## 216. Neutron Transport Process on Bounded Homogeneous Domain

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(Comm. by Kinjirô KUNUGI, M. J. A., Nov. 12, 1970)

1. The neutron transport process has been studied by Harris ([1]) and Mullikin ([5]) as an application of the theory of discrete-time branching processes. The main problems are the asymptotic behavior of the number of neutrons, the extinction probability and the rate of convergence of the extinction probability at time t to the extinction probability. In this paper we consider similar problems for a monoenergetic and isotropic neutron transport process on a bounded homogeneous domain. We will formulate the model as a continuous-time branching process and apply the general theory of such processes ([2]). Main results are the theorems  $1 \sim 5$  below. It will be seen that the expected number of new-born neutrons plays an essential role in the above problems. This is a typical property of branching processes, which is well known for Galton-Watson processes.

2. Let D be a bounded closed convex domain in the three-dimensional Euclidian space  $\mathbf{R}^3$  with a smooth boundary and  $\Omega$  be the unit sphere in  $\mathbf{R}^3$ . We denote by G the product space  $D \times \Omega$  and  $\partial G$  the set  $(x, \omega)$  where x belongs to the boundary of D and  $\omega$  is a direction exiting the domain; i.e.,  $(\omega, n_x) \ge 0$  where  $n_x$  is the direction of the outernormal at x. We formulate our model of neutron transport process as a continuous-time branching process as follows; a particle at  $x \in D$ starting with unit speed in the direction  $\omega^{(*)}$  will, at a random time T which is exponentially distributed with mean  $\sigma^{-1}$ , be absorbed, scattered, or multiplied by fission. If it leaves the domain D before T, then it is absorbed. The direction of new particles is supposed to be isotropically distributed. Each of new particles, independently each other, performs a similar motion as the original one. We can construct such a branching process on a suitable probability space ([2]) and every probabilistic argument below is based on this process.

Let  $F[\xi] = \sum_{n=0}^{\infty} p_n \xi^n$  where  $p_n$  is the probability that *n* neutrons are produced when fission occurs. (In particular  $p_0$  is the probability of absorption and  $p_1$  the probability of scattering.) We will assume  $F'[1] < \infty$  and  $p_0 + p_1 < 1$ . The first assumption guarantees that the

<sup>\*)</sup> This statement will be simplified below as "starting at  $(x, \omega)$ ."