# 6. On the Intersection Number of the Path of a Diffusion and Chains 

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1. We are concerned with the following problem which was already considered by H. P. McKean [4] for the Brownian motion: in what manner does the path of a diffusion on a manifold wind around a fixed point or a hole asymptotically? For this purpose, we shall define a stochastic version of the intersection number. As is wellknown, the usual intersection number can be represented by the integral of a differential double 1-form with singularity ([1]). Although the path of the diffusion is not smooth, we can define its intersection number with a chain by using the integral of the 1-form along the path defined in [2] (see also [3]). We then study the asymptotic behaviors of such random intersection numbers to get some solutions of the above mentioned problem.
2. Let $M$ be a $d$-dimensional connected orientable Riemannian manifold with a Riemannian metric $g$ and $\Delta$ be the Laplace-Beltrami operator corresponding to $g$. Let $L=\Delta / 2+b$, where $b$ is a $C^{\infty}$ vector field on $M$. Consider the minimal diffusion process $X=\left(X_{t}, P_{x}\right)$ on $M$ corresponding to $L$. For any continuous mapping $c:[0, t] \rightarrow M$, we denote by $c[0, t]$ the curve determined by $c: c[0, t]=\{c(s) ; 0 \leqq s \leqq t\}$. We regard $c[0, t]$ as a singular 1-chain ([5]).

To define the intersection number, we prepare some notations. We principally use the notations of de Rham's book ([1]). Let $\bar{D}$ be the space of square integrable currents. Set $\overline{\mathscr{D}}_{1}=\{T \in \overline{\mathscr{D}} ; T$ is homologous to zero $\}, \overline{\mathcal{D}}_{2}=\{T \in \overline{\mathscr{D}} ; T$ is cohomologous to zero $\}$ and $\mathscr{D}_{3}=\{T \in \overline{\mathscr{D}} ; T$ is harmonic $\}$. Then $\overline{\mathscr{D}}=\overline{\mathscr{D}}_{1}+\overline{\mathscr{D}}_{2}+\mathscr{D}_{3}$. Let $H_{1}, H_{2}, H_{3}$ be the projections on $\overline{\mathscr{D}}_{1}, \overline{\mathscr{D}}_{2}, \mathscr{D}_{3}$ respectively. For any 1-current $T$ which is continuous in mean at infinity, we define $H_{i} T$ by $\left(H_{i} T, \phi\right)=\left(T, H_{i} \phi\right), \phi \in C^{\infty} \cap \overline{\mathscr{D}}$, $i=1,2,3$. Then $T$ can be decomposed uniquely as follows: $T=H_{1} T$ $+H_{2} T+H_{3} T$. Denote by $h_{i}(x, y)$ the kernel of $H_{i}, i=1,2,3$. Let $e(x, y)$ $=*_{y} h_{1}(x, y)$ be the adjoint form of $h_{1}$ (as 1-form of $y$ ). Then $e$ is $C^{\infty}$ if $x \neq y$. It is known that $e(x, y)$ can be written locally as follows. Let $\Delta$ be the Hodge-Kodaira's Laplacian acting on 1-forms. We can choose a domain $U$ on which a fundamental solution $\gamma(x, y)$ for $\Delta \alpha=\beta$ exists. Let $\sigma(x, y)$ be a $C^{\infty}$ function supported in $U \times U$ with (i) $0 \leqq \sigma \leqq 1$, (ii)

