# 66. Studies on Holonomic Quantum Fields. IX 

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In this note we shall give a symplectic version of the 2-dimensional operator theory, previously expounded in the orthogonal case [2], [5], [6]. Of particular interest is the neutral theory discussed in § 4. Corresponding to the bose field $\varphi^{F}(a)$ [1], there arises a strongly interacting fermi field $\varphi^{B}(\alpha)={ }^{t}\left(\varphi_{+}^{B}(\alpha), \varphi_{-}^{B}(\alpha)\right)$. These two fields $\varphi^{F}(a)$ and $\varphi^{B}(a)$ are shown to share the same $S$-matrix in common, and their $\tau$-functions are related to each other through simple formulas (34), (36), (38)-(39) (cf. IV-(49) [2]).

We remark that the 1-dimensional Riemann-Hilbert problem [4], [8] is also dealt with in the symplectic framework.

We follow the notations used throughout this series [1]-[6].
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1. Let $W$ be an $N$-dimensional complex vector space equipped with a skew-symmetric inner product $\langle$,$\rangle . Let A(W)$ be the algebra generated by $W$ with the defining relation $w w^{\prime}-w^{\prime} w=\left\langle w, w^{\prime}\right\rangle$. Denote by $S(W)$ the symmetric tensor algebra over $W$. As in the orthogonal case [3], [7], the norm map

$$
\begin{equation*}
\mathrm{Nr}: A(W) \xrightarrow{\sim} S(W) \tag{1}
\end{equation*}
$$

and the expectation value $\langle a\rangle$ of $a \in A(W)$ are defined analogously, by specifying a bilinear form $\left(w, w^{\prime}\right) \rightarrow\left\langle w w^{\prime}\right\rangle$ on $W$ such that $\left\langle w w^{\prime}\right\rangle$ $-\left\langle w^{\prime} w\right\rangle=\left\langle w, w^{\prime}\right\rangle\left(w, w^{\prime} \in W\right)$.

Now let $v_{1}, \cdots, v_{N}$ be a basis of $W$, and set $K=\left(\left\langle v_{\mu} v_{\nu}\right\rangle\right), H=\left(\left\langle v_{\mu}\right.\right.$, $\left.\left.v_{\nu}\right\rangle\right)=K-{ }^{t} K$. Consider an element $g$ of the form
(2) $\quad \operatorname{Nr}(g)=\langle g\rangle e^{\rho / 2}, \quad \rho=\sum_{\mu, \nu=1}^{N} R_{\mu \nu} v_{\mu} v_{\nu}=v R^{t} v$
with $v=\left(v_{1}, \cdots, v_{N}\right)$. Contrary to the orthogonal case, $e^{\rho / 2}$ no longer belongs to $S(W)$. So we let $R_{\mu \nu}=R_{\nu \mu} \in t \cdot C[[t]]$, and regard $g$ (resp. $e^{\rho / 2}$ ) as an element of $A(W)[[t]]$ (resp. $S(W)[[t]])$, the formal power series ring with coefficients in $A(W)$ (resp. $S(W)$ ). The norm map (1) is uniquely extended there. (This formulation is due to T. Miwa.) Most of the formulas in the orthogonal case are valid for $g$ of the form (2), if we replace ${ }^{t} K$ by $-{ }^{t} K$. We tabulate below formulas corresponding to (1.5.5)-(1.5.6), (1.5.7)-(1.5.8) and (1.4.6)-(1.4.7) of [7].

$$
\begin{equation*}
\mathrm{Nr}(w g)=\left(\sum_{\mu, \nu=1}^{N} v_{\mu}\left(1+R^{t} K\right)_{\mu \nu} c_{\nu}\right) \cdot\langle g\rangle e^{\rho / 2} \tag{3}
\end{equation*}
$$

