

3. A Construction of the Fundamental Solution for the Schrödinger Equations

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§ 1. Introduction. The aim of this note is to improve the results of [6], that is, to show that the main results of [6] hold even if we substitute the amplitude function $a(\lambda, t, s, x, y)$ of (10) in [6] by the constant function 1. We shall consider the Schrödinger equation

$$(1) \quad \frac{\partial}{\lambda \partial t} u(t, x) + \frac{1}{2} \sum_{j=1}^n \left(\frac{\partial}{\lambda \partial x_j} \right)^2 u(t, x) + V(t, x) u(t, x) = 0,$$

$$(t, x) \in \mathbf{R} \times \mathbf{R}^n$$

and the initial condition

$$(2) \quad u(s, x) = \varphi(x).$$

Here $\lambda = ih^{-1}$ is a pure imaginary parameter and h is a small parameter $0 < h \leq 1$. The potential $V(t, x)$ is assumed to satisfy the following two conditions;

(V-I) $V(t, x)$ is real valued. For any fixed $t \in \mathbf{R}$, $V(t, x)$ is a C^∞ function of $x \in \mathbf{R}^n$. $V(t, x)$ is measurable in $(t, x) \in \mathbf{R} \times \mathbf{R}^n$.

(V-II) For any multi-index α with length $|\alpha| \geq 2$, the non-negative measurable function of t defined by

$$(3) \quad M_\alpha(t) = \sup_{x \in \mathbf{R}^n} \left| \left(\frac{\partial}{\partial x} \right)^\alpha V(t, x) \right| + \sup_{|x| \leq 1} |V(t, x)|$$

is essentially bounded on every compact interval of \mathbf{R}^1 .

We fix $L \geq 10(m+n+10)$. We put $T = \infty$ if $\text{ess. sup.}_{|\alpha|=L, t \in \mathbf{R}} M_\alpha(t) < \infty$.

Otherwise we let T denote an arbitrarily fixed positive number. Every discussion will be made in the interval $(-T, T)$ throughout this paper.

We shall consider the integral transformation

$$(4) \quad E(\lambda, t, s)\varphi(x) = \left(\frac{-\lambda}{2\pi(t-s)} \right)^{(1/2)n} \int_{\mathbf{R}^n} e^{iS(t, s, x, y)} \varphi(y) dy,$$

where $S(t, s, x, y)$ is the classical action along the classical orbit starting the point y at the time s and reaching the point x at the time t . (If $|t-s|$ is small enough, such an orbit is uniquely determined. See Proposition 1 below.) The integral transformation (4) is exactly the same transformation as Feynman used in [3] and [4].

Let $[s, t] \subset (-T, T)$ be an arbitrary interval. Let

$$\Delta; s = t_0 < t_1 < t_2 < \dots < t_{L-1} < t_L = t$$

be an arbitrary subdivision of the interval $[s, t]$. We put