

## Research Article

# Fourteen Limit Cycles in a Seven-Degree Nilpotent System

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Center conditions and the bifurcation of limit cycles for a seven-degree polynomial differential system in which the origin is a nilpotent critical point are studied. Using the computer algebra system Mathematica, the first 14 quasi-Lyapunov constants of the origin are obtained, and then the conditions for the origin to be a center and the 14th-order fine focus are derived, respectively. Finally, we prove that the system has 14 limit cycles bifurcated from the origin under a small perturbation. As far as we know, this is the first example of a seven-degree system with 14 limit cycles bifurcated from a nilpotent critical point.

## 1. Introduction

In the qualitative theory of planar differential equations, the center-focus problem and bifurcation of limit cycles for nilpotent system

$$\begin{aligned} \frac{dx}{dt} &= y + \sum_{k+j=2}^{\infty} a_{kj} x^k y^j = X(x, y), \\ \frac{dy}{dt} &= \sum_{k+j=2}^{\infty} b_{kj} x^k y^j = Y(x, y), \end{aligned} \quad (1)$$

are known as a difficult problem. Some advance of this problem can be dated back to [1–3]. In recent years, due to the improvement of research method and development of computer symbolic computation, the problem has attracted more and more scholars' attention and has received a lot of results. For instance, in [4, 5], the center conditions of the nilpotent critical points were obtained for several systems. In [6] the center conditions and the bifurcations of limit cycles were investigated for a quintic and a nine-degree nilpotent systems. The center and the limit cycles problems of a quintic nilpotent system were also solved in [7]. And in [8], the authors gave a recursive method to calculate quasi-Lyapunov constants of the nilpotent critical point. The nilpotent center problem and limit cycles bifurcations were performed also in

[9]. It is interesting how many limit cycles can be bifurcated from the nilpotent critical point. Let  $N(n)$  be the maximum possible number of limit cycles bifurcated from a nilpotent critical point of system (1) when  $X$  and  $Y$  are of degree at most  $n$ . The known results of  $N(n)$  are: Andreev et al. given have  $N(3) \geq 2$ ,  $N(5) \geq 5$ ,  $N(7) \geq 9$ , see [5]. Y. Liu and J. Li showed  $N(3) \geq 4$ ,  $N(3) \geq 7$ ,  $N(3) \geq 8$ , see [8, 10–12]. Li et al. found  $N(7) \geq 12$  in [13]. Recently, Li et al. [14] obtained  $N(7) \geq 13$ .

In this paper, we study the bifurcation of limit cycles for a seven-degree nilpotent system with the following form:

$$\begin{aligned} \frac{dx}{dt} &= \delta x + y + a_{30} x^3 + a_{12} x y^2 + a_{32} x^3 y^2 + a_{14} x y^4 \\ &\quad + a_{05} y^5 + a_{06} y^6 + a_{15} x y^5 + a_{24} x^2 y^4 + a_{33} x^3 y^3 \\ &\quad + a_{51} x^5 y + a_{07} y^7 + a_{16} x y^6 + a_{25} x^2 y^5 \\ &\quad + a_{34} x^3 y^4 + a_{43} x^4 y^3 + a_{61} x^6 y, \\ \frac{dy}{dt} &= 2\delta y - 2x^3 + x y^2 + b_{33} x^3 y^3 + a_{51} x^4 y^2. \end{aligned} \quad (2)$$

By the computation of the quasi-Lyapunov constants, we prove that its perturbed system has 14 small-amplitude limit cycles bifurcated from the origin, namely,  $N(7) \geq 14$  which improves the result in [14].

In Section 2, we give some preliminary knowledge concerning the nilpotent critical point. In Section 3, we obtain the first 14 quasi-Lyapunov constants and derive the sufficient and necessary conditions of the origin to be a center and a 14th-order fine focus. At the end, it is proved that there exist 14 limit cycles in the neighborhood of the origin of the system.

## 2. Focal Values and Quasi-Lyapunov Constants

In order to discuss limit cycles of the system, we state some preliminary results given by [8].

According to [2], the origin of system is a 3th-order monodromic critical point and a center or a focus if and only if  $b_{20} = 0, (2a_{20} - b_{11})^2 + 8b_{30} \leq 0$ . Without loss of generality, we assume that  $a_{20} = \mu, b_{20} = 0, b_{11} = 2\mu, b_{30} = -2$ , otherwise let  $(2a_{20} - b_{11})^2 + 8b_{30} = -16\lambda^2, 2a_{20} + b_{11} = 4\lambda\mu$ .

Under the substitutions

$$\eta = \lambda y + \frac{1}{4}(2a_{20} - b_{11})^2 \lambda x^2 \quad \xi = \lambda x, \quad (3)$$

system (1) becomes

$$\begin{aligned} \frac{dx}{dt} &= y + \mu x^2 + \sum_{k+2j=3}^{\infty} a_{kj} x^k y^j = X(x, y), \\ \frac{dy}{dt} &= -2x^3 + 2\mu xy + \sum_{k+2j=4}^{\infty} b_{kj} x^k y^j = Y(x, y). \end{aligned} \quad (4)$$

By the transformation of the generalized polar coordinates,

$$x = r \cos \theta \quad y = r^2 \cos \theta, \quad (5)$$

system (4) is transformed into

$$\frac{dr}{d\theta} = \frac{\cos \theta R_1(\theta)}{Q_1(\theta)} + o(r), \quad (6)$$

where

$$R_1(\theta) = \sin \theta (1 - 2\cos^2 \theta) + \mu (\cos^2 \theta + 2\sin^2 \theta), \quad (7)$$

$$Q_1(\theta) = -2 (\cos^4 \theta + \sin^2 \theta) < 0.$$

For sufficiently small  $h$ , let

$$r = \tilde{r}(\theta, h) = \sum_{k=1}^{\infty} \nu_k(\theta) h^k \quad (8)$$

be a solution of (6) satisfying the initial value condition  $r|_{\theta=0} = h$ , where

$$\begin{aligned} \nu_1(\theta) &= (\cos^4 \theta + \sin^2 \theta)^{-1/4} \\ &\times \exp\left(\left(\frac{-\mu}{2}\right) \arctan\left(\frac{\sin \theta}{\cos^2 \theta}\right)\right), \quad (9) \\ \nu_1(k\pi) &= 1, \quad k = 0, \pm 1, \pm 2 \dots \end{aligned}$$

Because for all sufficiently small  $r$ , there is  $d\theta/dt < 0$ , in a small neighborhood; we obtain the Poincaré return map of (6) in a small neighborhood of the origin as follows:

$$\Delta(h) = \tilde{r}(-2\pi, h) - h = \sum_{k=2}^{\infty} \nu_k(-2\pi) h^k. \quad (10)$$

**Lemma 1.** For any positive integer  $m, \nu_{2m+1}(-2\pi)$  has the form

$$\nu_{2m+1}(-2\pi) = \sum_{k=1}^{\infty} \zeta_m^{(k)} \nu_{2k}(-2\pi), \quad (11)$$

where  $\zeta_m^{(k)}$  is a polynomial of  $\nu_i(\pi), \nu_i(2\pi), \nu_i(-2\pi), (i = 2, 3, \dots, 2m)$  with rational coefficients.

**Definition 2.** (i) For any positive integer  $m, \nu_{2m}(-2\pi)$  is called the  $m$ th-order focal value of system (4) at the origin; (ii) if  $\nu_2(-2\pi) \neq 0$ , the origin of system (4) is called an 1th-order weak focus; if there is an integer  $m > 1$  such that  $\nu_2(-2\pi) = \nu_4(-2\pi) = \dots = \nu_{2m-2}(-2\pi) = 0, \nu_{2m}(-2\pi) \neq 0$ , then the origin of system (4) is called a  $m$ th-order weak focus; (iii) if for all positive integer  $m$ , we have  $\nu_{2m}(-2\pi) = 0$ , the origin of system (4) is called a center.

**Lemma 3.** For system (4), one can derive successively the formal series

$$M(x, y) = x^4 + y^2 + o(r^4) \quad (12)$$

such that

$$\begin{aligned} \left(\frac{\partial X}{\partial x} + \frac{\partial Y}{\partial y}\right) M - (s+1) \left(\frac{\partial M}{\partial x} X + \frac{\partial M}{\partial y} Y\right) \\ = \sum_{m=1}^{\infty} \lambda_m [(2m - 4s - 1) x^{2m+4} + o(r^{2m+4})]. \end{aligned} \quad (13)$$

**Lemma 4.** If there exists a natural number  $s$  and formal series

$$M(x, y) = x^4 + y^2 + o(r^4) \quad (14)$$

such that (13) holds, then

$$\nu_{2m}(-2\pi) \sim \sigma_m \lambda_m, \quad m = 1, 2, 3, \dots, \quad (15)$$

where

$$\begin{aligned} \sigma_m &= \frac{1}{2} \int_0^{2\pi} (1 + \sin^2 \theta) \cos^{2m+4} \theta \\ &\times \left( (\sin^4 \theta + \sin^2 \theta)^{2m+7/4} \right. \\ &\times \exp\left(\left(2m - \frac{1}{2}\right) \mu \arctan \frac{\sin \theta}{\cos \theta}\right) \left. \right)^{-1} d\theta > 0. \end{aligned} \quad (16)$$

In (15),  $\sim$  is the symbol of algebraic equivalence, meaning that there exists  $\xi_m^{(k)} (k = 1, 2, \dots, m-1)$ , polynomial functions of the coefficients of system (4), such that

$$\nu_{2m+1}(-2\pi) = \sigma_m \lambda_m + \sum_{k=1}^{m-1} \xi_m^{(k)} \lambda_k. \quad (17)$$

**Definition 5.** In Lemma 4,  $\lambda_m$  is called the  $m$ th-order quasi-Lyapunov constant of the origin of system (4).

**Lemma 6.** For system (4), one can derive successively the formal series

$$M(x, y) = y^2 + \sum_{\alpha+\beta=3}^{\infty} c_{\alpha\beta} x^\alpha y^\beta \tag{18}$$

such that

$$\begin{aligned} & \left( \frac{\partial X}{\partial x} + \frac{\partial Y}{\partial y} \right) M - (s+1) \left( \frac{\partial M}{\partial x} X + \frac{\partial M}{\partial y} Y \right) \\ &= \sum_{m=3}^{\infty} \omega_m(s, \mu) x^m, \end{aligned} \tag{19}$$

where  $c_{00} = c_{10} = c_{01} = c_{20} = c_{11} = 0, c_{02} = 1$ . For  $\alpha \geq 1, \alpha + \beta \geq 3, c_{\alpha\beta}$ , and  $\omega_m(s, \mu)$  are determined by the following recursive formulas:

$$\begin{aligned} c_{\alpha\beta} &= \frac{1}{(s+1)\alpha} (A_{\alpha-1, \beta+1} + B_{\alpha-1, \beta+1}), \\ \omega_m(s, \mu) &= A_{m,0} + B_{m,0}, \end{aligned} \tag{20}$$

where

$$\begin{aligned} A_{\alpha\beta} &= \sum_{k+j=2}^{\alpha+\beta-1} [k - (s+1)(\alpha - k + 1)] a_{kj} c_{\alpha-k+1, \beta-j}, \\ B_{\alpha\beta} &= \sum_{k+j=2}^{\alpha+\beta-1} [j - (s+1)(\beta - j + 1)] b_{kj} c_{\alpha-k, \beta-j+1}. \end{aligned} \tag{21}$$

By choosing  $\{c_{0\beta}\}$  such that

$$\omega_{2k+1}(s, \mu) = 0, \quad k = 1, 2, \dots, \tag{22}$$

one has

$$\lambda_m = \frac{\omega_{2m+4}(s, \mu)}{2m - 4s - 1}. \tag{23}$$

One considers the perturbed system of system (4)

$$\begin{aligned} \frac{dx}{dt} &= \delta x + y + \mu x^2 + \sum_{k+2j=3}^{\infty} a_{kj} x^k y^j, \\ \frac{dy}{dt} &= 2\delta y - 2x^3 + 2\mu xy + \sum_{k+2j=4}^{\infty} b_{kj} x^k y^j. \end{aligned} \tag{24}$$

For system (24) $_{\delta=0}$ , from Lemma 4, we know that the first nonvanishing quasi-Lyapunov constant  $\lambda_m$  is positive constant times as much as the first nonvanishing focal value, so the former shows the same effect as the latter in the study of bifurcation of limit cycles. From [10, Theorem 4.7], we have the following.

**Theorem 7.** For the system (27) $_{\delta=0}$ , assume that the quasi-Lyapunov constants of the origin  $\lambda_i$  ( $i = 1, 2, \dots$ ) have  $k$

independent parameters  $\gamma = (\gamma_1, \gamma_2, \dots, \gamma_k)$ ; that is,  $\lambda_i = \lambda_i(\gamma_1, \gamma_2, \dots, \gamma_k)$ . If  $\gamma = \gamma_0$ , the origin of the system (4) is an  $m$ th-order weak focus ( $m \leq k$ ), and the Jacobian determinant

$$\left. \frac{\partial(\lambda_1, \lambda_2, \dots, \lambda_{m-1})}{\partial(\gamma_1, \gamma_2, \dots, \gamma_{m-1})} \right|_{\gamma=\gamma_0} \neq 0, \tag{25}$$

then, the perturbed system (24) exists  $m$  small amplitude limit cycles bifurcated from the origin.

### 3. Criterion of Center Focus and Bifurcation of Limit Cycles

Applying the recursive formulas in Lemma 6, we compute the quasi-Lyapunov constants of the origin of system (2) $_{\delta=0}$  with the computer algebra system Mathematica and obtain the following result.

**Theorem 8.** For system (2) $_{\delta=0}$ , the first 14 quasi-Lyapunov constants are as follows:

$$\begin{aligned} \lambda_1 &= a_{30}, \\ \lambda_2 &= \frac{2}{5} a_{12}, \\ \lambda_3 &= \frac{2}{7} a_{32}, \\ \lambda_4 &= \frac{4}{15} a_{14}, \\ \lambda_5 &= \frac{12}{77} a_{34}, \\ \lambda_6 &= \frac{2}{195} (20a_{16} + 3a_{51}b_{33}), \\ \lambda_7 &= \frac{1}{385} b_{33} (35a_{51} - 8a_{33}), \\ \lambda_8 &= \frac{7}{13260} b_{33} (128a_{15} - 355a_{51}), \\ \lambda_9 &= \frac{3}{33440} b_{33} a_{51} (1385 + 64a_{61}), \\ \lambda_{10} &= \frac{1}{278460} b_{33} a_{51} \\ &\quad \times (-192495 + 12320a_{05} + 1904a_{43}), \\ \lambda_{11} &= \frac{9}{1184444800} b_{33} a_{51} \\ &\quad \times (317763455 + 1688064a_{43} + 1158080a_{51}^2), \\ \lambda_{12} &= \frac{1}{505504614521088000} b_{33} a_{51} \\ &\quad \times (424870735079675775 - 8480461063976518a_{51}^2 \\ &\quad \quad - 164955456258816b_{33}^2), \end{aligned}$$

$$\lambda_{13} = \frac{1}{2497759223828804812800} \times b_{33} a_{51} \left( 1154557205782671354192175 - 25287050037965301847744a_{51}^2 \right),$$

$$\lambda_{14} = -\frac{1}{1926846314779614102444810240000} b_{33} a_{51} \times \left( 1913839774991447312487020909964625 - 38616043776955260227746202006848a_{51}^2 + 457974511144735287048192000a_{51}^4 \right). \tag{26}$$

Here, every  $\lambda_k$  ( $k = 1, 2, \dots, 14$ ) was computed under the assumption  $\lambda_1 = \lambda_2 = \dots = \lambda_{k-1} = 0$ .

It is easy to obtain the following Theorem.

**Theorem 9.** For system (2)| $_{\delta=0}$ , the first 14 quasi-Lyapunov constants at the origin are all zero if and only if the following condition is satisfied:

$$a_{30} = a_{12} = a_{32} = a_{14} = a_{34} = a_{51} = a_{33} = a_{15} = a_{16} = 0. \tag{27}$$

If  $\delta = 0$  and the condition (27) holds, system (2) becomes

$$\frac{dx}{dt} = y + a_{05}y^5 + a_{06}y^6 + a_{24}x^2y^4 + a_{07}y^7 + a_{25}x^2y^5 + a_{43}x^4y^3 + a_{61}x^6y, \tag{28}$$

$$\frac{dy}{dt} = -2x^3 + xy^2 + b_{33}x^3y^3,$$

which is symmetric with respect to the  $y$ -axis, one has the following.

**Theorem 10.** The origin of system (2) is a center if and only if  $\delta = 0$  and (27) holds.

By  $\lambda_1 = \lambda_2 = \dots = \lambda_{13} = 0, \lambda_{14} \neq 0$ , one has the following.

**Theorem 11.** The origin of system (2) is a 14th-order weak focus if and only if

$$\delta = a_{30} = a_{12} = a_{32} = a_{14} = a_{34} = 0,$$

$$a_{61} = -\frac{1385}{64},$$

$$a_{05} = \frac{30075794600575314214479775}{606889200911167244345856},$$

$$a_{43} = -\frac{66625696625444520068811785}{303444600455583622172928},$$

$$b_{33}^2 = \frac{10913994716347225847247003725}{4779252457175442049223616},$$

$$a_{51}^2 = \frac{1154557205782671354192175}{25287050037965301847744},$$

$$a_{16} = -\frac{3}{20}a_{51}b_{33}, \quad a_{33} = \frac{35}{8}a_{51}, \quad a_{15} = \frac{355}{128}a_{51}. \tag{29}$$

By computing carefully, we obtain that the Jacobian determinant

$$\frac{\partial(\lambda_1, \lambda_2, \lambda_3, \lambda_4, \lambda_5, \lambda_6, \lambda_7, \lambda_8, \lambda_9, \lambda_{10}, \lambda_{11}, \lambda_{12}, \lambda_{13})}{\partial(a_{30}, a_{12}, a_{32}, a_{14}, a_{34}, a_{16}, a_{33}, a_{15}, a_{61}, a_{05}, a_{43}, a_{51}, b_{33})} \Big|_{(29)}$$

$$= -\frac{11259131158497337756164795883686035195310097999201613627491381814272a_{51}^4 b_{33}^6}{110636634525265639383282317978327920684865639296808353136757452754003615234375}$$

$$\approx -2526.4563514134 \neq 0. \tag{30}$$

From (30) and Theorem 7, one has the following.

**Theorem 12.** For system (2), under the condition (29), by small perturbations of the parameter group  $(\delta, a_{30}, a_{12}, a_{32}, a_{14}, a_{34}, a_{16}, a_{33}, a_{15}, a_{61}, a_{05}, a_{43}, a_{51}, b_{33})$ , then there are 14 small amplitude limit cycles bifurcated from the origin.

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**References**

- [1] A. F. Andreev, "Investigation of the behaviour of the integral curves of a system of two differential equations in the neighbourhood of a singular point," *American Mathematical Society Translations*, vol. 8, pp. 183–207, 1958.
- [2] V. V. Amel'kin, N. A. Lukashevich, and A. P. Sadovskii, *Non-linear Oscillations in Second Order Systems*, Belarusian State University, Minsk, Russia, 1982, (Russian).
- [3] V. G. Romanovskii, "On the cyclicity of the equilibrium position of the center or focus type of a certain system," *Vestnik St. Petersburg University: Mathematics*, vol. 19, pp. 51–56, 1986.
- [4] M. J. Álvarez and A. Gasull, "Monodromy and stability for nilpotent critical points," *International Journal of Bifurcation and Chaos*, vol. 15, no. 4, pp. 1253–1265, 2005.

- [5] A. F. Andreev, A. P. Sadovskii, and V. A. Tsikalyuk, "The center-focus problem for a system with homogeneous nonlinearities in the case of zero eigenvalues of the linear part," *Differential Equations*, vol. 39, no. 2, pp. 155–164, 2003.
- [6] M. J. Álvarez and A. Gasull, "Generating limit cycles from a nilpotent critical point via normal forms," *Journal of Mathematical Analysis and Applications*, vol. 318, no. 1, pp. 271–287, 2006.
- [7] A. Algaba, C. García, and M. Reyes, "Local bifurcation of limit cycles and integrability of a class of nilpotent systems of differential equations," *Applied Mathematics and Computation*, vol. 215, no. 1, pp. 314–323, 2009.
- [8] Y. Liu and J. Li, "On third-order nilpotent critical points: integral factor method," *International Journal of Bifurcation and Chaos*, vol. 21, no. 5, pp. 1293–1309, 2011.
- [9] M. Han and V. G. Romanovski, "Limit cycle bifurcations from a nilpotent focus or center of planar systems," *Abstract and Applied Analysis*, vol. 2012, Article ID 720830, 28 pages, 2012.
- [10] Y. Liu and J. Li, "New study on the center problem and bifurcations of limit cycles for the Lyapunov system. I," *International Journal of Bifurcation and Chaos*, vol. 19, no. 11, pp. 3791–3801, 2009.
- [11] Y. Liu and J. Li, "New study on the center problem and bifurcations of limit cycles for the Lyapunov system. II," *International Journal of Bifurcation and Chaos*, vol. 19, no. 9, pp. 3099–3807, 2009.
- [12] Y. Liu and J. Li, "Bifurcations of limit cycles and center problem for a class of cubic nilpotent system," *International Journal of Bifurcation and Chaos*, vol. 20, no. 8, pp. 2579–2584, 2010.
- [13] F. Li, Y. Liu, and Y. Wu, "Center conditions and bifurcation of limit cycles at three-order nilpotent critical point in a seventh degree Lyapunov system," *Communications in Nonlinear Science and Numerical Simulation*, vol. 16, no. 6, pp. 2598–2608, 2011.
- [14] F. Li, Y. Liu, and H. Li, "Center conditions and bifurcation of limit cycles at three-order nilpotent critical point in a septic Lyapunov system," *Mathematics and Computers in Simulation*, vol. 81, no. 12, pp. 2595–2607, 2011.