# 3. Studies on Holonomic Quantum Fields. I 

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To understanding the mathematical structure of quantized fields or systems with infinite freedom, non trivial but exactly calculable models would be of great help [1]. In this and subsequent notes we present, both in the continuum and in the lattice, 2-dimensional soluble models of neutral scalar massive field theory whose $\tau$-functions exhibit a non trivial singularity structure.

In the present article we deal with the continuum case. We introduce an auxiliary free fermi/bose field and construct the field operator by specifying its induced rotation in the space of wave functions. Making use of the "theory of rotation" (2 cf. [2]) developed recently by the first author, we express this field operator in the normal product form of these free fields. We also calculate the asymptotic fields and the $S$-matrix of the field $\varphi^{F}$ defined in 3 . Next we give explicit formulae for $\tau$-functions of these models and study their holonomy structure.

The lattice field theory will be discussed in a subsequent paper. Specifically we shall show that our model $\varphi^{F} / \varphi_{F}$ coincide with the scaling limit of the Ising model from above/below the critical temperature. Main part of these results has been announced in [3].

We use the following notations. The space-time and the energymomentum co-ordinates are denoted by $x=\left(x^{0}, x^{1}\right)$ and $p=\left(p^{0}, p^{1}\right)$. We also use $x^{ \pm}=\left(x^{0} \pm x^{1}\right) / 2$ and $p^{ \pm}=p^{0} \pm p^{1}$. The mass-shell $\left\{p \in \boldsymbol{R}^{2} \mid p^{2}=\left(p^{0}\right)^{2}\right.$ $\left.-\left(p^{1}\right)^{2}=m^{2}\right\}(m>0)$ is denoted by $M$. For $p \in M$ we set $u^{ \pm 1}=p^{ \pm} / m$, $\underline{d u}=d u / 2 \pi|u|$.

1. Let $\psi(u)^{\dagger}$ and $\psi(u)(u>0)$ be the creation and annihilation operators of auxiliary fermion. If we define $\psi(u)=\psi(-u)^{\dagger}$ for $u<0$, their anti-commutation relation reads $\left[\psi(u), \psi\left(u^{\prime}\right)\right]_{+}=2 \pi|u| \delta\left(u+u^{\prime}\right)$. Likewise we define auxiliary bosons $\phi(u)$ with the commutation relation $\left[\phi(u), \phi\left(u^{\prime}\right)\right]_{-}=2 \pi u \delta\left(u+u^{\prime}\right)$. In two dimensional space-time these two are in fact equivalent. Namely

$$
\begin{equation*}
\phi_{ \pm}(u)=: \psi(u) \exp \int_{0}^{\infty}(-2) \theta\left( \pm\left(|u|-u^{\prime}\right)\right) \psi\left(u^{\prime}\right)^{\dagger} \psi\left(u^{\prime}\right) \underline{d u^{\prime}}: \tag{1}
\end{equation*}
$$

satisfy the commutation relation $\left[\phi_{ \pm}(u), \phi_{ \pm}\left(u^{\prime}\right)\right]_{-}=2 \pi u \delta\left(u+u^{\prime}\right)$, and conversely $\psi(u)$ is given by

$$
\begin{equation*}
\psi(u)=: \phi_{ \pm}(u) \exp \int_{0}^{\infty}(-2) \theta\left( \pm\left(|u|-u^{\prime}\right)\right) \phi_{ \pm}\left(u^{\prime}\right)^{\dagger} \phi_{ \pm}\left(u^{\prime}\right) \underline{d u^{\prime}}: . \tag{2}
\end{equation*}
$$

