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Integrality of the Monopole Number in SU(2) Yang-Mills-Higgs Theory on $\mathbb{R}^3 \star$

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Abstract. We prove that in classical SU(2) Yang-Mills-Higgs theories on \mathbb{R}^3 with a Higgs field in the adjoint representation, an integer-valued monopole number (magnetic charge) is canonically defined for any finite-action $L^2_{1, \text{loc}}$ configuration. In particular the result is true for smooth configurations. The monopole number is shown to decompose the configuration space into path components.

Introduction

In classical Yang-Mills-Higgs theories over a Riemannian manifold M, one fixes a principal bundle

P $\downarrow^{_{G}}$ M

and studies the action functional

$$a(A, \Phi) = \frac{1}{2} \int_{M} (|F_A|^2 + |d_A \Phi|^2).$$
(1)

This functional is defined on the configuration space

$$\mathscr{C} = \{ (A, \Phi) \in \mathscr{A} \times \mathscr{E} | a(A, \Phi) < \infty \},$$
(2)

where \mathscr{A} is the space of $L^2_{1, \text{loc}}$ connections on *P*, and \mathscr{E} is the space of $L^2_{1, \text{loc}}$ sections of some (fixed) associated vector bundle *E*. (This is the most general configuration space. Often one considers only *smooth* connections *A* and sections Φ , but for many applications this restriction is inconvenient. If a potential term is included in the action, the configuration space again is smaller.) The symbols $F_A, d_A \Phi$ denote the curvature of $A \in \mathscr{A}$ and the covariant derivative of $\Phi \in \mathscr{E}$, respectively; norms

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