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## 87. Infinitely Divisible Distributions and Ordinary Differential Equations

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1. Introduction. In the study of the limit distributions of multi-type Galton-Watson processes, S. Sugitani [2] has discovered that if a nonnegative function  $\psi(t, \lambda)$ , defined for  $t \ge 0$ ,  $\lambda \ge 0$ , satisfies the ordinary differential equation having  $\lambda$  as a parameter

(1)  $\psi' = -B\psi^2 + \lambda p(t), \qquad \psi(0, \lambda) = m\lambda,$ 

where B>0,  $m\geq 0$  and p(t) is a polynomial with positive coefficients, then there exists for each t>0 an infinitely divisible distribution  $\nu_t$  on  $[0, \infty)$  such that

(2) 
$$\exp\left\{-\int_0^t \psi(s,\lambda)ds\right\} = \int_0^\infty e^{-\lambda x} \nu_t(dx).$$

Further information on  $\nu_t$  is given in [3].

In this note we will prove a stronger result that  $\exp\{-\psi(t,\lambda)\}$  is the Laplace transform of some infinitely divisible distribution  $\mu_t$  on  $[0, \infty)$ . Our proof is quite elementary and can be applied to more general equations.

2. A heuristic argument. Given f(x),  $g(t, \lambda)$  and  $h(\lambda)$  defined for  $x \in (-\infty, \infty)$ ,  $t \in [0, T]$ ,  $\lambda \in [0, \infty)$ , consider the following ordinary differential equation having  $\lambda$  as a parameter;

(3)  $\psi' = f(\psi) + g(t, \lambda), \qquad \psi(0, \lambda) = h(\lambda).$ 

For the moment, we assume that equation (3) has a unique solution  $\psi(t, \lambda)$  in  $[0, T] \times [0, \infty)$ . Here and after we will write  $\psi'$  for  $D_t \psi(t, \lambda)$ ,  $f^{(n)}$  for  $D_x^n f$ ,  $g_{n\lambda}(t, \lambda)$  for  $D_\lambda^n g(t, \lambda)$  and so on. We now seek a suitable condition in order that  $\psi_{\lambda}(t, \lambda)$  is completely monotonic in  $\lambda \in (0, \infty)$  for each  $t \ge 0$ . The essential part of our condition is that  $-f^{(2)}(\cdot), g_{\lambda}(t, \cdot)$  for each  $t \in [0, T]$  and  $h_{\lambda}(\cdot)$  are completely monotonic in  $(0, \infty)$ . To show the above assertion, differentiating (3) with respect to  $\lambda$ , we have

$$\psi_{\lambda}' = f^{(1)}(\psi)\psi_{\lambda} + g_{\lambda}(t,\lambda), \qquad \psi_{\lambda}(0,\lambda) = h_{\lambda}(\lambda).$$

Since  $g_{\lambda}(t, \lambda) \ge 0$  and  $h_{\lambda}(\lambda) \ge 0$ , it follows that  $\psi_{\lambda}(t, \lambda) \ge 0$ . Similarly, *n*-times differentiation of (3) leads us to

$$\psi_{n\lambda} = f^{(1)}(\psi)\psi_{n\lambda} + f^{(2)}(\psi) \sum_{\substack{k_1+k_2=n\\1\le k_1\le k_2}} c_{k_1,k_2} \psi_{k_1\lambda} \psi_{k_2\lambda} \\ + \dots + f^{(j)}(\psi) \sum_{\substack{k_1+\dots+k_j=n\\1\le k_1\le \dots\le k_j}} c_{k_1,\dots,k_j} \psi_{k_1\lambda} \dots \psi_{k_j\lambda}$$