

Multicomponent Field Theories and Classical Rotators

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Abstract. It is shown that a D -component Euclidean quantum field, $\varphi=(\varphi^1, \dots, \varphi^D)$, with $\lambda|\varphi|^4 + \beta|\varphi|^2$ interaction, can be obtained as a limit of (ferromagnetic) classical rotator models; this extends a result of Simon and Griffiths from the case $D=1$. For these Euclidean field models, it is then shown that a Lee-Yang theorem applies for $D=2$ or 3 and that Griffiths' second inequality is valid for $D=2$; a complete proof is included of a Lee-Yang theorem for plane rotator and classical Heisenberg models. As an application of Griffiths' second inequality for $D=2$, an interesting relation between the "parallel" and "transverse" two-point correlations is obtained.

1. Introduction

Consider a multicomponent scalar field, $\varphi(x, t)=(\varphi^1(x, t), \dots, \varphi^D(x, t))$, in d -dimensional space-time with Hamiltonian,

$$\int_{\mathbb{R}^d-1} [\frac{1}{2} \sum_{i=1}^D ((\pi^i)^2 + |\nabla \varphi^i|^2 + m_0^2 (\varphi^i)^2) + \tilde{Q}(\sum_{i=1}^D (\varphi^i)^2) + \sum_{i=1}^D A^i \varphi^i] dx, \quad (1.1)$$

where $\pi = \partial \varphi / \partial t$, \tilde{Q} is a polynomial with positive highest coefficient, and $A(x, t)$ is an external field interacting with φ . The problem of constructing a corresponding quantum field has been considerably simplified in recent years by the probabilistic methods of Euclidean field theory (see, for example, the articles in [1]); it has been particularly realized that the associated Euclidean field is closely related (via a "lattice approximation") to certain models of ferromagnets from classical statistical mechanics [2].

The type of model we are concerned with consists of a family of random D -dimensional "spin" vectors $\{S_j = (S_j^1, \dots, S_j^D): j=1, \dots, N\}$ with joint probability distribution on $(\mathbb{R}^D)^N$,

$$\frac{1}{Z} \exp(\sum_{j=1}^N \mathbf{a}_j \cdot \mathbf{s}_j + \sum_{j,k=1}^N \sum_{i=1}^D J_{jk}^i s_j^i s_k^i) \prod_{j=1}^N \varrho_j(\mathbf{s}_j), \quad (1.2)$$

where

$$Z = Z(\{\mathbf{a}_j\}, \{J_{jk}^i\}) = \int_{(\mathbb{R}^D)^N} \exp(\sum_j \mathbf{a}_j \cdot \mathbf{s}_j + \sum_{j,k,i} J_{jk}^i s_j^i s_k^i) \prod_j d\varrho_j(\mathbf{s}_j), \quad (1.3)$$

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