_{H^{s-1}}^{2} \leq c\left(\delta^{1 / 2}+\|w\|_{H^{s}}+\|w\|_{s-1}\right)\left\|w_{t}\right\|_{H^{s-1}}+\varepsilon\left\|w_{t}\right\|_{H^{s-1}}^{2} \tag{3.52}
\end{equation*}
$$

Hence

$$
\begin{align*}
\frac{1}{2}\left\|w_{t}\right\|_{H^{s-1}}^{2} & \leq(1-2 \varepsilon)\left\|w_{t}\right\|_{H^{s-1}}^{2}  \tag{3.53}\\
& \leq c\left(\delta^{1 / 2}+\|w\|_{H^{s}}+\|w\|_{H^{s-1}}\right)\left\|w_{t}\right\|_{H^{s-1}}
\end{align*}
$$

which results in

$$
\begin{equation*}
\frac{1}{2}\left\|w_{t}\right\|_{H^{s-1}} \leq c\left(\delta^{1 / 2}+\|w\|_{H^{s}}+\|w\|_{H^{s-1}}\right) \tag{3.54}
\end{equation*}
$$

It follows from (3.40) and (3.50) that $w_{t} \rightarrow 0$ as $\varepsilon, \delta \rightarrow 0$ in the $H^{s-1}$ norm. This implies that $u_{\varepsilon}$ is a Cauchy sequence in the spaces $C\left([0, T) ; H^{s}(R)\right)$ and $C\left([0, T) ; H^{s-1}(R)\right)$, respectively. The proof is completed.

## 4. Proof of the Main Result

We consider the problem

$$
\begin{gather*}
(1-\varepsilon) u_{t}-\varepsilon u_{t x x}=\left(1-\partial_{x}^{2}\right)^{-1}\left[-\frac{k}{m+1}\left(u^{m+1}\right)_{x}-\frac{m+3}{m+2}\left(u^{m+2}\right)_{x}+\frac{1}{m+2} \partial_{x}^{3}\left(u^{m+2}\right)\right. \\
\left.-(m+1) \partial_{x}\left(u^{m} u_{x}^{2}\right)+u^{m} u_{x} u_{x x}\right]+\lambda u  \tag{4.1}\\
u(0, x)=u_{\varepsilon 0}(x) .
\end{gather*}
$$

Letting $u(t, x)$ be the limit of the sequence $u_{\varepsilon}$ and taking the limit in problem (4.1) as $\varepsilon \rightarrow 0$, from Lemma 3.13, we know that $u$ is a solution of the problem

$$
\begin{gather*}
u_{t}=\left(1-\partial_{x}^{2}\right)^{-1}\left[-\frac{k}{m+1}\left(u^{m+1}\right)_{x}-\frac{m+3}{m+2}\left(u^{m+2}\right)_{x}+\frac{1}{m+2} \partial_{x}^{3}\left(u^{m+2}\right)\right. \\
\left.-(m+1) \partial_{x}\left(u^{m} u_{x}^{2}\right)+u^{m} u_{x} u_{x x}\right]+\lambda u  \tag{4.2}\\
u(0, x)=u_{0}(x),
\end{gather*}
$$

and hence $u$ is a solution of problem (4.2) in the sense of distribution. In particular, if $s \geq 4$, $u$ is also a classical solution. Let $u$ and $v$ be two solutions of (4.2) corresponding to the same
initial value $u_{0}$ such that $u, v \in C\left([0, T) ; H^{s}(R)\right)$. Then, $w=u-v$ satisfies the Cauchy problem

$$
\begin{gather*}
w_{t}=\left(1-\partial_{x}^{2}\right)^{-1}\left\{\partial _ { x } \left[-\frac{k}{m+1} w g_{m}-\frac{m+3}{m+2} w g_{m+1}+\frac{1}{m+2} \partial_{x}^{2}\left(w g_{m+1}\right)\right.\right. \\
\left.\left.-\partial_{x}\left(u^{m+1}\right) \partial_{x} w-\partial_{x}\left(u^{m+1}-v^{m+1}\right) \partial_{x} v\right]+u^{m} u_{x} u_{x x}-v^{m} v_{x} v_{x x}\right\}+\lambda w \\
w(0, x)=0 \tag{4.3}
\end{gather*}
$$

For any $1 / 2<q<\min \{1, s-1\}$, applying the operator $\Lambda^{q} w \Lambda^{q}$ to both sides of (4.3) and integrating the resultant equation with respect to $x$, we obtain the equality

$$
\begin{gather*}
\frac{1}{2} \frac{d}{d t}\|w\|_{H^{q}}^{2}=\int_{R}\left(\Lambda^{q} w\right) \Lambda^{q-2}\left\{\partial _ { x } \left[-\frac{k}{m+1} w g_{m}-\frac{m+3}{m+2} w g_{m+1}+\frac{1}{m+2} \partial_{x}^{2}\left(w g_{m+1}\right)\right.\right. \\
\left.-\partial_{x}\left(u^{m+1}\right) \partial_{x} w-\partial_{x}\left(u^{m+1}-v^{m+1}\right) \partial_{x} v\right]+u^{m} u_{x} u_{x x}  \tag{4.4}\\
\left.-v^{m} v_{x} v_{x x}\right\} d x+|\lambda|\|w\|_{H^{q}}^{2}
\end{gather*}
$$

By the similar estimates presented in Lemma 3.13, we have

$$
\begin{equation*}
\frac{d}{d t}\|w\|_{H^{q}}^{2} \leq \tilde{c}\|w\|_{H^{q}}^{2} \tag{4.5}
\end{equation*}
$$

Using the Gronwall inequality leads to the conclusion that

$$
\begin{equation*}
\|w\|_{H^{q}} \leq 0 \times e^{\tilde{c} t}=0 \tag{4.6}
\end{equation*}
$$

for $t \in[0, \widetilde{T})$. This completes the proof.

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