PAPERS COMMUNICATED

72. Asymptotic Expansions in the Heaviside's Operational Calculus.

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1. Let p denote d/dx and F(z) be a power series $\sum a_n z^n$. Operating $F(p)p^{1/2}$ on 1 we have formally

$$\frac{1}{\sqrt{\pi x}}\left(a_0 - \frac{a_1}{(2x)} + \frac{1 \cdot 3 a_2}{(2x)^2} + \dots (-1)^n \frac{1 \cdot 3 \cdot 5 \dots (2n-1) a_n}{(2x)^n} + \dots\right), \quad (1)$$

which is in general divergent. Denoting this series by S(x) and putting

$$F(p)p^{1/2}$$
. $1=h(x)$,

we have

$$h(x) \sim S(x)$$
,

where \sim means that S(x) is an asymptotic expansion of h(x). This is one of the Heaviside's rules.

To give a rigorous basis for this Heaviside's rule, Mr. Carson introduced Laplace operator and defined h(x) by the relation

$$\frac{F(p)}{p^{1/2}} = \int_0^\infty e^{-pt} h(t) \, dt \tag{2}$$

and proceeded as follows.1)

He defined g(t) by

$$F(p) = \int_0^\infty e^{-pt} g(t) dt$$

and deduced

$$a_n = \frac{(-1)^n}{n!} \int_0^\infty g(t)t^n dt$$
 $(n=0, 1, 2, ...)$

by assuming the existence of $\int_0^\infty g(t)t^n dt$ and the termwise integrability of $\sum \frac{(-1)^n t^n}{n!} g(t)$. Then by making use of the relation

$$\frac{1}{p^{1/2}} = \int_0^\infty e^{-pt} \frac{dt}{\sqrt{\pi t}}$$

and the so-called "Faltungssatz," he deduced further

$$h(x) = \int_0^x \frac{g(t)}{\sqrt{\pi(x-t)}} dt$$

$$= \frac{1}{\sqrt{\pi x}} \int_0^x g(t) \left(\sum \frac{1 \cdot 3 \cdot 5 \dots (2n-1)}{n! (2x)^n} t^n \right) dt.$$

If the termwise integrability be assumed, it results

¹⁾ See Carson, Electric circuit theory and the operational calculus, 1926, Chap. V.