# 145. Analogue of Fourier's Method for Korteweg - de Vries Equation 

By Shunichi Tanaka<br>Department of Mathematics, Osaka University<br>(Comm. by Kôsaku Yosida, m. J. A., Nov. 13, 1972)

1. Introduction. In this paper we study the Kortewegde Vries (KdV) equation
(1) $\quad u_{t}-6 u u_{x}+u_{x x x}=0 \quad u=u(t)=u(x, t)-\infty<x, t<\infty$
for rapidly decreasing initial data. Gardner, Greene, Kruskal and Miura (G.G.K.M.) [2] have associated one dimensional Schrödinger operators $L_{u(t)}=-(d / d x)^{2}+u(t)$ to a solution of (1). They have found a simple formula describing the time variation of scattering data of $L_{u(t)}$. This paper is concerned with converse statement which may be viewed as a non-linear analogue of Fourier's method for solving linear partial differential equations of mathematical physics: Given the initial value one determine the scattering data of $L_{u(0)}$. Define scattering data for each $t$ according to the formula of G.G.K.M. Using inverse scattering theory, one can construct potential $u(x, t)$ with prescribed scattering data for each $t$. Then $u(x, t)$ is a solution of (1).

Throughout the paper subscripts with independent variables denote partial differentiations. Integrations are taken over ( $-\infty, \infty$ ) unless explicitly indicated.
2. Preparation from scattering theory. Consider one dimensional Schrödinger equation
( 2 ) $-\phi_{x x}+u(x) \phi=\zeta^{2} \phi$.
Under the assumption that $(1+|x|) u(x)$ is integrable, the inverse scattering theory for (2) has been solved by Marchenko for the half line ( $0, \infty$ ) and then the case of the infinite interval has been treated by Faddeev [1]. We follow [1] in this paper.

For each $\zeta=\xi+i \eta, \eta \geq 0$, there exist unique solutions $f_{ \pm}(x, \zeta)$ which behave like $\exp ( \pm i \zeta x)$ as $x \rightarrow \pm \infty$. They are called Jost solutions of (2). Jost solutions are analytic in $\zeta$, $\operatorname{Im} \zeta>0$. If $\zeta=\xi$ non-zero real, then $f_{+}$and its complex conjugate $f_{+}^{*}$ are independent solutions of (2). One can express $f_{-}$as $f_{-}=a(\xi) f_{+}^{*}+b(\xi) f_{+} . a(\xi)$ is limiting value of a function $\alpha(\zeta)$ analytic in $\zeta$, $\operatorname{Im} \zeta>0$. The (right) reflection coefficient $r(\xi)=b(\xi) a(\xi)^{-1}$ is defined for $\xi \neq 0$ and its absolute value is bounded by 1. $a(\zeta)$ has only a finite numbers of zeros. They are all simple and purely imaginary. We denote them by $i \eta_{1}, \cdots, i \eta_{N} . f_{ \pm}$are linearly dependent for $\zeta=i \eta_{j}$ and are square integrable because of the asymptotic

