NEW CONNECTION METHOD ACROSS MORE GENERAL TURNING POINTS¹

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The basic oscillator or Schroedinger equation

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$${}^{2}d^{2}w/dz^{2} + q^{2}w(z) = 0, \quad \epsilon \downarrow 0,$$
(1)

with q(z) analytic, except for a root or singular point at z = 0, poses a wellknown connection problem. *WKBJ* or *L*-*G* approximations

$$w \sim Aq^{-\frac{1}{2}}\exp(i\xi/\epsilon) + Bq^{-\frac{1}{2}}\exp(-i\xi/\epsilon),$$

$$dw/z \sim (iq^{\frac{1}{2}}/\epsilon)[A \exp(i\xi/\epsilon) - B \exp(-i\xi/\epsilon)]$$
(2)

with

$$\xi = \int_0^z q(s) \ ds \tag{3}$$

and constants A, B are valid in suitable regions, but not in a neighborhood of the transition point. The constants A_p , B_l in (2) valid to the left of such a point generally differ from those, A_p , B_p , valid to the right. Langer [1] found one constant pair in terms of the other for the class of "fractional turning points" at which $z^{-\nu}q(z)$ is analytic for real $\nu > -1$. But, if branch points, even of infinite order, be thus admitted, why not logarithms of z and still other branch points?

The reason lies in his "central connection" method by uniform approximation even near z = 0. For fractional turning points such approximands to w(z)are furnished by Bessel functions, but for logarithmic turning points they are necessarily less tractable. Uniform approximation, however, is not a prerequisite for connection, nor is it needed for the most important physical applications, e.g., in scattering.

A new method will now be summarized which solves the problem for a more general class of equations (1) characterized (a precise statement is found in the Appendix) by

$$\rho(\xi) := \frac{1}{2}q^{-2}dq/dz = \xi^{-1}(\gamma + o(1)) \text{ as } \xi \to 0$$
(4)

with real $\gamma < \frac{1}{2}$. This includes the fractional turning points, where $\gamma = \frac{1}{2}\nu/(1 + \nu)$ (< $\frac{1}{2}$ since $\nu > -1$), and also logarithmic turning points, where $q(z) \sim z^{\nu}(\log z)^{\mu}$

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