# A NEW PROOF AND AN EXTENSION OF HARTOG'S THEOREM ${ }^{1}$ 

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Let $R$ denote $n$ dimensional real euclidean space and let $\Omega_{0}$ be a shell in $R$, by this we mean that there exist open sets $\Omega_{1}, \Omega_{2}$ where $\Omega_{1}$ is relatively compact and has its closure contained in $\Omega_{2}$, and $\Omega_{0}$ $=\Omega_{2}-$ closure $\Omega_{1}$. Call $\Gamma_{j}$ the boundary of $\Omega_{j}$. Let $D=\left(D_{1}, \cdots, D_{r}\right)$ be a sequence of linear partial differential operators with constant coefficients on $R$ with $r>1$. For a function $f$ on $R$ we write $D f=0$ if $D_{j} f=0$ for $j=1,2, \cdots, r$. We want to determine the conditions on $D$ in order that the following property should hold: If $f$ is an indefinitely differentiable function on $\Omega_{0}$ with $D f=0$ then there exists a unique indefinitely differentiable function $h$ on $\Omega_{2}$ with $D h=0$ and $h=f$ on $\Omega_{0}$. Hartog's theorem asserts that such an extension of $f$ is possible if $R$ is complex euclidean space of complex dimension $n / 2=m>1$ and $\Omega_{1}$ and $\Omega_{2}$ are topological balls, and $D_{j}=\partial / \partial x_{2 j-1}+i \partial / \partial x_{j}$ for $j=1,2, \cdots, m$ where $x=\left(x_{1}, \cdots, x_{n}\right)$ are the coordinates on $R$. An extension of Hartog's theorem has been found by S. Bochner in [1] by a different method.

We can find a function $g$ defined and $C^{\infty}$ on $\Omega_{2}$ such that $g=f$ on $\Omega_{0}$ except on an arbitrarily small neighborhood $N\left(\Gamma_{1}\right)$ in $\Omega_{0}$. (We choose $N\left(\Gamma_{1}\right)$ so small that its closure does not meet $\Gamma_{2}$ ) Call $\Omega_{3}=\Omega_{1} \cup N\left(\Gamma_{1}\right)$. We have $D g=0$ on $\Omega-\Omega_{3}$. We set $g_{j}=D_{j} g$, so $g_{j}$ are $C^{\infty}$ and have their supports in the closure of $\Omega_{3}$; in particular the $g_{j}$ are of compact support. For any $j, k$,

$$
\begin{equation*}
D_{k} g_{j}=D_{j} g_{k} \tag{1}
\end{equation*}
$$

since both sides are equal to $D_{k} D_{j} g$ in $\Omega_{3}$ and zero outside.
Next we take the Fourier transforms: Call $P_{k}$ the Fourier transform of $D_{k}$ and $G_{k}$ that of $g_{k} ; P_{k}$ is a polynomial and $G_{k}$ an entire function of exponential type on $C$ (complex $n$-space); the exponential type of $G_{k}$ is determined by the convex hull $K$ of $\Omega_{3}$. Moreover, $G_{k}$ decreases on the real part of $C$ faster than the reciprocal of any polynomial (see [5]). Relation (1) becomes

$$
\begin{equation*}
P_{k}(z) G_{j}(z)=P_{j}(z) G_{k}(z) \tag{2}
\end{equation*}
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[^0]:    ${ }^{1}$ Work supported by ONR 432 JLP.

