

101. *Calculus on Gaussian White Noise. III*

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In the previous parts of this series [11], [12], we have given a systematic treatment of calculus on Gaussian white noise, which is a reformulation of Hida's works [1], [2]. In this part we will show further relations between Hida's approach and ours. We will use the same notations and definitions as in Part I and Part II.

**§ 8. Multiple Wiener integrals.** Here we assume that the Borel measure  $\nu$  on  $T$  has no atoms. Let  $\mathcal{C} \subset E_0 = L^2(T, \nu) \subset \mathcal{C}^*$  be a triplet as in § 5 of Part II, and let  $\mu$  be the measure of Gaussian white noise on  $\mathcal{C}^*$  with characteristic functional  $\exp[-\|\xi\|_0^2/2]$ . The multiple Wiener integral  $I_n(F_n)$  of  $F_n$  in  $L^2(T^n, \nu^n)$  is defined as follows:

First,  $I_1(F_1)$  is the limit of  $\langle x, \xi_k \rangle$  in  $(L^2) = L^2(\mathcal{C}^*, \mu)$ , where  $\{\xi_k\}$  is any sequence in  $\mathcal{C}$  with  $\|\xi_k - F_1\|_0 \rightarrow 0$ , as  $k \rightarrow \infty$ . Specially, put  $W(B) \equiv I_1(I_B)$ , where  $I_B$  denotes the indicator function of a Borel set  $B$  with  $\nu(B) < \infty$ . Secondary, let  $\alpha = \{B_j\}$  be a countable Borel partition of  $T$  with  $\nu(B_j) < \infty$  and let  $\alpha^n$  be the collection of all subsets of  $T^n$  of the form  $C = B_{j(1)} \times B_{j(2)} \times \cdots \times B_{j(n)}$ ,  $B_{j(k)} \in \alpha$ ,  $B_{j(k)} \cap B_{j(m)} = \emptyset$  for  $k \neq m$ . For such a set  $C$  in  $\alpha^n$ , define

$$I_n(I_C) \equiv \prod_{k=1}^n W(B_{j(k)}).$$

Define  $I_n(G_n) \equiv \sum a_k I_n(I_{C_k})$  for  $G_n = \sum a_k I_{C_k}$  with  $C_k \in \alpha^n$ . Then we can define  $I_n(F_n)$  by

$$(8.1) \quad I_n(F_n) \equiv \text{l.i.m.}_{\alpha \uparrow} I_n(F_n^\alpha), \quad F_n^\alpha \equiv \sum \nu^{-1}(C)(F_n, I_C)I_C,$$

where  $\alpha \uparrow$  means refinements.

**Theorem 8.1.** (i) For  $F_n \in L^2(T^n, \nu^n)$ , put  $\varphi(x) = I_n(F_n)$ , then we have

$$(\mathcal{S}\varphi)(\xi) = \int_{T^n} F_n(u_1, \dots, u_n) \xi(u_1) \cdots \xi(u_n) d\nu^n(u_1, \dots, u_n).$$

(ii) For any  $\psi$  in  $(L^2)$ , there exist  $F_n \in \hat{L}^2(T^n, \nu^n)$ ,  $n \geq 0$ , such that  $\psi(x)$  is decomposed into the following orthogonal sum;

$$\psi(x) = \sum_{n=0}^{\infty} I_n(F_n) \quad \text{and} \quad \|\psi\|_{(L^2)}^2 = \sum_{n=0}^{\infty} n! \|F_n\|_{L^2(T^n, \nu^n)}^2.$$

We now remark that the symmetric  $L^2$ -space  $\hat{L}^2(T^n, \nu^n)$  is naturally identified with the symmetric tensor product space  $E_0^{\hat{\otimes} n}$ . By Theorems 6.3 and 6.5, we have